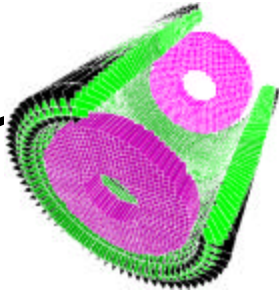


**Bottomonium studies  
via  $\Upsilon(3S)$  decays**

**First observation of  $\Upsilon(1D)$**

*Tomasz Skwarnicki*

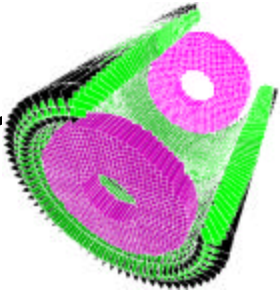
*Representing  
the CLEO collaboration*



# Onia

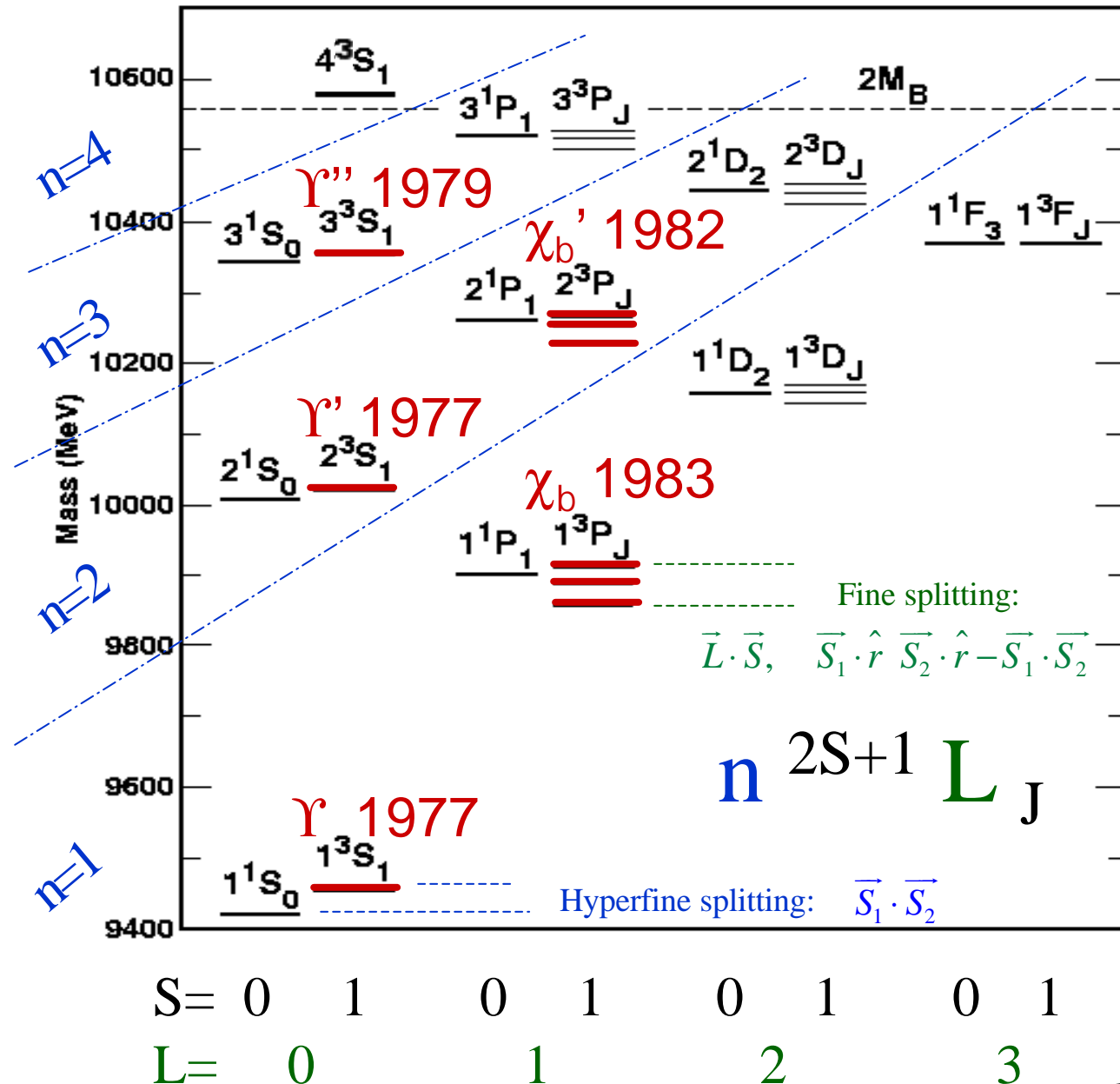
FORCES		System	$(v/c)^2$	Ground triplet state $1^3S_1$		Number of states below dissociation energy		
binding	decay			Name	$\Gamma$ (MeV)	$n^3S_1$	all	
<b>POSITRONIUM</b>								
EM	EM	$e^+e^-$	$\sim 0.0$	Ortho-	$5 \cdot 10^{-15}$	2	8	
<b>QUARKONIUM</b>								
S T R O N G	S T R O N G		$u\bar{u}, d\bar{d}$	$\sim 1.0$	$\rho$	150.00	0	0
			$s\bar{s}$	$\sim 0.8$	$\phi$	4.40	"1"	"2"
		E M	$c\bar{c}$	$\sim 0.25$	$\psi$	0.09	2	8
			$b\bar{b}$	$\sim 0.08$	$\Upsilon$	0.05	3	30
	weak	$t\bar{t}$	$< 0.01$			3000.00	0	0

