

Topics in Meson Spectroscopy

- S-wave mesons below 1 GeV

Kaminski, Lesniak and Rybicki [ABS288] hep-ph/0109268

Van Beveren & Rupp [ABS22] hep-ph/0207022

Furman and Lesniak [ABS284] hep-ph/0203255

- Heavy Quarkonium
- Exotic Hybrids, present & future
- Summary



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S-wave Mesons below 1 GeV

For a recent review see: Close & Tornqvist hep-ph/0204205

The phenomenology of 0^{++} sector reflects the many different components of the Fock space:

$q\bar{q}$

$qq\bar{q}\bar{q}$

Glueballs

MM molecules

Threshold effects

- So it's a real challenge to understand this sector
- 2 steps
 - Analyze the raw data
 - Interpret in terms of underlying models

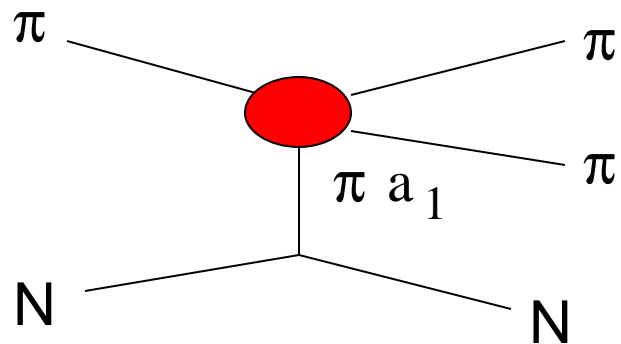


Information on low energy S-wave $\pi\pi$ scattering from

$p^- p \rightarrow p^+ p^- n$ CERN-Munich experiment [NP B75, 189 (1974)]

$p^- p \rightarrow p^0 p^0 n$ CERN-Cracow-Munich [NP B150, 301 (1979); B151, 46 (1979)]

Brookhaven E852 [PR D64, 07003 (2001)]

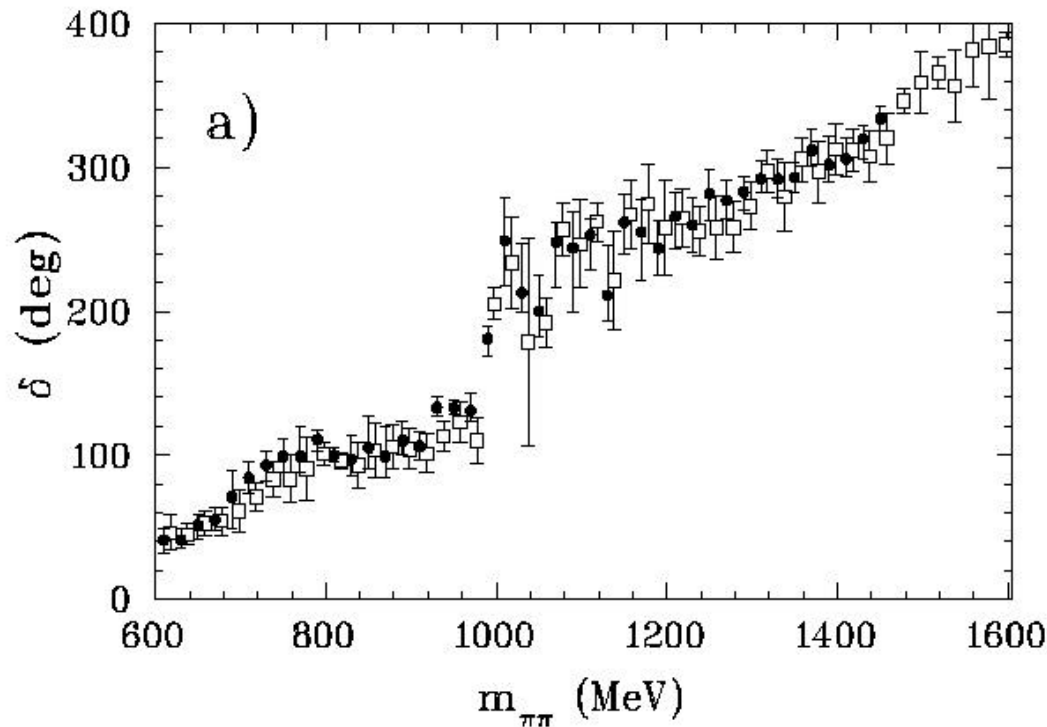


- Experiments provide fewer observables than needed to unambiguously describe partial waves
- Therefore make assumptions in PWA
 - Ignore role of nucleon spin
 - Ignore a_1 exchange amplitude
- Results in ambiguities in extraction of phase shifts from PWA

- Kaminski, Lesniak and Rybicki studied this problem

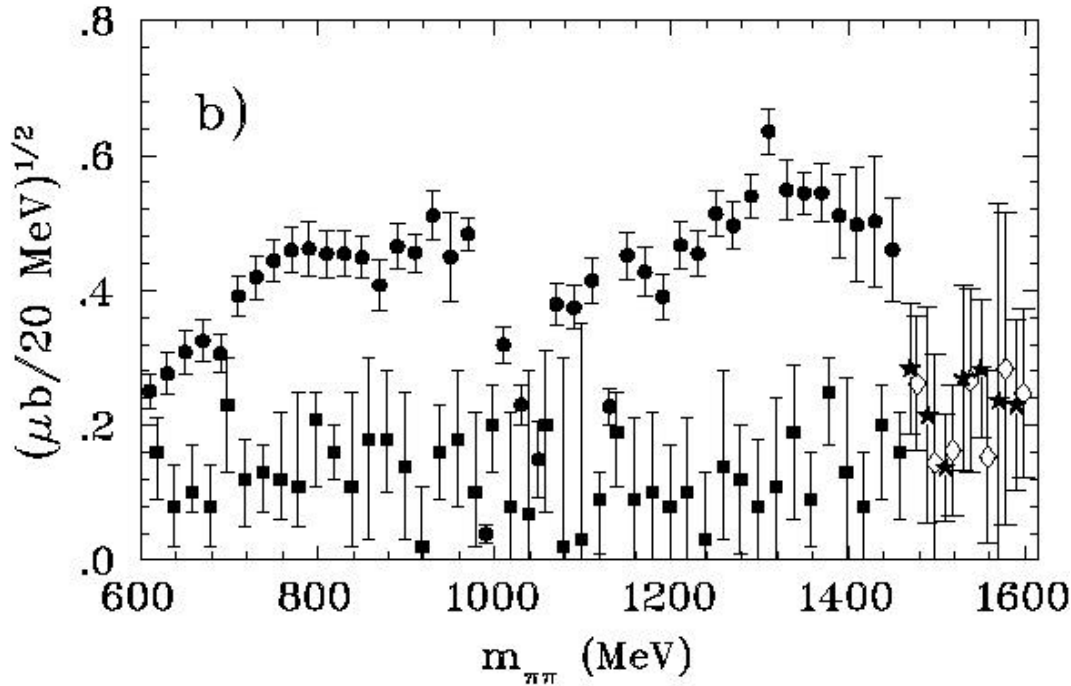
[ABS288: hep-ph/0109268]

