

# COLLIDER IMPLICATIONS OF MODELS WITH EXTRA DIMENSIONS

ICHEP 2002, SESSION 10

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: Brief review of extra dimension

Symmetric  
Compactification

argm.  
Compact.

: Models with  $\text{TeV}^{-1}$  scale extra dim

: Collider signals

: only gauge bosons in extra dim

: all SM particles in extra dim  
(Universal Extra Dim)

: Conclusions

①

EXTRA DIMENSIONS  $\rightarrow$  motivated by string theory  
but may be more general

Superstring Theory  $\rightarrow$  1 + 9  
                                   $\uparrow$                    $\uparrow$   
                                  time                  space

$\Rightarrow$  6 extra space-like dim  
 $\Rightarrow$  must be compact

: Physics depends on how compactified  
and at what scale

TWO BROAD POSSIBILITIES

all sizes  $\sim M_p^{-1}$

$\downarrow$   
Very little  
phenomenological  
implications

(Witten, Lykken)

Some or all much larger  
than  $M_{pl}^{-1}$

$\sim$  submm,  $TeV^{-1}$ , .....

$\Downarrow$   
many interesting phenomeno-  
logical consequences

- : modify Newton's law of gravity
- : effects in low energy astrophysical phenomena
- : new physics in high energy colliders.

(2)

High Energy Colliders  $\Rightarrow$  new physics from the KK excitations of gravity and/or SM fields.

: Where do the SM particles live??

4 + N DIM SPACE

TWO SCENARIOS

SM particles do not see any extra dim (confined to  $D_3$  brane)  
 $\Rightarrow$  only gravity propagate to extra dim.  
 $\Rightarrow$  ADD scenario

Some or all SM particles do see one or more extra dimensions  
(Antoniadis, 1991)  
Dienes, Dudas, Ghazizadeh, Lykken, Nandi)

: For time limitation, I will not talk about ADD or RS Models.

Contribution by Kim, Lee and Song

$\Rightarrow$  interesting enhancement of Higgs pair productions in ADD  
 $\hookrightarrow$  RS models at the LHC

3

: how extra dims are compactified?

$$D = 4 + N$$

Symmetric Compactification:

All extra dim. sym., same size  $R$

Scaling law:  $M_{\text{PL}}^2 = M_D^{n+2} R^n$

$N=1$  excluded,  $N=2 \sim \text{sub mm}$

$N=6, \sim (10 \text{ MeV})^{-1}$

$\Rightarrow$  all SM particles must be confined to  $D_3$ -brane

Asymmetric compactification:

$N$  extra dim diff sizes and shapes

Simplest possibility: Two sizes

$r \rightarrow \text{sub mm}$

$R \rightarrow \text{TeV}^{-1}$

Scaling law:  $M_{\text{PL}}^2 = M_D^{n+m+2} R^n r^m$

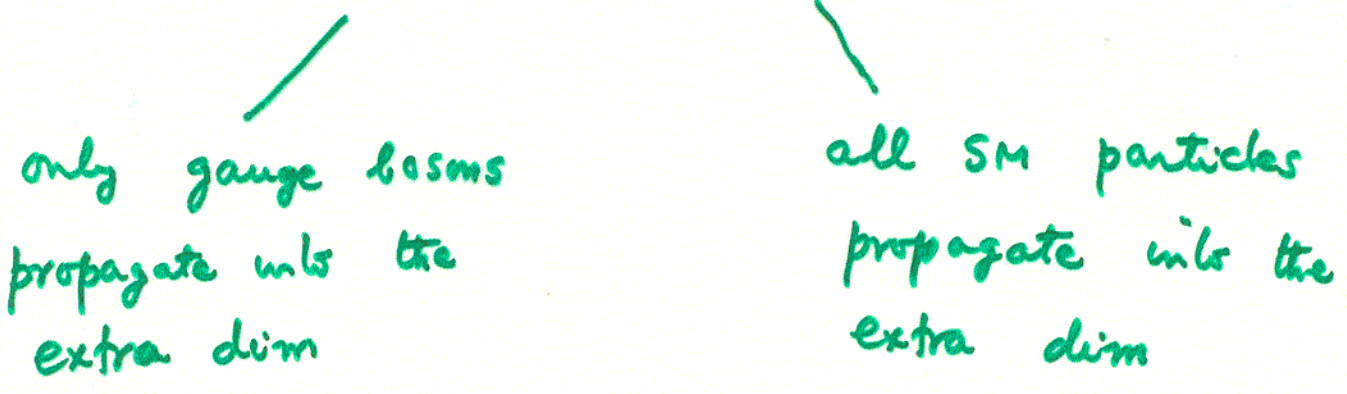
$$n + m \leq 6$$

For simplicity, we consider

- :  $1 \text{ TeV}^{-1}$  extra dim where SM particles propagate
- : 5 much larger extra dim where gravity propagate

$1 \text{ TeV}^{-1}$  scale extra dim

Two Possibilities



⇒ SM will have KK towers of excitations with masses starting at  $\sim \text{TeV}$



Collider implications due to the productions or off shell effects of those KK excitations

- ⇒ : new signals
- : modification of SM signals

5

SCENARIO 1:

: Only gauge bosons see the extra dim,  
fermions confined to D3 brane

gauge bosons  $\Rightarrow$  KK excitations

$$g \rightarrow g_n^* , \quad W, Z, \gamma \rightarrow W_n^*, Z_n^*, \gamma_n^*$$

$$M_n^2 = \frac{n^2}{R^2} + m_0^2$$

$\rightarrow$  mass of  $g, W, Z, \gamma$

: 5 dim momentum or KK # not conserved  
(Brane can absorb momenta in the 5th dir.)

new couplings:  $\bar{e}e g^* , \bar{e}e W^* , \dots$

new collider processes:

$$\begin{array}{ccc}
 \bar{e}e \rightarrow g_n^* \rightarrow \bar{e}e & , & \bar{e}e \rightarrow g g_n^* & , & \bar{e}e \rightarrow g^* g^* \\
 \downarrow & & \downarrow & & \downarrow \\
 \text{dijet} & & \text{2j} & & \text{4jets} \\
 & & \downarrow & & \\
 & & \text{3jet} & & 
 \end{array}$$

observable signal above background

$\Rightarrow$  modification of dijet crosssections at high  $p_T$   
at Tevatron & LHC