- Heavy quarkonia hold a promise of playing a similar role for QCD as positronium did for QED
  - Upsilon states are the most non-relativistic (i.e. simplest) states among all long-lived quarkonia states
  - The Upsilon system also has the largest number of stable states
- ⇒ Upsilon states play a special role in probing the strong interactions (tests of lattice QCD, potential models)

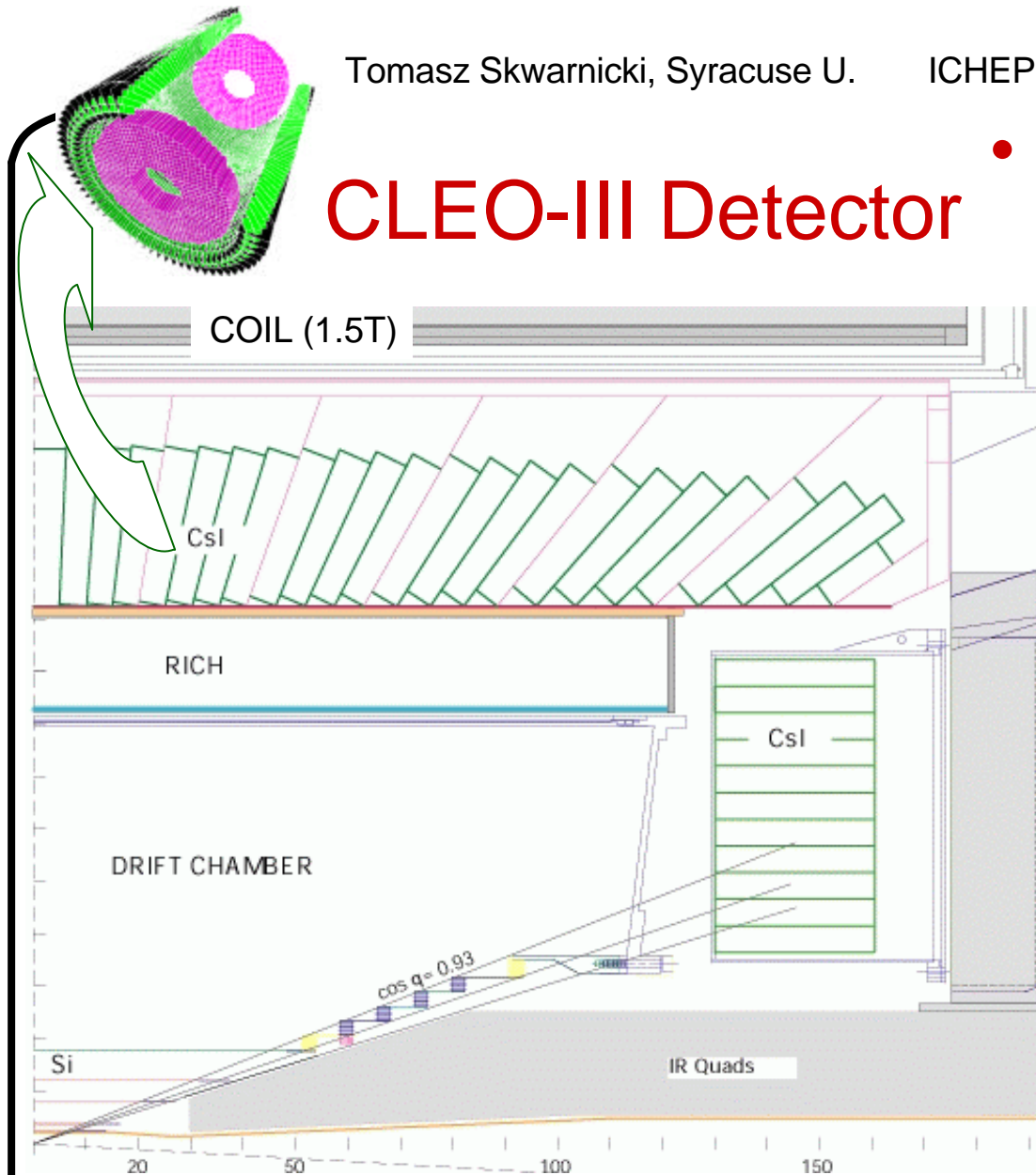


# Upsilon States

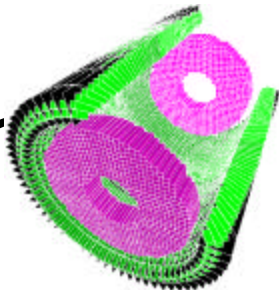
- Only 9 out of 30 narrow states observed so far
- No spin-singlet states observed
- No new states observed in 19 years



# CLEO-III Detector

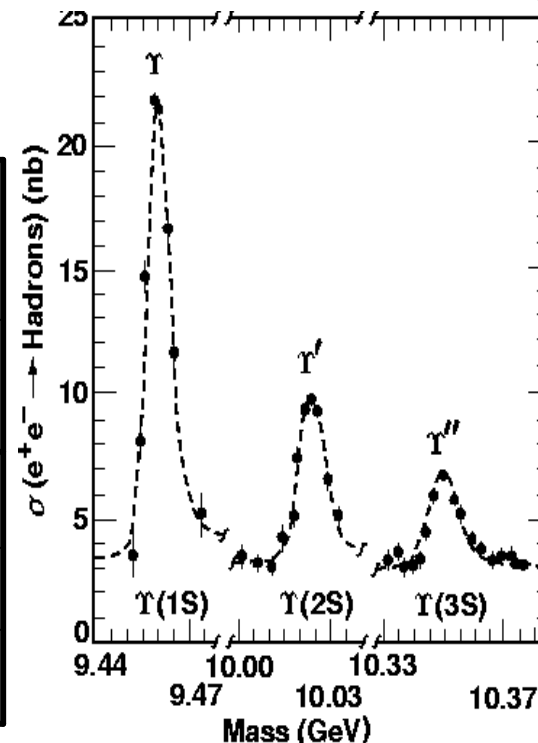


- **EM calorimeter** - Essential for this work
  - ~8000 CsI(Tl) crystals + photo-diodes
  - First crystal calorimeter in magnetic field
  - **Changes since CLEO-II:**
    - **Low-mass DR endplate!**
    - Re-stacked endcaps, moved away
    - New readout electronics
    - Some light loss due to the deteriorating glue used to attach the photo-diodes



# CLEO-III $\Upsilon$ Data

Reso- nance	CLEO-III Integrated Luminosity ( $\text{pb}^{-1}$ )			Number of resonance decays ( $10^6$ )		
	ON	OFF	Scan	CLEOIII	CLEOII	CUSB (C.Ball)
$\Upsilon(3S)$	1150	128	100	4.7	0.46	1.3
$\Upsilon(2S)$	550	150	50	~3.6	0.49	(0.19)
$\Upsilon(1S)$	1230	200	100	~29.0	1.9	(0.48)