- to constrain allowed solutions
 - Apply unitarity
 - Non-observation of inelastic scattering below the $f_0(980)$
- Related $\pi^+\pi^-$ to $\pi^0\pi^0$ amplitudes



- new result
- previous results





Pseudoscalar (circles) and
pseudovector (squares)
amplitudes

Conclude that a_1 contributions are significantly different
from zero (ratio is 0.2-0.3)



- How to describe these phase shifts in terms of true resonances or from the dynamics between KK , $K\pi$, $\pi\pi$?

Van Beveren & Rupp hep-ph/0207022 [ABS22]

- Must look at the meson-meson scattering
- Extract the phase shifts not resonance positions
- Need multichannel approach which includes
resonances (from constituent qq channels)
meson-meson interactions

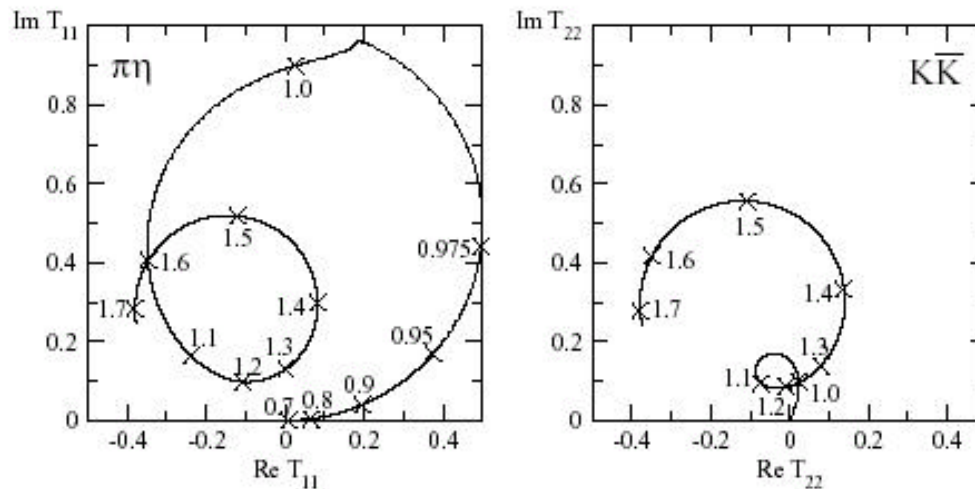
(An interesting and important consequence is that the pole parameters are dependent on the coupling strengths)



Coupled channel model of two a_0 resonances decaying into $\pi\eta$, and $K\bar{K}$ mesons using Lippman-Schwinger equation:

$$\langle \mathbf{p} | T | \mathbf{q} \rangle = \langle \mathbf{p} | V | \mathbf{q} \rangle + \int \frac{d^3 s}{(2\pi)^3} \langle \mathbf{p} | V | \mathbf{s} \rangle \langle \mathbf{s} | G | \mathbf{s} \rangle \langle \mathbf{s} | T | \mathbf{q} \rangle$$

where T, V , and G are 2×2 matrices



Find significantly different widths in the two channels in agreement with E852 and Crystal Barrel (also used to describe $I=0$ sector)

- These multichannel approaches

Furman & Lesniak hep-ph/0203255 [ABS284]

Van Beveren & Rupp hep-ph/0207022 [ABS22]

(and others)

obtain a good description of the low lying S-wave spectrum

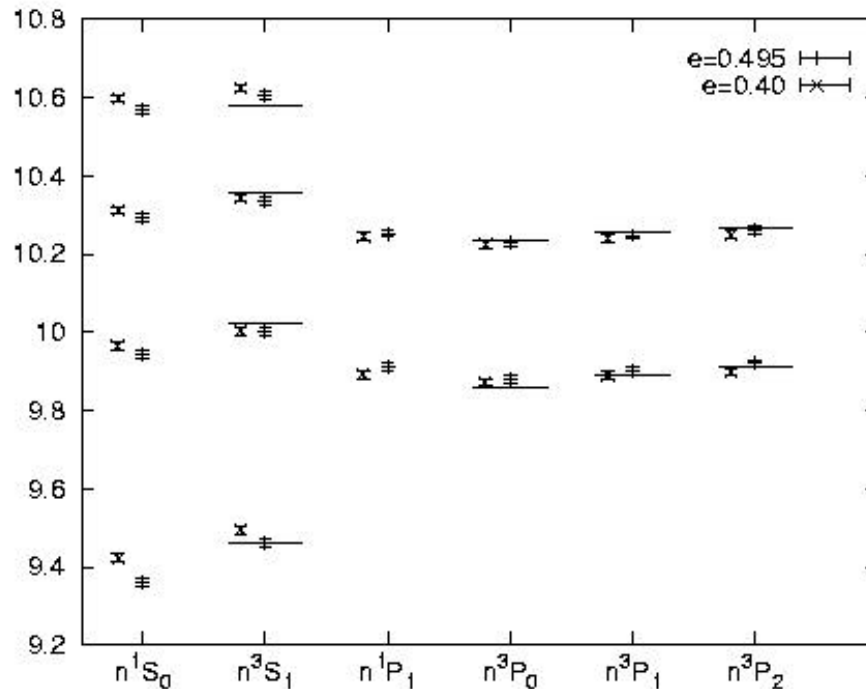
My impression is that workers in the field have come to a consensus on the general description if not the details.