Run 1 bounds,  $1/R \equiv \mu \gtrsim \text{TeV}$

Run 2  $\Rightarrow \mu \gtrsim 2 \text{TeV}$

LHC  $\Rightarrow \mu \gtrsim 7 \text{TeV}$

Figs

Dicus, McMullen + Nandi: PRD

56 S

TEVATRON RUN 2,  $\sqrt{s} = 2\text{TeV}$

DIJETS +  $\cancel{E}_T$  signal

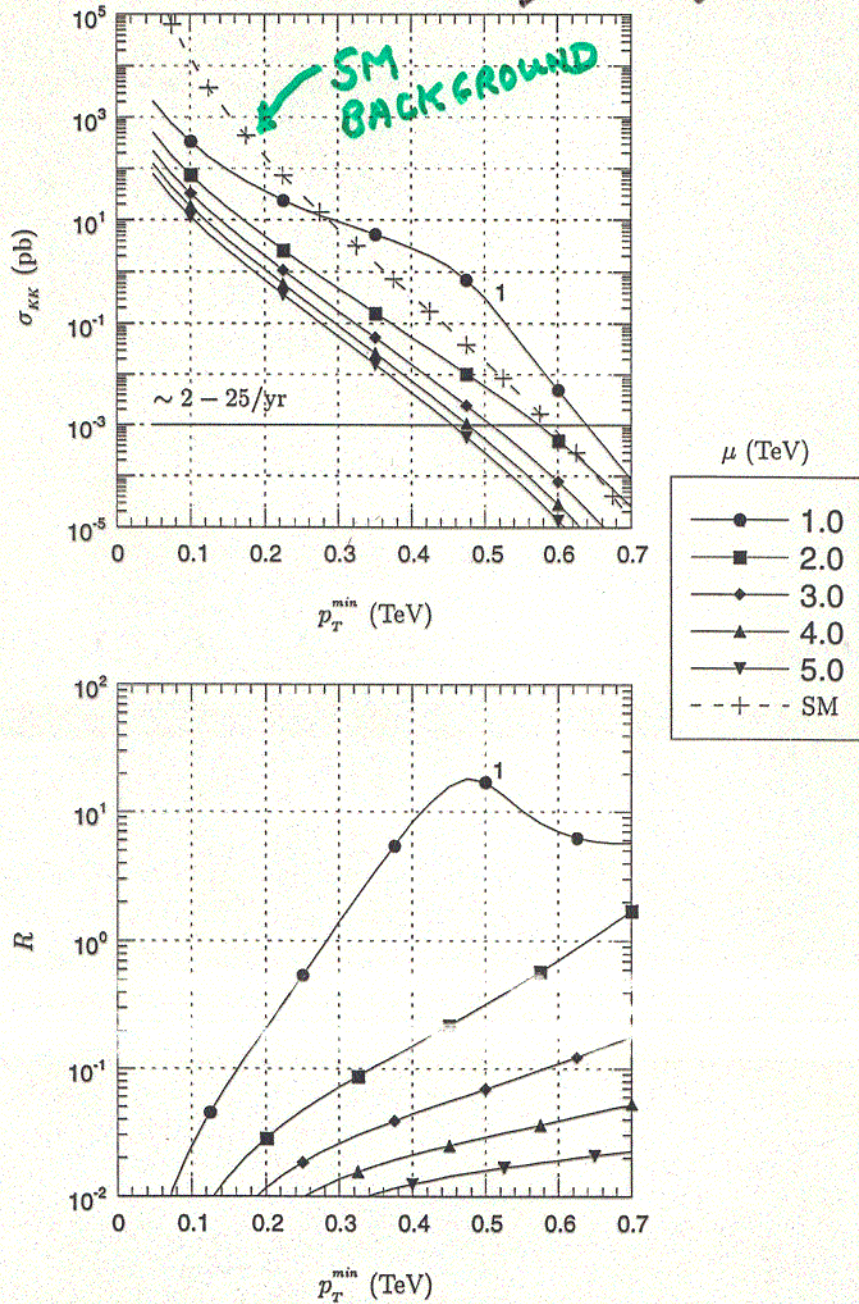


Figure 7: The contributions of the virtual exchanges of  $g^*$ 's to the Tevatron dijet production cross section,  $\sigma_{KK} = \sigma - \sigma_{SM}$ , (top) and the ratio of the KK contribution to the SM background,  $R = \sigma_{KK}/\sigma_{SM}$ , (bottom) are illustrated as a function of the minimum transverse momentum  $p_T^{\min}$  for fixed values of the compactification scale  $\mu$ . The solid horizontal line represents  $\sim 2$  (25) events/yr at the projected initial (final) Run 2 integrated luminosity. Discernible bumps in regions for which  $p_T^{\min} = k\mu/2$  are indicated by the corresponding value of  $k \in \{1, 2, \dots\}$ .