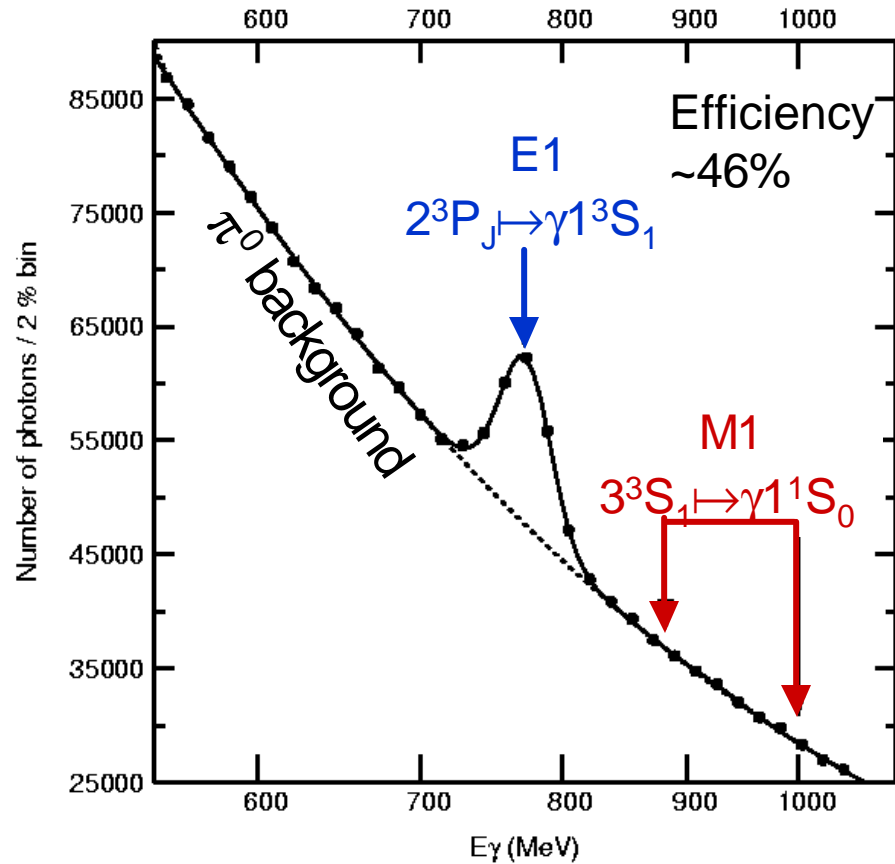
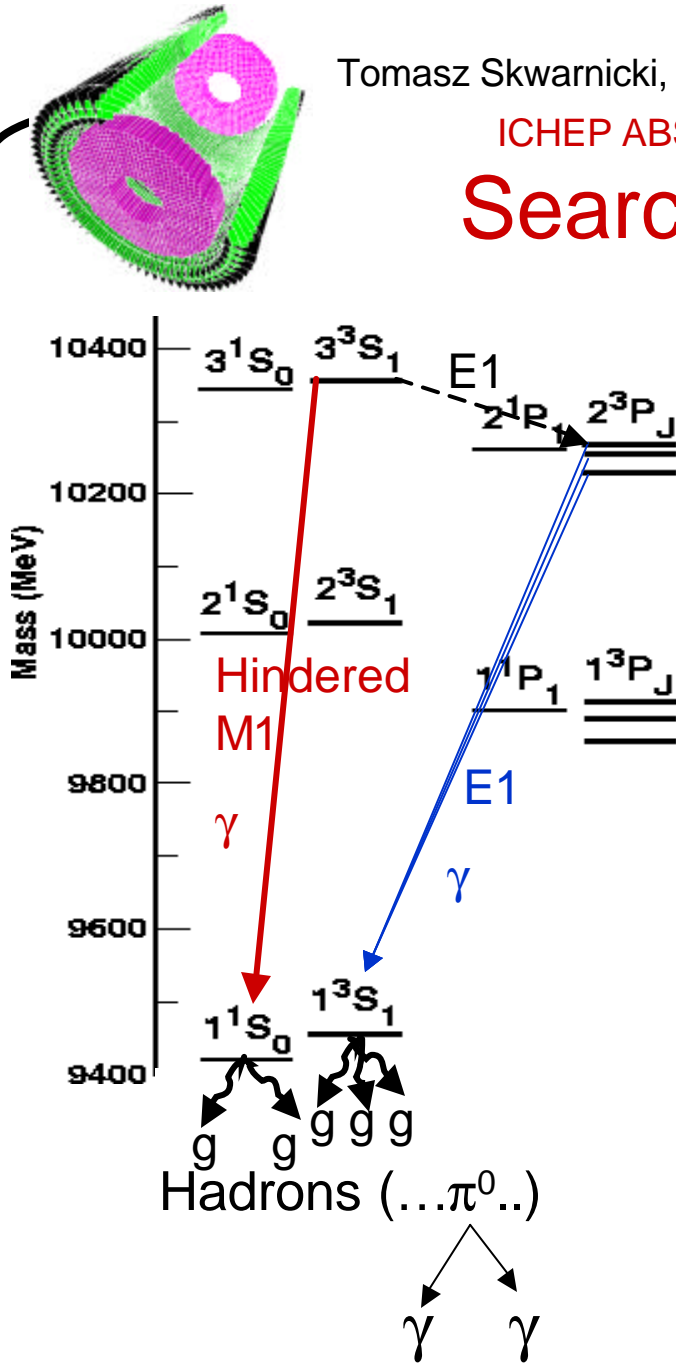