Heavy Quarkonium

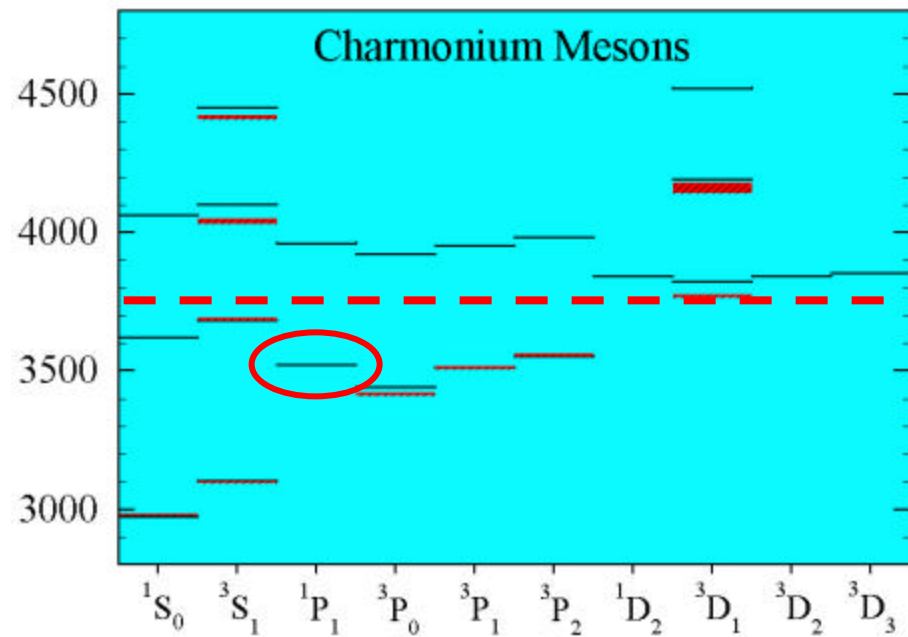
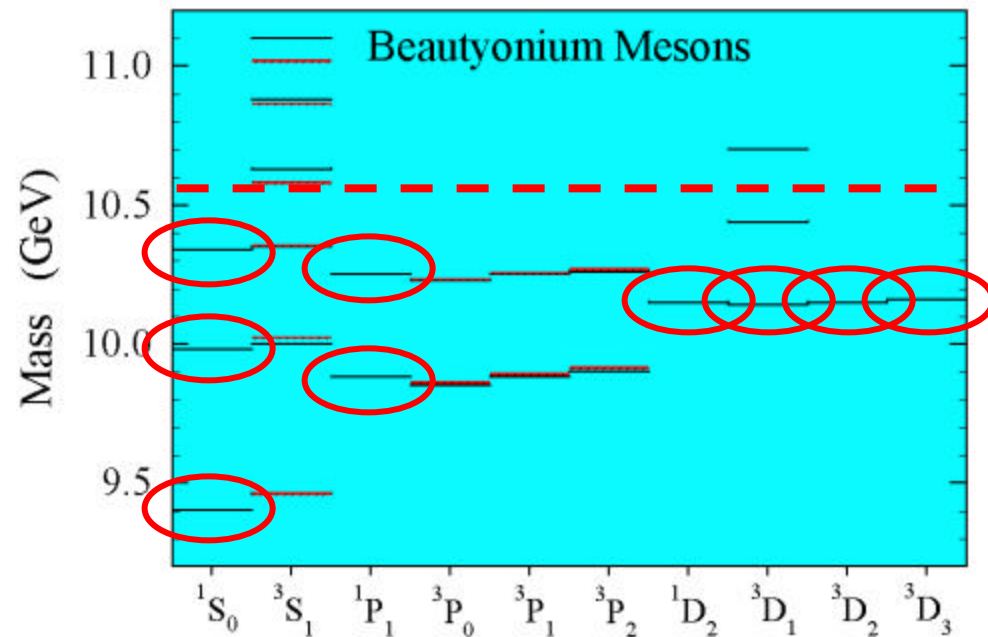
- Recent interest due to
 - CLEO/CESR run on $Y(3S)$
 - Belle observation of $\eta_c(2S)$ in B decay
- Lattice QCD starting to make quantitative predictions for the hadron mass spectrum
- Need to test Lattice calculations against experiment