LHC,  $\sqrt{s} = 14 \text{ TeV}$   
 Dijet +  $E_T$  signal

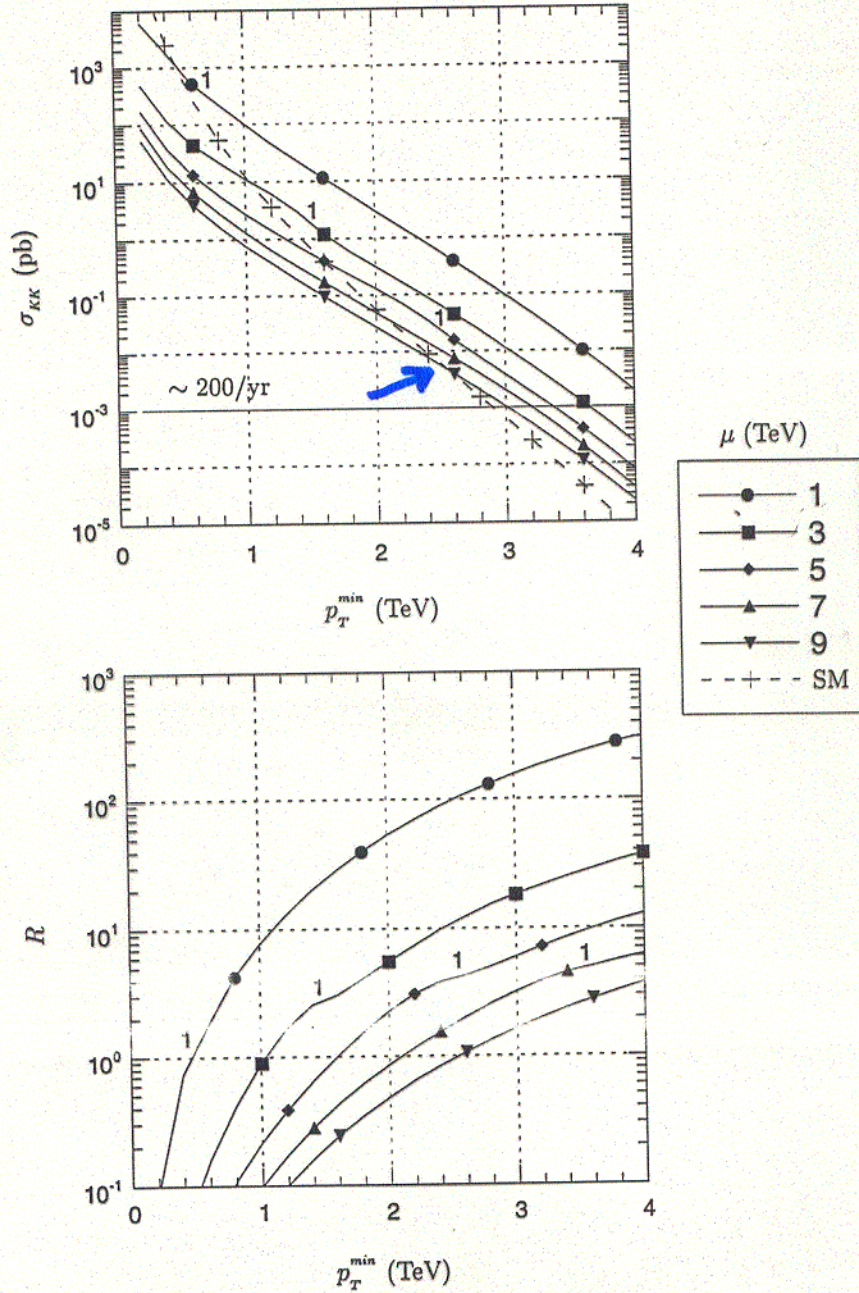


Figure 3: The contributions of the virtual exchanges of  $g^*$ 's to the LHC dijet production cross section,  $\sigma_{KK} = \sigma - \sigma_{SM}$ , (top) and the ratio of the KK contribution to the SM background,  $R = \sigma_{KK}/\sigma_{SM}$ , (bottom) are illustrated as a function of the minimum transverse momentum  $p_T^{\min}$  for fixed values of the compactification scale  $\mu$ . The solid horizontal line represents  $\sim 200$  events/yr at the projected integrated luminosity. Discernible bumps in regions for which  $p_T^{\min} = k\mu/2$  are indicated by the corresponding value of  $k \in \{1, 2, \dots\}$ .



6

## SCENARIO 2

: All SM particles propagate into  $\text{TeV}^{-1}$  size extra dim.  $\Rightarrow$  Universal Extra Dim (UED)  
 (Appelquist, Cheng, Dobrescu)  
 $\Rightarrow$  all SM have KK excitations

: consider one  $\text{TeV}^{-1}$  size extra dim.

: 5D fermions are 4-compo, vector-like Dirac spinors

$\Downarrow$

use  $S_1/Z_2$  orbifold to project out the fermions with wrong chirality zero modes.

$$Q(x,y) = \frac{1}{\sqrt{\pi R}} \left\{ \begin{pmatrix} u(x) \\ d(x) \end{pmatrix}_L + \sqrt{2} \sum_{n=1}^{\infty} \left[ Q_L^{(n)}(x) \cos \frac{ny}{R} + Q_R^{(n)}(x) \sin \frac{ny}{R} \right] \right\}$$

$$U(x,y) = \frac{1}{\sqrt{\pi R}} \left\{ U_R(x) + \sqrt{2} \sum_{n=1}^{\infty} \left[ U_R^{(n)}(x) \cos \frac{ny}{R} + U_L^{(n)} \sin \frac{ny}{R} \right] \right\}$$

$$Q_L^{(n)}, Q_R^{(n)} \Rightarrow Q^{(n)}, \quad U_R^{(n)}, U_L^{(n)} \rightarrow U^{(n)}$$

$\Rightarrow$  two KK towers for each flavor  $q$ .

$$M_n = \sqrt{n^2/R^2 + m_q^2}$$

$\swarrow$  zero mode mass  
 $\searrow$  from Yukawa couplings

$$Q^{(n)} \rightarrow q^{\circ}, \quad U^{(n)} \rightarrow q_n^{\circ}$$

$g_n^*$  → KK excitation of gluons,  $W_n^*$ ,  $Z_n^*$ ,  $\gamma_n^*$

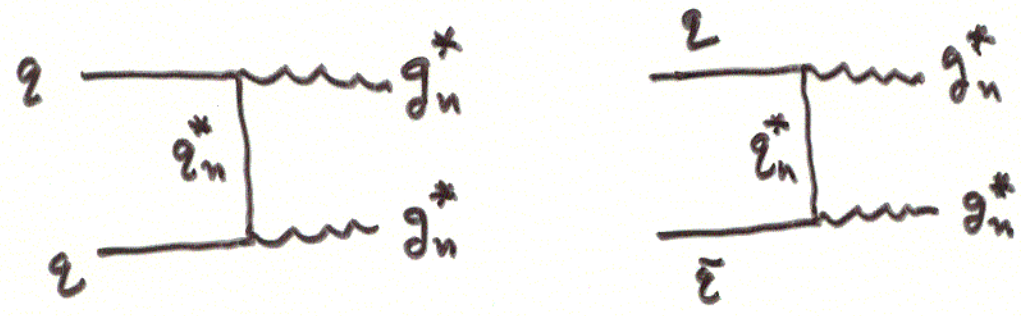
$q_n^0, q_n^0$  → two towers for each flavor of quarks (also leptons)