- About 10-fold increase over the CLEO-II statistics
- About 4-times more  $\Upsilon(3S)$  data than analyzed by CUSB + 2.5-4.5 times higher efficiency for the final states analyzed here
- The  $\Upsilon(3S)$  data already processed off-line
- We will take more  $\Upsilon$  data before lowering the beam energy to the charm threshold region next year ( $\rightarrow$  CLEO-c)

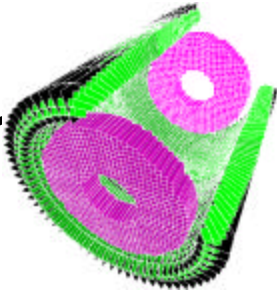
ICHEP ABS947, CLEO CONF 02-05

# Search for $\eta_b(1^1S_0)$

**Inclusive g spectrum  
in multi-hadronic events**

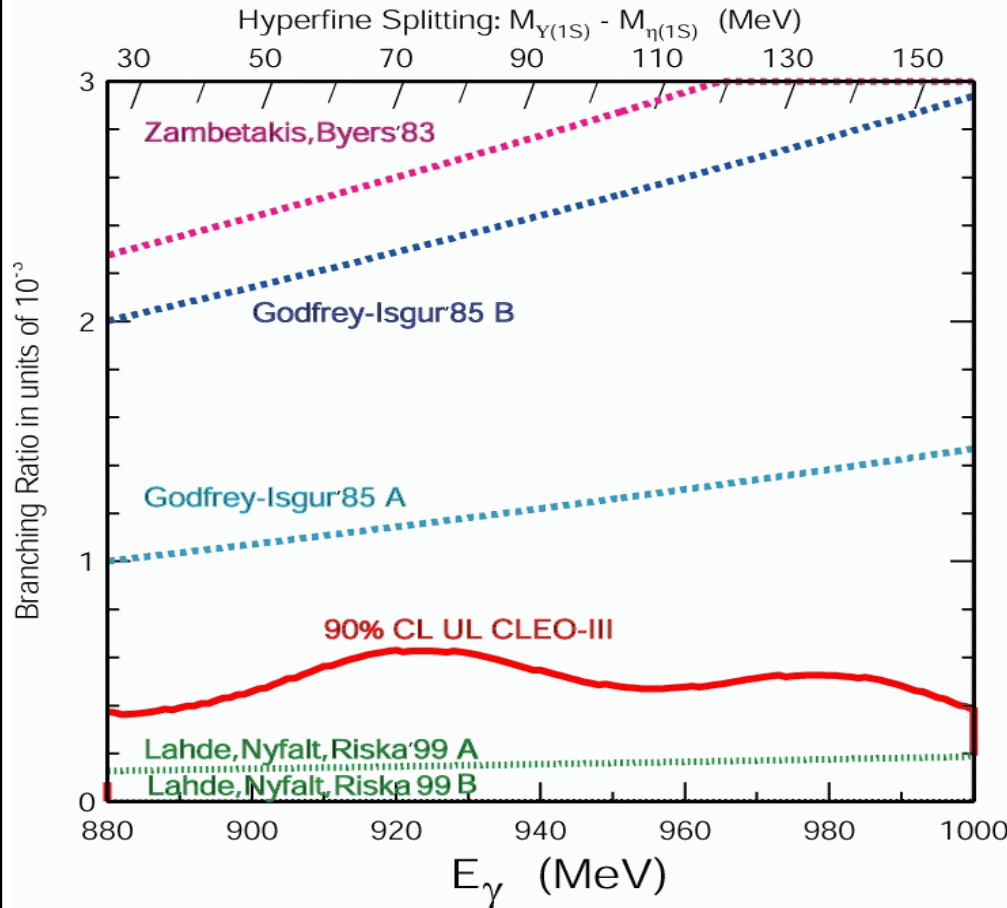


• **No signal found**



# Search for $\eta_b(1^1S_0)$

- Test potential model predictions for  $\Gamma_{M1}$



Models from the compilation by Godfrey&Rosner PR D64, 074011 (2001) [scaled here by phase-space]

$$\Gamma_{M1} \propto \frac{e_b^2}{m_b^2} \left| \langle n_f L | n_i L \rangle \right|^2 E_g^3$$