$b\bar{b}$ Mass

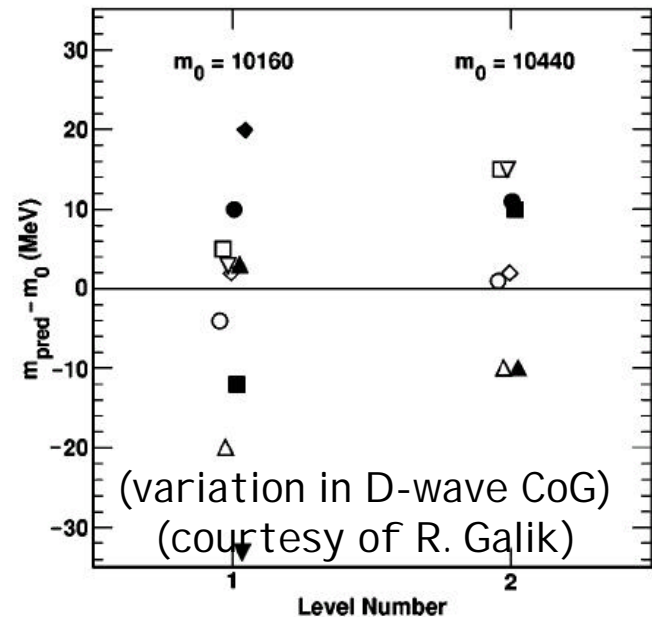


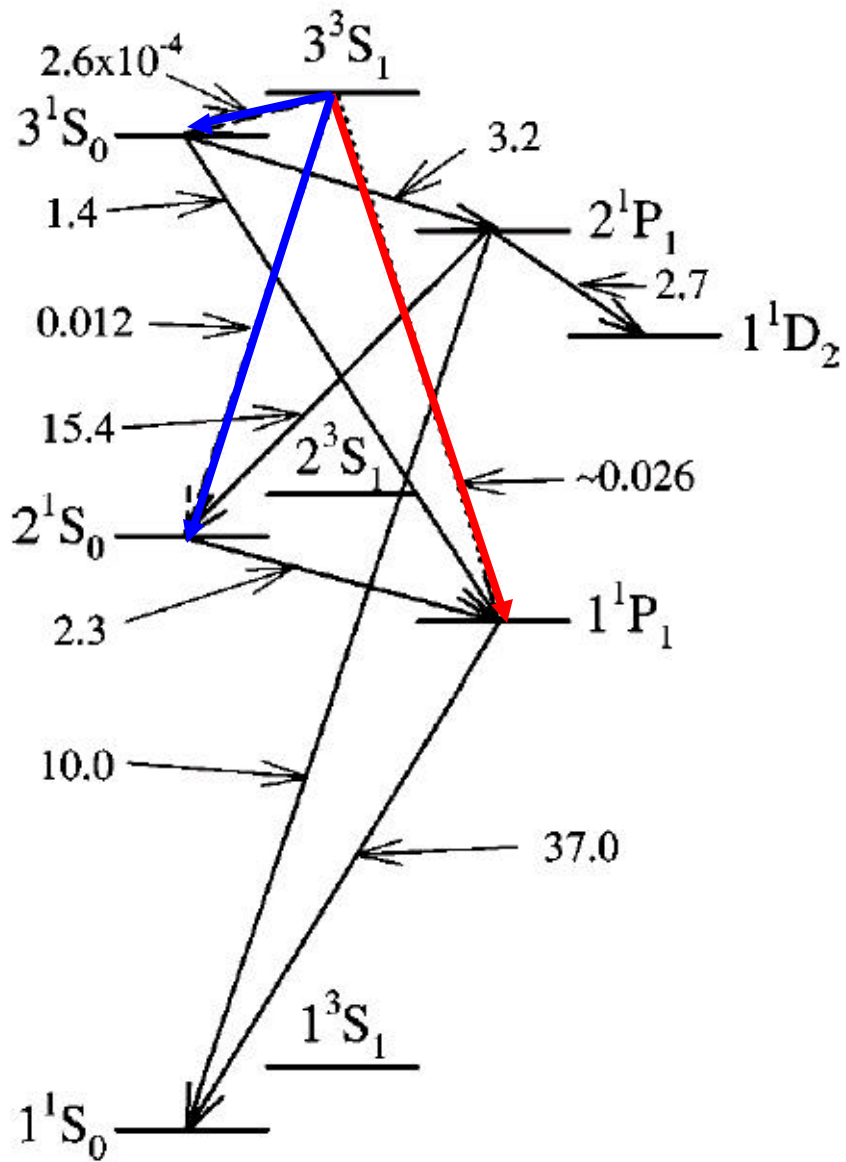
Bali, Schilling and Wachter
Hep-ph/9611226





- Numerous states below threshold
- Spin triplet states have been observed
- Few spin singlet states have been seen
- Wide variation in splittings
- Their observation would test the various Calculations
- Expect many of these states to be found in
 - The recent CESR/CLEO run
 - B-decays at B-factories
 - At future CLEO-c/CESR-c





- CESR/CLEO has just completed a high statistics run at the Y(3S)
- Expect very rich spectroscopy
- Estimate the radiative widths and BR using quark model

Production of the $h_b(nS)$ states

S.G + J. Rosner, Phys Rev D64, 074011 (2001)

Proceeds via magnetic dipole (M1) transitions:

$$Y(nS) \rightarrow \eta(n'S) + \gamma$$

$$\Gamma(^3S_1 \rightarrow ^1S_0 + \mathbf{g}) = \frac{4}{3} \mathbf{a} \frac{e_Q^2}{m_Q^2} \left| \langle f | j_0(kr/2) | i \rangle \right|^2 \mathbf{w}^3$$

- Hindered transitions have large phase space
- Relativistic corrections resulting in differences in 3S_1 and 1S_0 wavefunctions due to hyperfine interaction



	Transition	BR (10^{-4})
Y(3S)		
($\Gamma_{\text{tot}}=52.5$ keV)	$\rightarrow 3^1S_0$	0.10
	$\rightarrow 2^1S_0$	4.7
	$\rightarrow 1^1S_0$	25
Y(2S)	$\rightarrow 2^1S_0$	0.21
($\Gamma_{\text{tot}}=44$ keV)	$\rightarrow 1^1S_0$	13
Y(1S)	$\rightarrow 1^1S_0$	2.2
($\Gamma_{\text{tot}}=26.3$ keV)		

- Expect substantial rate to produce η_b 's
- Also $Y(3S) \rightarrow h_b(1P_1) \pi\pi \rightarrow \eta_b + \gamma + \pi\pi$
BR=0.1-1% BR = 50%

[Kuang & Yan PRD24, 2874 (1981); Voloshin Yad Fiz 43, 1571 (1986)]



Production of the singlet P-wave states

S.G + J. Rosner, PR D66 in press

Two interesting cascades:

M1

E1

E1

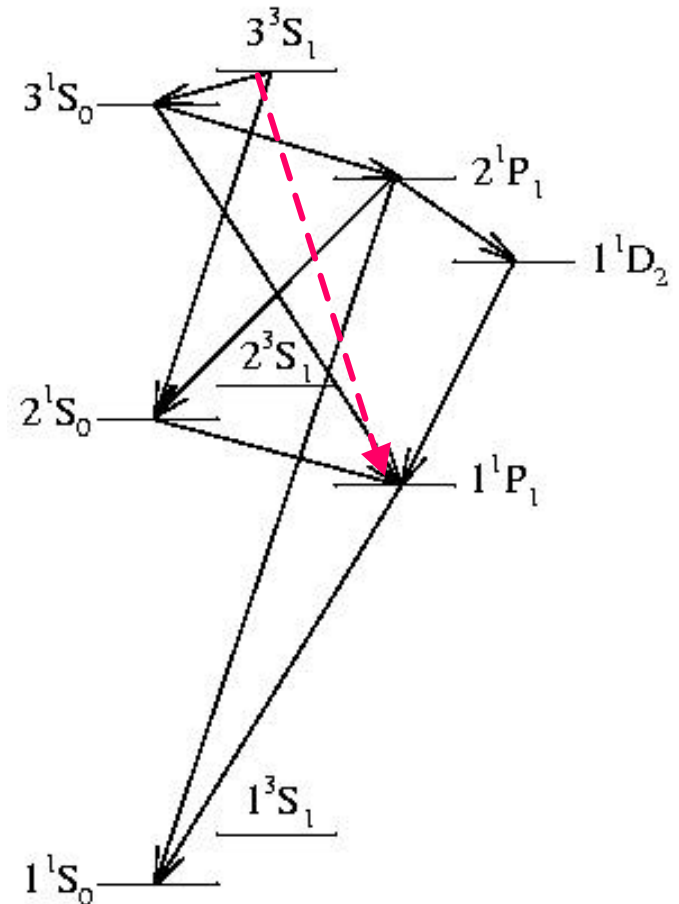
$$Y(3S) \rightarrow \eta_b(2S) + \gamma \rightarrow h_b + \gamma\gamma \rightarrow \eta_b + \gamma\gamma\gamma$$