new ⇒ KK number is conserved at tree level



⇒ forbidden

⇒ KK particles must be pair produced



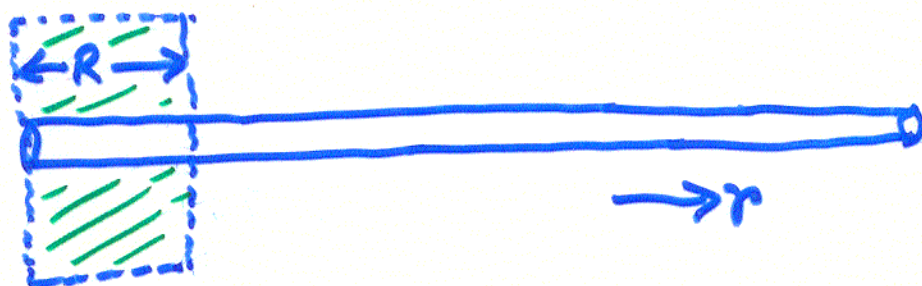
⇒ Mass bounds get reduced

from 1 TeV to  $\sim \sqrt{1 \text{ TeV}^2} \sim 300 \text{ GeV}$

9

# SIGNALS WITH KK NUMBER VIOLATING TREE LEVEL INTERACTION :

: Use "FAT BRANE SCENARIO"



$$r \sim \text{Submm}$$

$$1/R \sim \text{TeV}$$

: coupling of the KK violating interaction is proportional to the overlap of the wave functions in the 5th dim. The extra momenta along the  $y$  dir from the KK violation is absorbed by the brane.

$\Rightarrow$  KK states can decay via grav. int.

$$g^* \rightarrow g G_m, \quad q^* \rightarrow q G_m$$

: Life-time depends on  $N = \#$  of extra dim grav.

$\rightarrow$  always fast enough for  $g^*, q^*$  propagate the detector.

Signal  $\Rightarrow$  dijets events with high  $p_T$  + large missing energy  $\gg$  Background

Run 1 bound:  $\sim 400$  GeV, Run 2:  $\sim 450-500$  GeV

LHC:  $\sim 3$  TeV

figs

95

Macesanu, McMullen + Nandori, hep-ph/0210300

UED, DIJETS  
TEVATRON RUN2

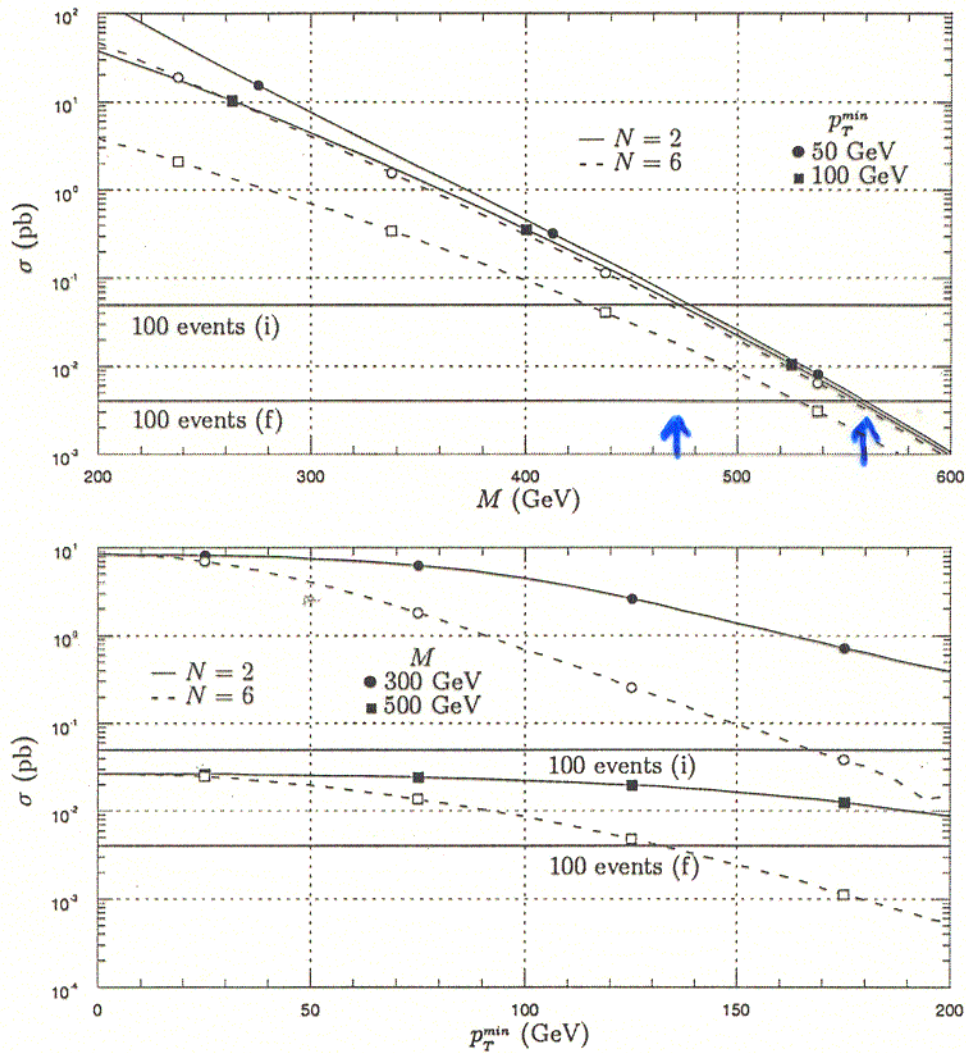


Figure 7: The total cross section for the dijet production plus missing energy from decaying KK final states at the Tevatron Run II energy is illustrated as a function of  $\mu$  for fixed  $p_T^{\min}$  (top) and as a function of the minimum transverse momentum  $p_T^{\min}$  for fixed values of the compactification scale  $\mu$  (bottom). Solid horizontal lines mark 100 events at the initial and final projected luminosities. In this and the following figures, we implement cuts on the  $p_T$ , rapidity, and separation of the jets.