DIRECT  $n_i = n_f$

$$\langle n_f L | n_i L \rangle = 1 \quad E_g^3 \text{ -tiny}$$

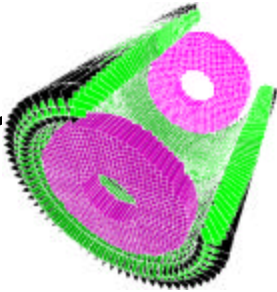
hopeless for  $b\bar{b}$

HINDERED  $n_i \neq n_f$

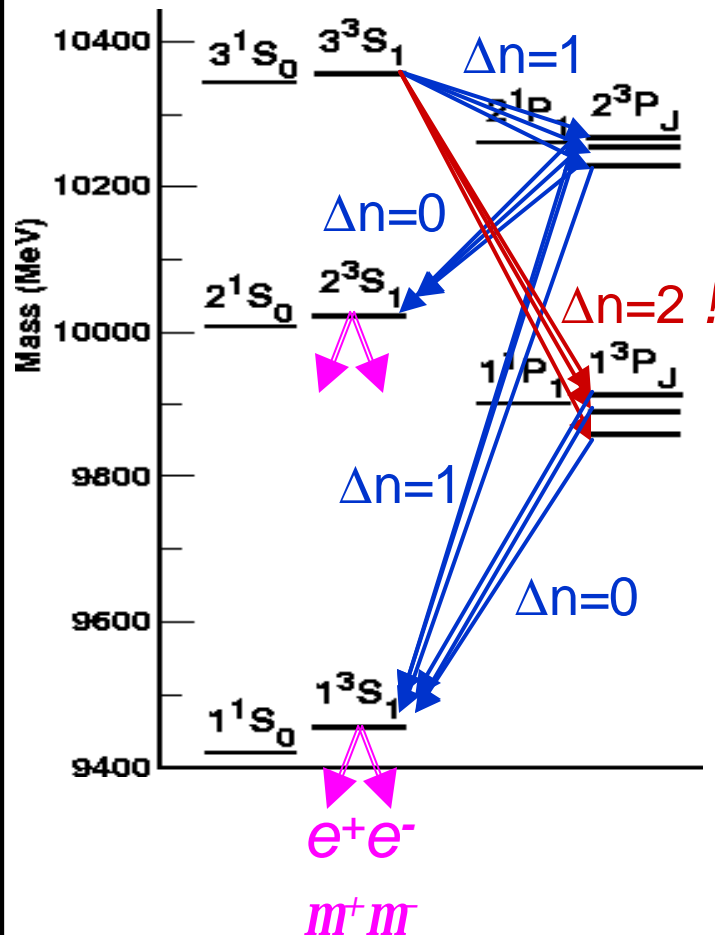
$$\langle n_f L | n_i L \rangle \approx 0 \quad E_g^3 \text{ -large}$$

difficult to predict

- Most of the calculations are ruled out!



# Exclusive $2\gamma$ -cascades



- $\gamma\gamma l^+l^-$  final states
- No  $\pi^0$  backgrounds from gluonic  $b\bar{b}$  annihilation
- Low product branching ratio (a few  $10^{-4}$ )
- Sensitivity to hadronic widths of triplet P-states:

$$B(P \mapsto \gamma S) = \Gamma_{E1} / (\Gamma_{E1} + \Gamma_{had})$$

- $3S \mapsto \gamma 1P$  is a  $\Delta n=2$  transition (rare)

$$\Gamma_{E1} \propto e_b^2 \left| \langle n_f L_f | r | n_i L_i \rangle \right|^2 E_g^3$$