$$\psi(2S) \rightarrow \eta_c(2S) + \gamma \rightarrow h_c + \gamma\gamma \rightarrow \eta_c + \gamma\gamma\gamma$$

$$Y(3S) \rightarrow h_b + \pi \rightarrow \eta_b + \gamma + \pi$$

$$\psi(2S) \rightarrow h_c + \pi \rightarrow \eta_c + \gamma + \pi$$

Need branching ratios and hence partial widths



$$\Gamma[\mathbf{h}(2^1S_0) \rightarrow h_b(1^1P_1) + \mathbf{g}] = \frac{4}{3} \mathbf{a} e_Q^2 \left| \langle 1^1P_1 | r | 1^1S_0 \rangle \right|^2 \mathbf{w}^3 = 2.3 \text{ keV}$$

$$\Gamma[h_b(1^1P_1) \rightarrow \mathbf{h}_b(1^1S_0) + \mathbf{g}] = \frac{4}{9} \mathbf{a} e_Q^2 \left| \langle 1^1S_0 | r | 1^1P_1 \rangle \right|^2 \mathbf{w}^3 = 37 \text{ keV}$$

$$\Gamma[\mathbf{h}_b(2^1S_0) \rightarrow gg] = \frac{27\mathbf{p}}{5(\mathbf{p}^2 - 9)\mathbf{a}_s} \times \Gamma[\mathbf{Y}(2^3S_1) \rightarrow ggg] = 4.1 \pm 0.7 \text{ MeV}$$

$$\text{BR}(3^3S_1 \gamma \rightarrow 2^1S_0 \gamma) = 4.7 \times 10^{-4} \quad \text{and} \quad \text{BR}(2^1S_0 \gamma \rightarrow 1^1P_1 \gamma) = 5.7 \times 10^{-5}$$

$$\text{BR}[\mathbf{Y}(3S) \rightarrow 2^1S_0 \gamma \rightarrow 1^1P_1 \gamma] = 2.6 \times 10^{-7} \Rightarrow 0.3 \text{ events}/10^6 \mathbf{Y}(3S)'s$$

Similarly

$$\text{BR}[\psi(2S) \rightarrow 2^1S_0 \gamma \rightarrow 1^1P_1 \gamma] = 10^{-6} \Rightarrow 1 \text{ event}/10^6 \mathbf{Y}(3S)'s$$

(A challenge for the experimentalists!)



A more promising approach utilizes:

$$\text{BR}[Y(3S) \rightarrow \pi 1^1P_1] = 0.1\%$$

$$\Gamma[h_b(1^1P_1) \rightarrow \mathbf{h}_b(1^1S_0) + \mathbf{g}] = \frac{4}{9} \mathbf{a}e_Q^2 \left| \langle 1^1S_0 | r | 1^1P_1 \rangle \right|^2 \mathbf{w}^3 = 37 \text{ keV}$$

$$\Gamma[h_b(1^1P_1) \rightarrow ggg] = \frac{5}{2n_f} \Gamma[\mathbf{c}_{b1}(1^3P_1) \rightarrow q\bar{q}g] = 50.8 \text{ keV}$$

$$\text{BR}[Y(3S) \rightarrow \pi 1^1P_1 \rightarrow 1^1S_0 \gamma] = 4 \times 10^{-4} \Rightarrow 400 \text{ events}/10^6 Y(3S)'s$$

Similarly

$$\text{BR}[\psi(2S) \rightarrow \pi 1^1P_1 \rightarrow 1^1S_0 \gamma] = 3.8 \times 10^{-4} \Rightarrow \sim 400 \text{ event}/10^6 \psi(2S)'s$$

Expect ~400 events!



Production of the D -wave states

- By direct scans in e^+e^- to produce 3D_1
- In e.m. cascades: $Y(3S) \rightarrow \gamma \chi'_b \rightarrow \gamma\gamma ^3D_J$
- Some 4 γ cascades with observable # of events/ 10^6 $Y(3S)$'s:

Cascade	Events
$3^3S_1 \rightarrow 2^3P_2 \rightarrow 1^3D_3 \rightarrow 1^3P_2 \rightarrow 1^3S_1$	7.8
$3^3S_1 \rightarrow 2^3P_2 \rightarrow 1^3D_2 \rightarrow 1^3P_1 \rightarrow 1^3S_1$	2.7
$3^3S_1 \rightarrow 2^3P_1 \rightarrow 1^3D_2 \rightarrow 1^3P_1 \rightarrow 1^3S_1$	20
$3^3S_1 \rightarrow 2^3P_1 \rightarrow 1^3D_1 \rightarrow 1^3P_1 \rightarrow 1^3S_1$	3.3

S.G + J. Rosner, Phys Rev D64, 097501 (2001)

Expect ~ 38 events / 10^6 $Y(3S)$ via 3D_J

- The e^+e^- final states leads to less background
- $\mu^+\mu^-$ final states also contribute if μ 's are identified



In the CESR run just completed expect to see evidence for the

$$2^1S_0, 1^1S_0, 1^1P_1, 1^3D_2$$

And maybe the

$$3^1S_0, 1^3D_1 \text{ and } 1^3D_3$$

Would represent a significant increase in our knowledge of quarkonium and provide an important benchmark against which to measure the results of lattice QCD