Macesanu, McMullen + Nandori

hep-ph/0210300

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UED, DIJETS

LHC

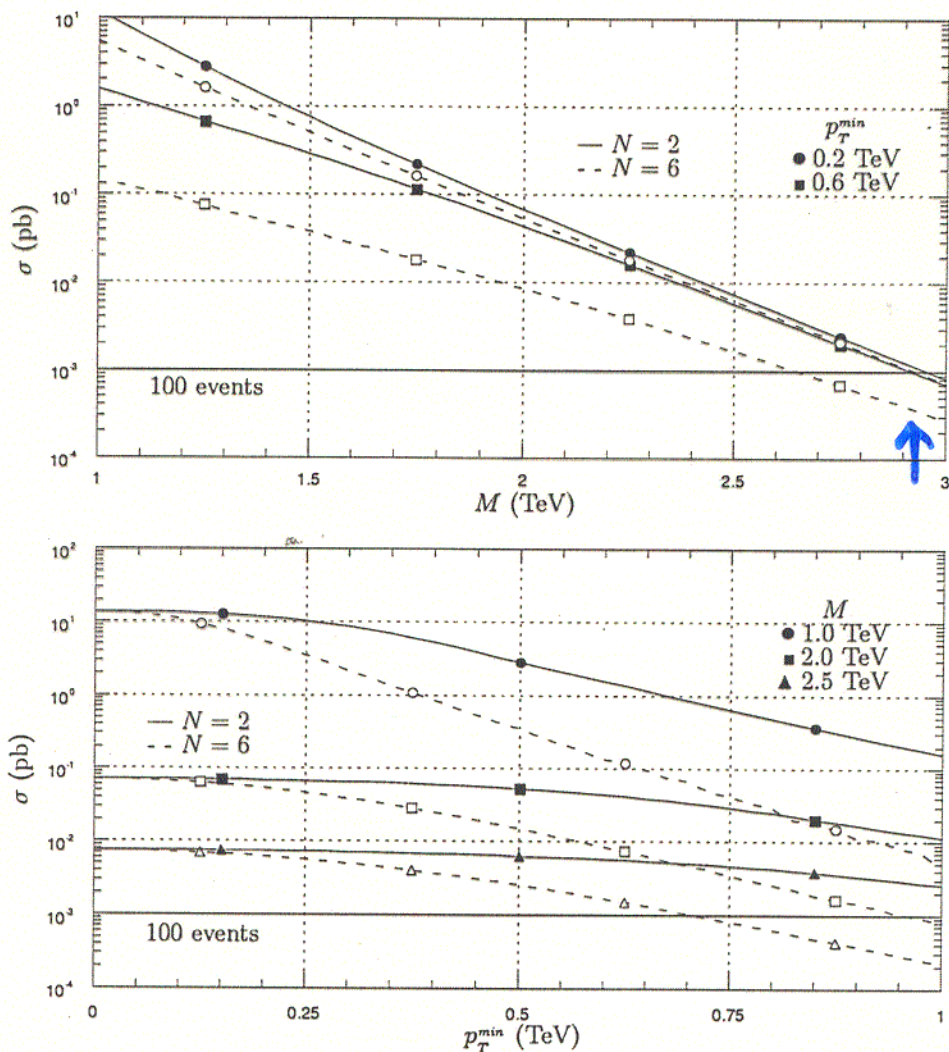


Figure 8: The same as Fig. 7, but for the LHC. The solid horizontal line marks 100 annual events at the projected luminosity.

SCENARIO B Include one loop correction to the KK masses, but KK violating interaction absent (Cheng, Matchov, Martin)

: degeneracy of KK excitations split

1st level:  $g_1^*$  → heaviest (20-30% correction)  
~ 150-200 GeV split for  $M R \sim 20$

next:  $q_1^0, q_1^+, W_1^+, Z_1^+, h_1^+, l_1^0, \dots$

Lightest KK particle  $\Rightarrow \gamma_1^*$  (LKP)

$\Rightarrow$  STABLE  $\hookrightarrow \sim B^*$

$\Rightarrow$  Decays allowed

$$g_1^* \rightarrow q_1^0 \bar{q}, q_1^+ \bar{q}$$

$$q_1^+ \rightarrow q Z_1^* \rightarrow q l l_1^* \rightarrow q l l \gamma_1^*, BR \sim 33\%$$

$$q_1^+ \rightarrow q W_1^* \rightarrow q e' l_1^* \rightarrow q e' l \gamma_1^*, BR \sim 65\%$$

: Pair produced at collider, after decay chain,

$\Rightarrow$  end up with  $\gamma_1^* \gamma_1^*$  + soft like quarks and leptons.