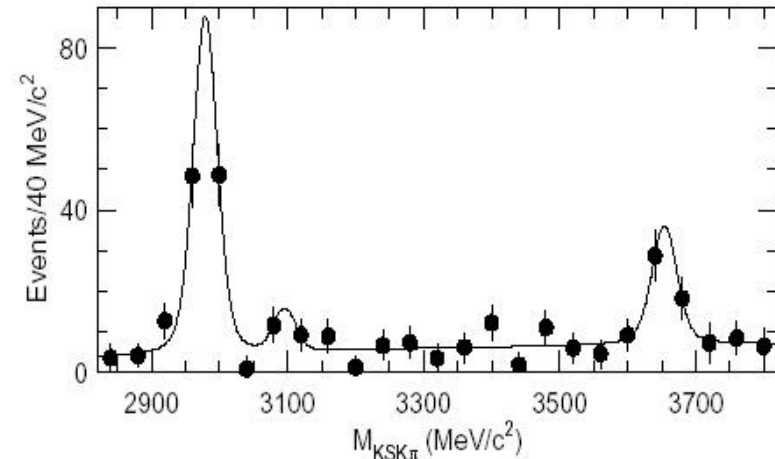


Charmonium in B decays

Recent observation by Belle of $\eta_c(2S)$ in: $B \rightarrow \eta_c(2S) K \rightarrow KK_S K \pi^+$

$M = 3654 \pm 6$ (stat) ± 8 (sys) MeV
 $\Gamma < 55$ MeV (90% C.L.)

[hep-ex/0206002]



Belle had previously reported the observation of

$$B^+ \rightarrow \chi_{c0} K^+$$

PRL 88, 031802 (2002)

$$B \rightarrow \chi_{c2} X$$

hep-ex/0202038

And $B \rightarrow \chi_{c1} K$ has been observed by both BaBar and Belle

Suzuki proposed to search for the h_c in $B \rightarrow h_c X$ [hep-ph/0204043]

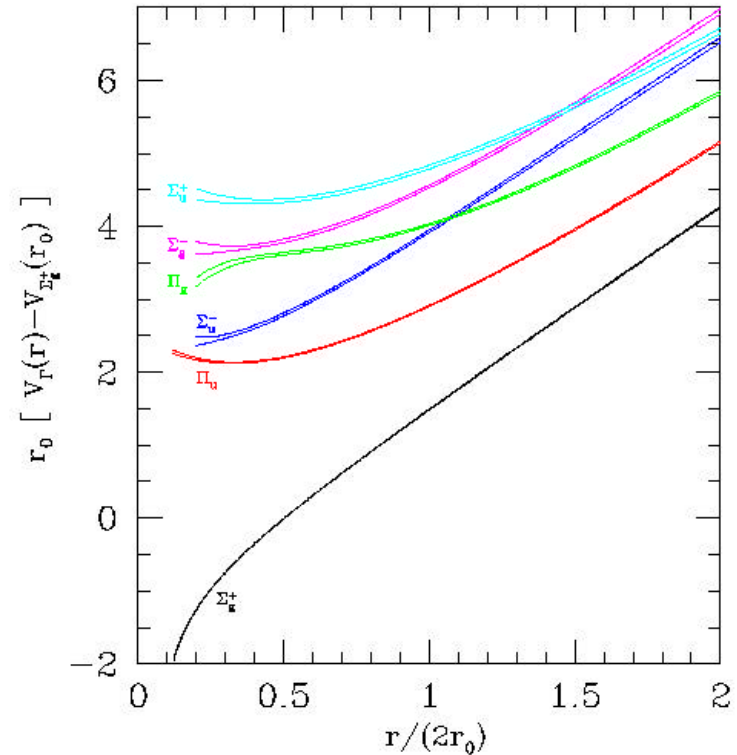
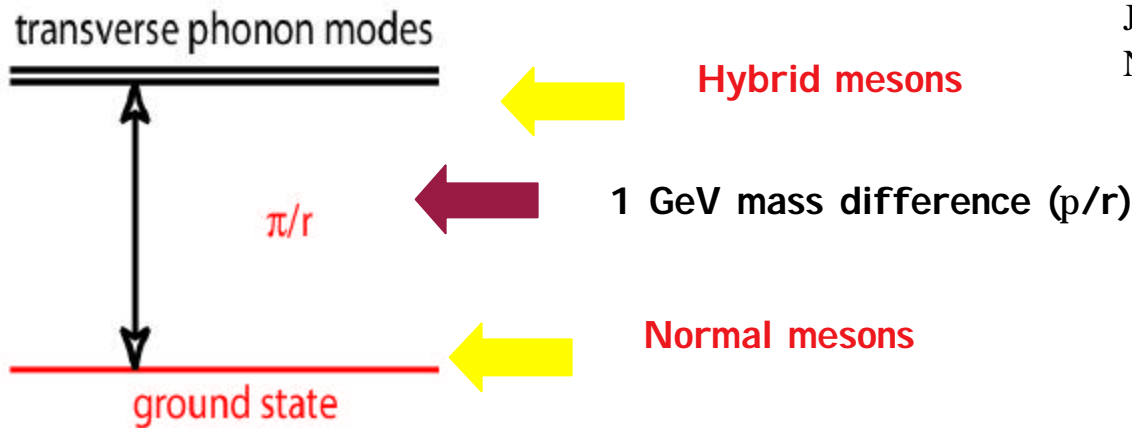
[see also Eichten Lane & Quigg hep-ph/0206018; Gu hep-ph/0206002]

Other related work by Hao Liu & Chao hep-ph/0206226



Exotic Hybrids

- Quarks move in effective potentials of adiabatically varying state of flux tubes
- Lowest excited adiabatic surface corresponds to transverse excitations



Juge, Kuti, and Morningstar,
Nucl. Phys. (Proc. Suppl.) **63A-C**, 326 (1998)

*Lowest mass hybrids
at ~ 1.9 GeV*

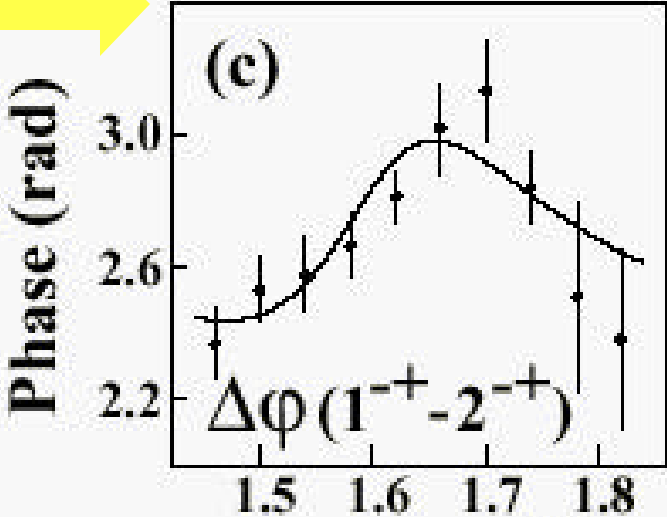
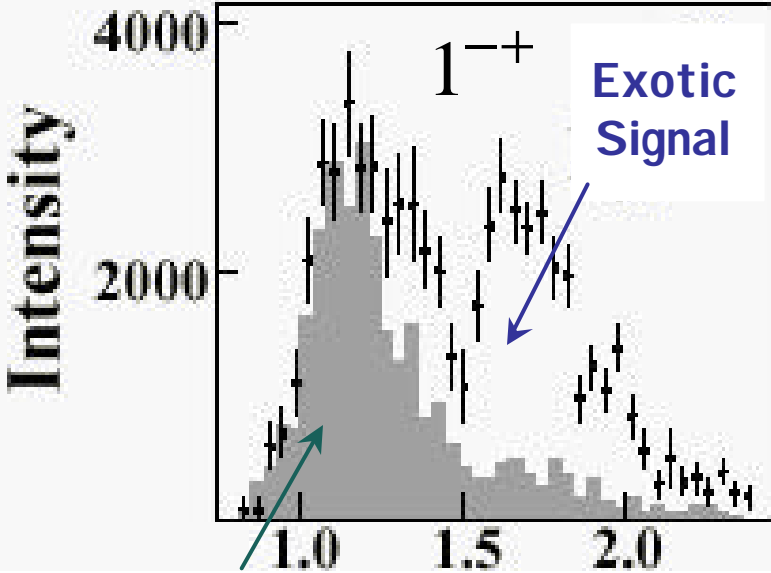
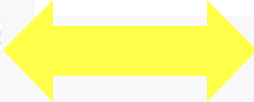
Doubly degenerate:

$$J^{PC} = 0^{+-} \ 0^{-+} \ 1^{+-} \ 1^{-+} \\ 2^{+-} \ 2^{-+} \ 1^{++} \ 1^{--}$$



An Exotic Signal in E852

Correlation of
Phase
&
Intensity



Leakage
From
Non-exotic Wave
due to imperfectly
understood acceptance

$M(\pi^+\pi^-\pi^-)$ [GeV / c^2]

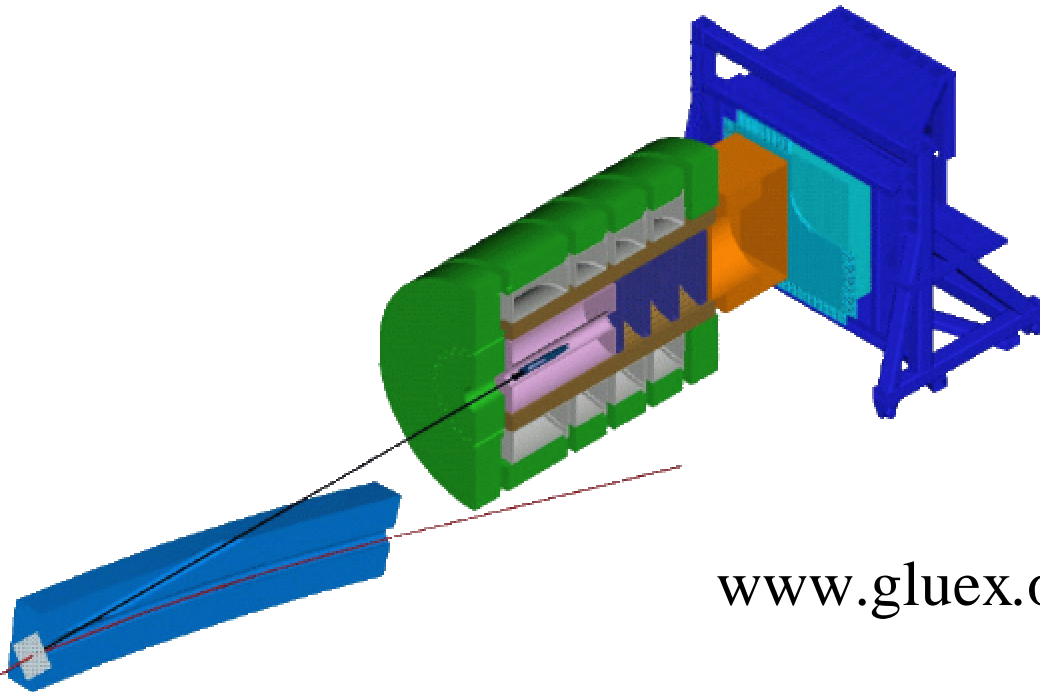


What about the future?