Signal  $\Rightarrow$  soft leptons + missing energy (like SUSY with almost deg. superpartners)

Reach: Run2: ~ 350 GeV  
LHC: ~ 1.5 TeV

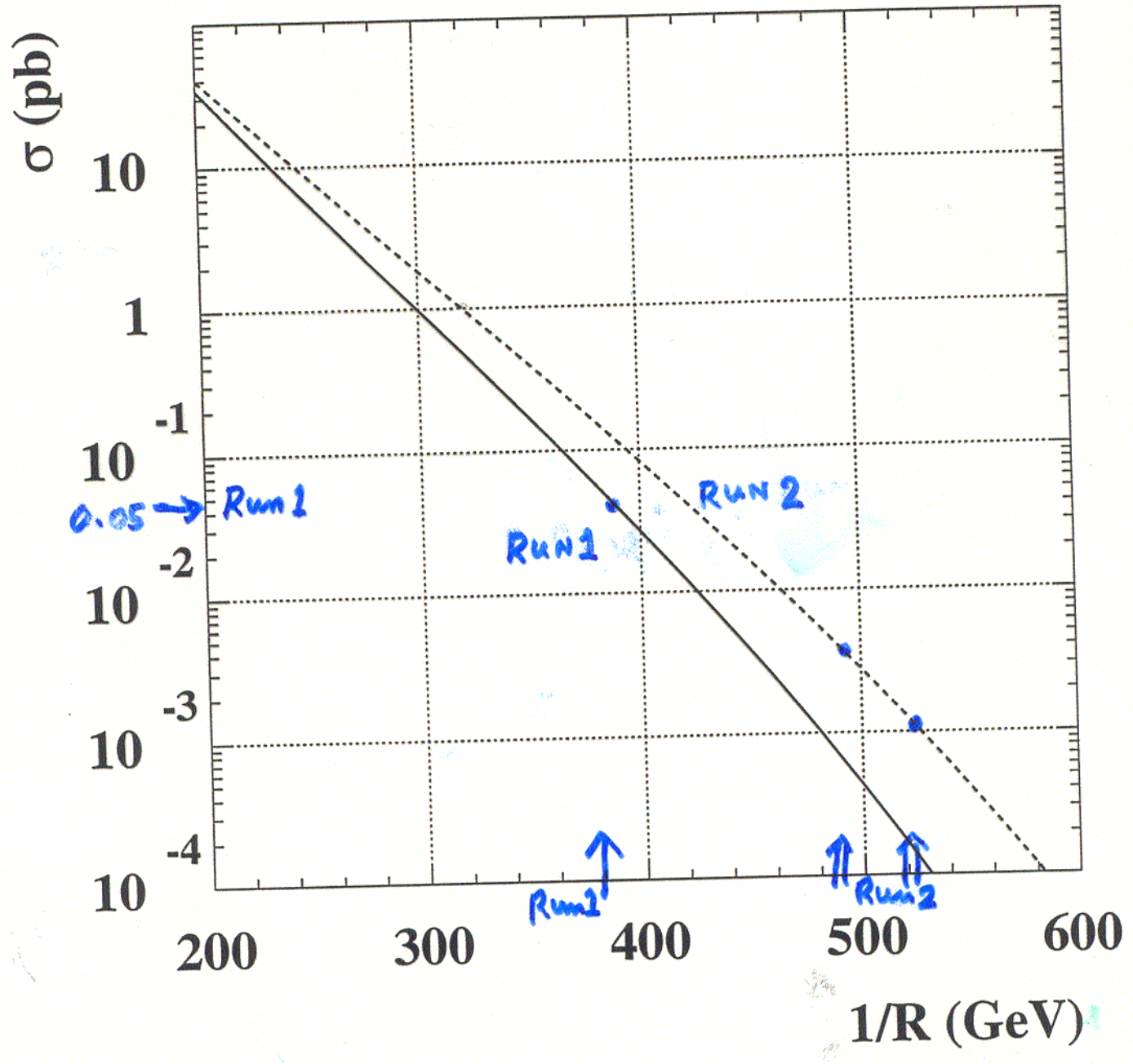
Macesanu, McMullen + Nandi, hep-ph/0207269