Hall D at Jefferson Lab

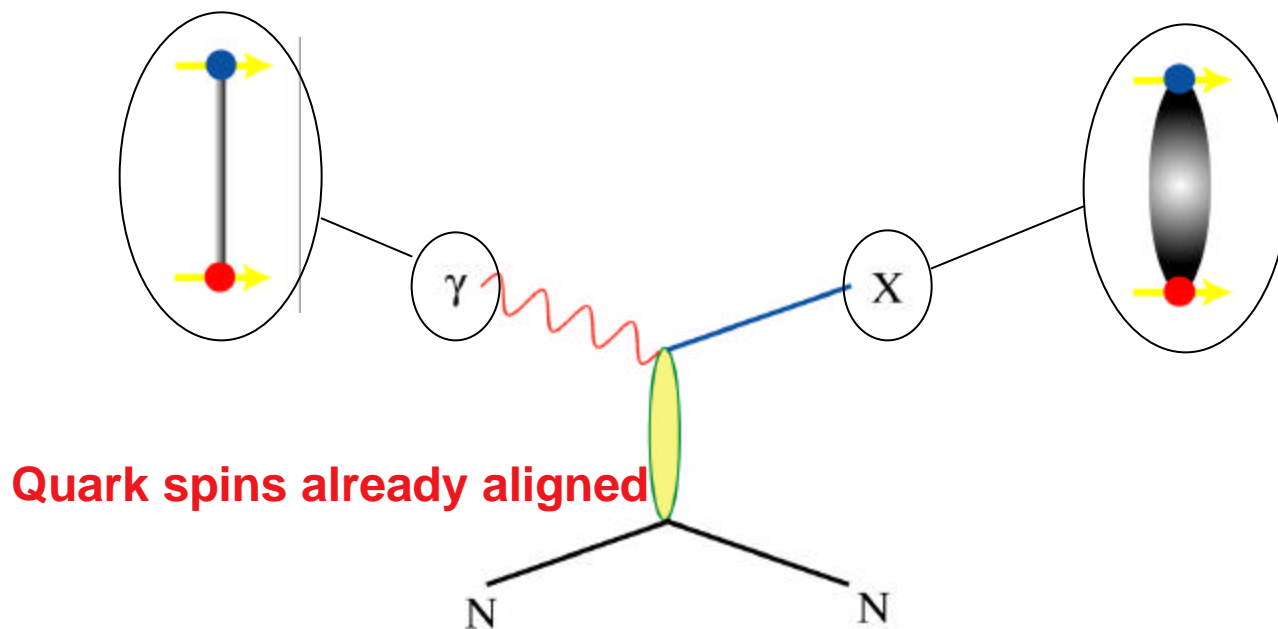
"Searching for Exotic Gluonic Excitations"



www.gluex.org



Photoproduction:



Quark spins already aligned

Production of exotic hybrids favored.
Almost no data available

Optimum photon energy is about 9 GeV

Proposal to upgrade CEBAF to 12 GeV
Produce photons through coherent bremsstrahlung

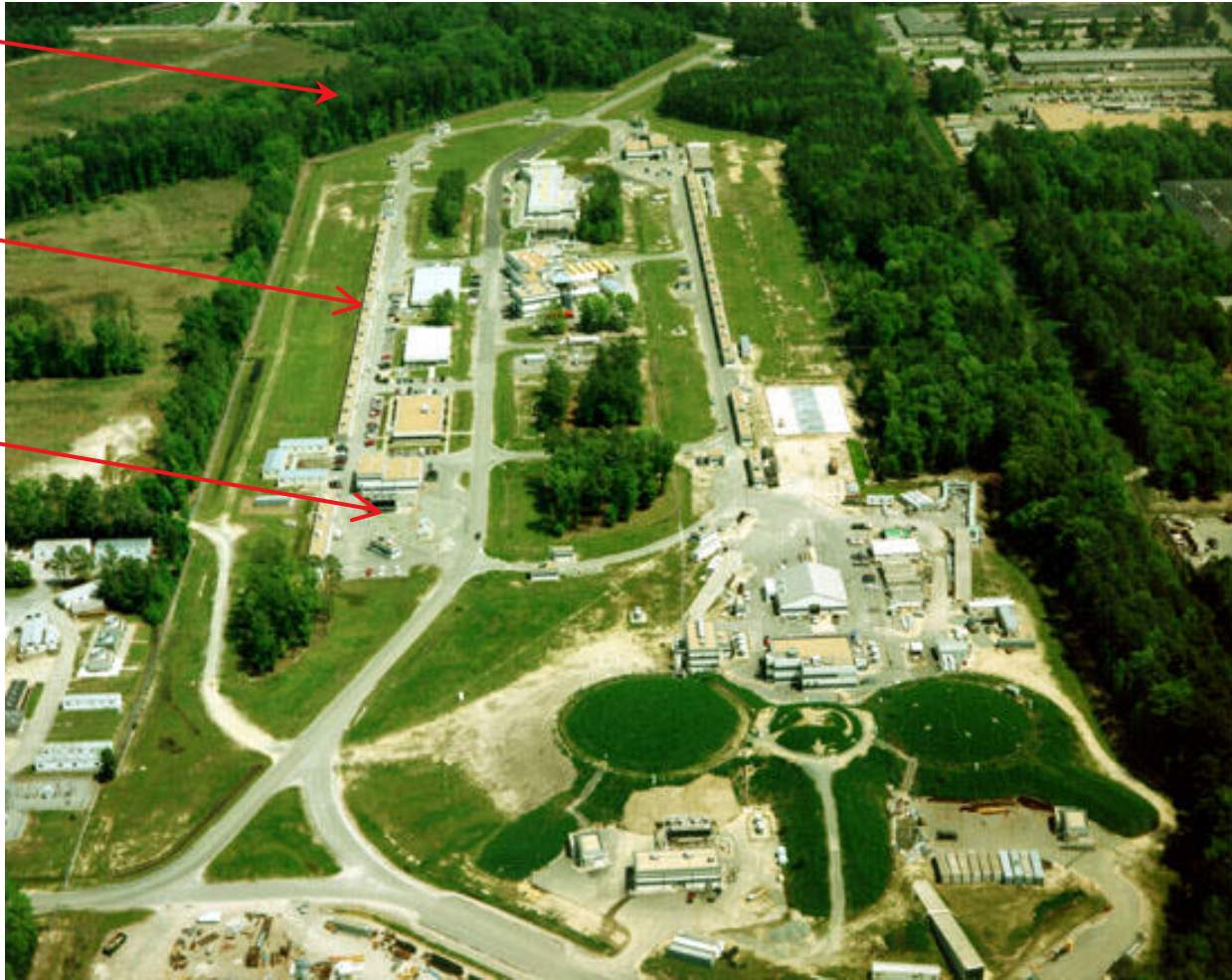


The JLab Accelerator Complex

Hall D will
be located
here

Linac

Arc



Construction start – 2004 **Physics - 2008**



Summary

- S-wave $\pi\pi$ scattering far from a closed book but definite progress in our understanding
- Expect great progress in heavy quarkonium spectroscopy
- Confirmation and mapping out of hybrid-meson spectrum is one of the most important qualitative question facing hadron spectroscopy
- Theory and Experiment go hand in hand to fully understand **Soft QCD**