TEVATRON,

UED, 1 LOOP MASS CORRECTIONS,

$\gamma\gamma + \cancel{e}e$  SIGNAL

BACKGD  $\sim 0.6$  fb

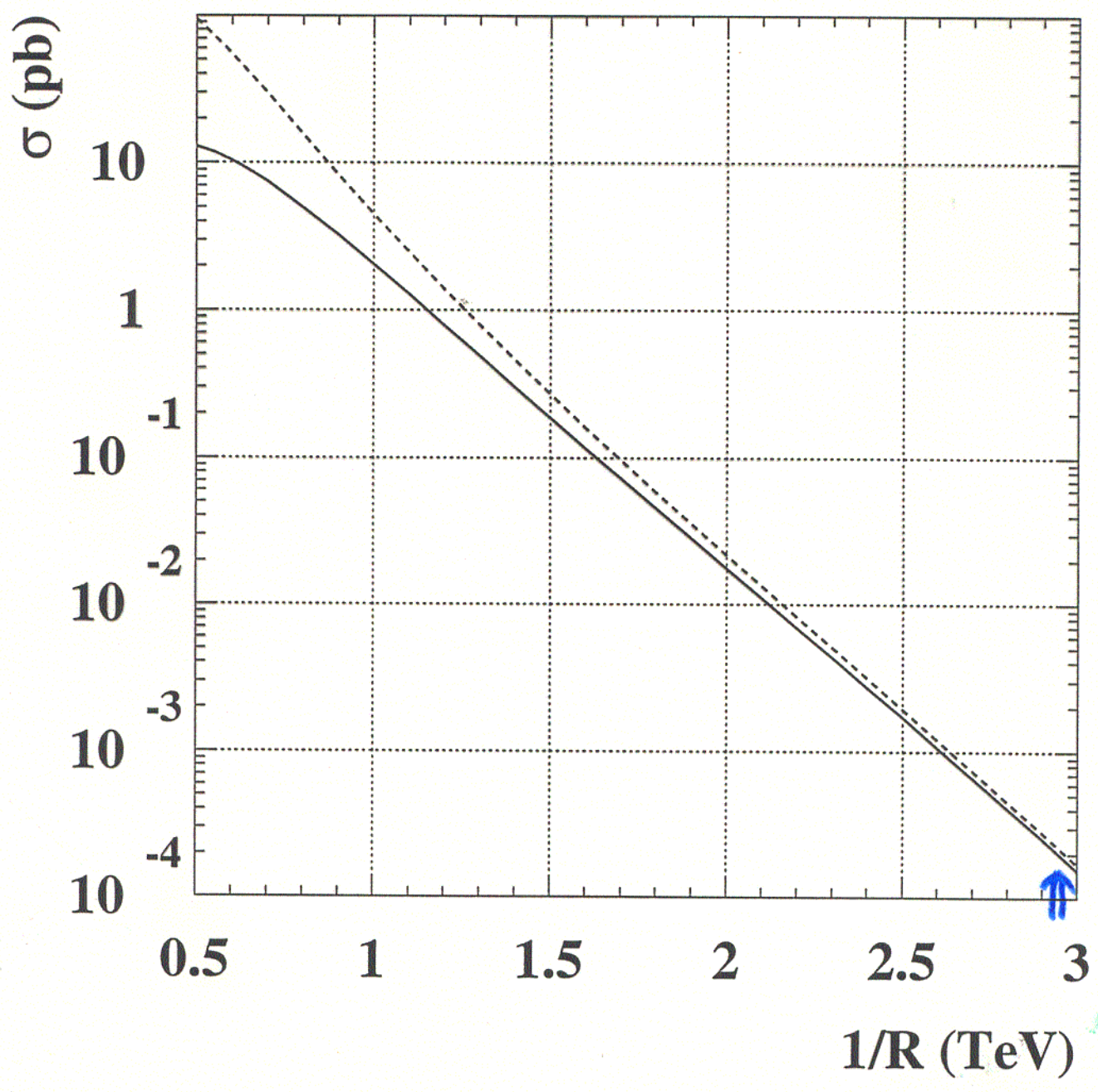


Macesanu, McMullen + Nandi, hep-ph/0207269

UED, 1 LOOP MASS CORRECTED

LHC

$\gamma\gamma + \cancel{E}_T$  SIGNAL  
BACKGROUND  $\sim 0.05$  fb





SCENARIO: Include one loop correction to <sup>(11)</sup>  
 KK masses AND KK violating  
 interactions (Macosanu, McMullen + Nandi)

Case 1: Assume decay width due to mass  
 splitting  $\gg$  KK violating grav. widths  
 (happens for  $N=6$ )  
 role of KK violating int:  $\gamma_1^+ \rightarrow \gamma G_m$

Signal:  $g g^+, g^+ z^+, z^+ z^+$   
 $\Rightarrow \gamma_1^+ \gamma_1^+ \rightarrow \underbrace{\gamma \gamma}_{\text{two photons with high } p_T} + \underbrace{G_m G_m}_{\text{missing energy}} + \text{soft particles}$

Backgrounds  $\Rightarrow W+\gamma, W+\text{jets}, Z \rightarrow ee, Z \rightarrow \mu\mu ee$   
 (with misidentified photons)  
 $\Rightarrow$  easily eliminated by  $p_T$  and  $\cancel{E}_T$  cuts.

figs

Run 1 bound:  $\sim 380$  GeV  
 Run 2:  $\sim 520$  GeV  
 LHC:  $\sim 3$  TeV

Opposite case: grav. dominate  $\Rightarrow$  Signal: dijet +  $\cancel{E}_T$   
 ( $N=2$ )

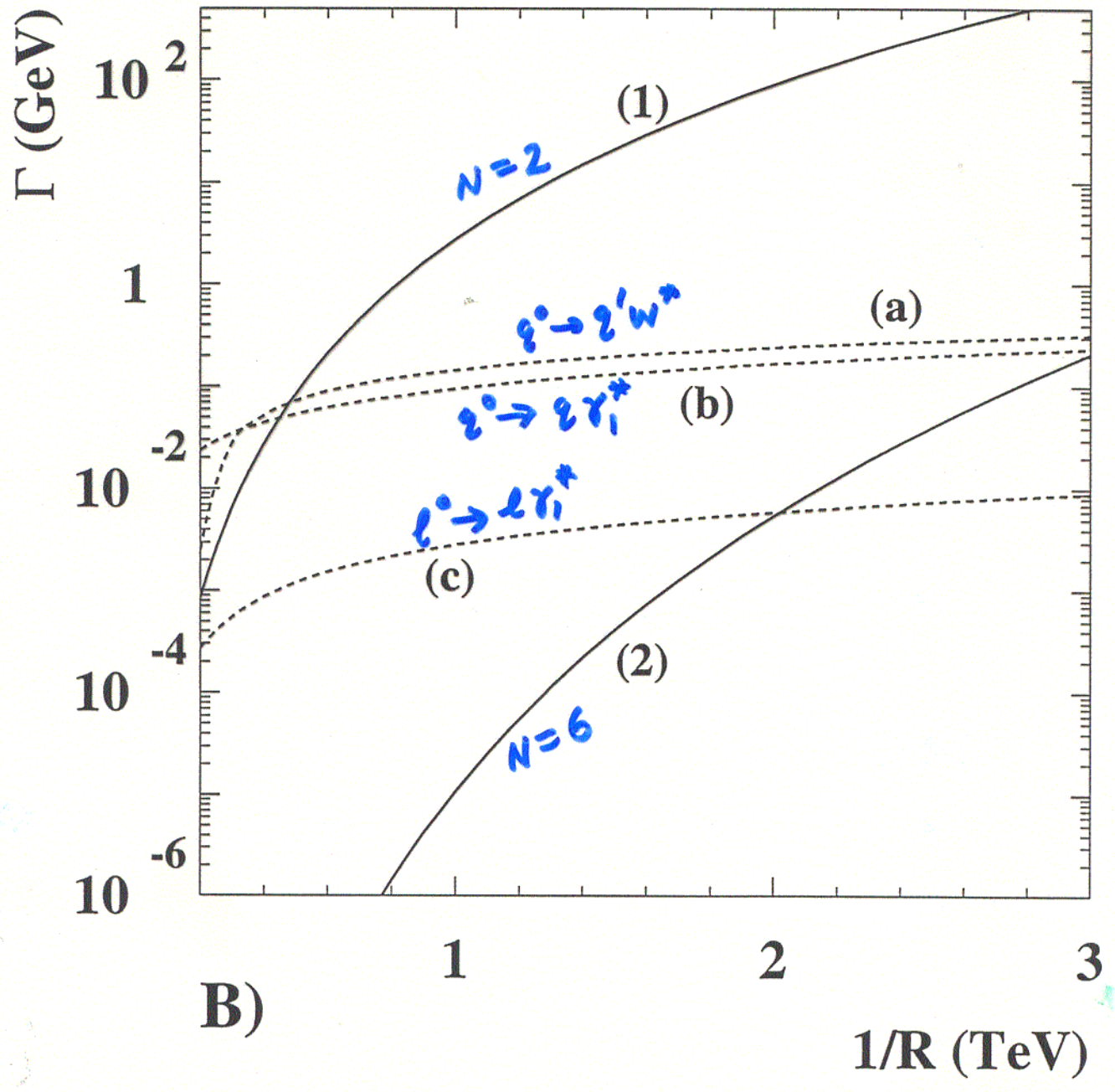
GENERAL CASE: Both decays comparable,

figs

Signal:  $\gamma\gamma + \cancel{E}_T, \text{dijet} + \cancel{E}_T$   
 BR depends on  $N, 1/R$ .

Macesanu, McMullen & Nandi, hep-ph/0207257

# UED, 1 LOOP MASS CORRECTED VS KK VIOLATING DECAYS

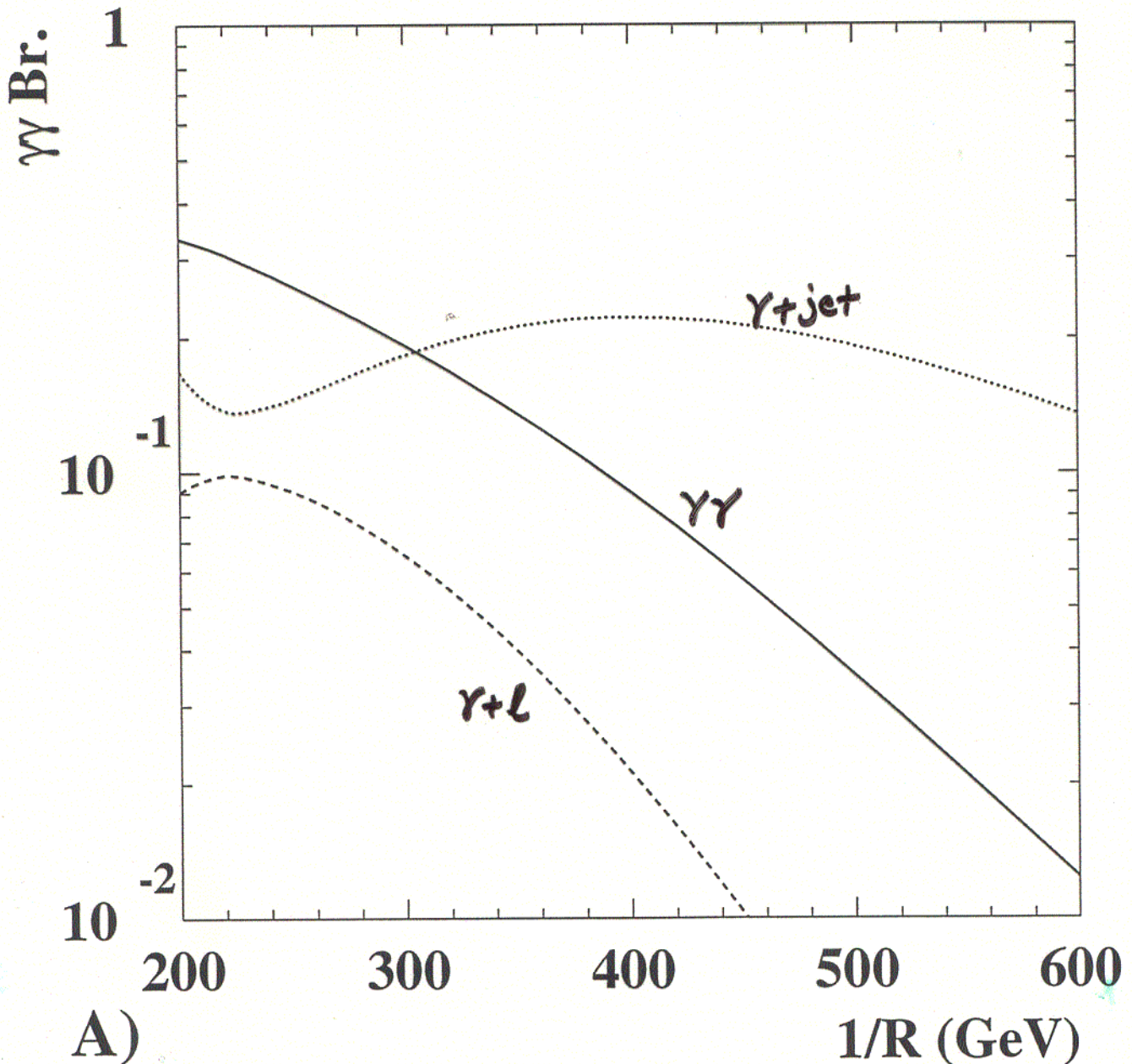


(11c)

Macesanu, McMullen + Nandi, hep-ph/0207269

UED, 1 Loop + KK Violating Decays  
 $\gamma\gamma$  BRANCHING RATIOS

$N=2$



CONCLUSIONS :

: Models with  $\text{TeV}^{-1}$  scale extra dim gives exciting new physics signals at colliders

: Signals

: high  $p_T$  dijets +  ~~$E_T$~~

: high  $p_T$   $\gamma\gamma$  +  ~~$E_T$~~

⋮

: Tevatron Run 2 : 500 GeV to  $\sim 2\text{TeV}$

LHC : 1.5 TeV to  $\sim 7\text{TeV}$

: Allowed scale low enough to see 2<sup>nd</sup> KK excitations at LHC

$\Rightarrow$  existence of extra dim.