much larger than  $\Gamma_{M1}$ , since no suppression by  $1/m_b^2$

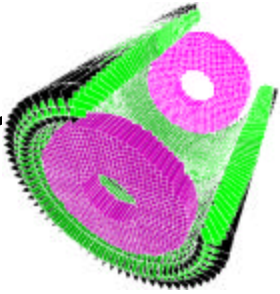
$$L_f = L_i \pm 1$$

as  $\Delta n = n_i - n_f$  increases

$$\langle n_f L_f | r | n_i L_i \rangle \text{ decreases}$$

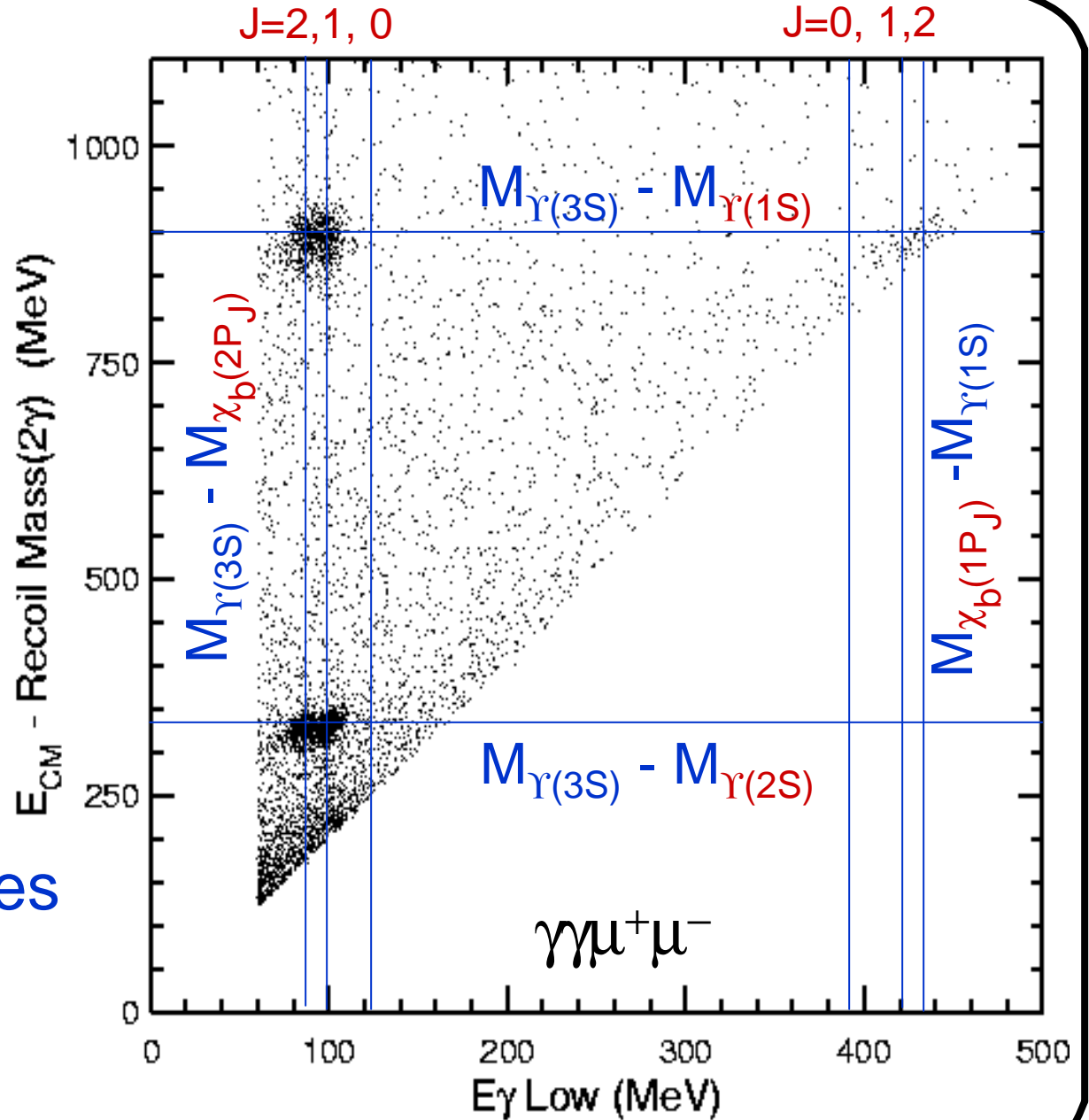
and  $\Gamma_{E1}$  becomes more difficult to predict

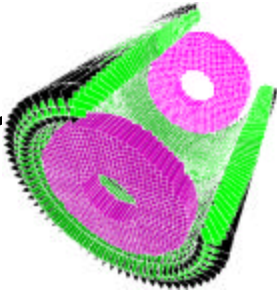




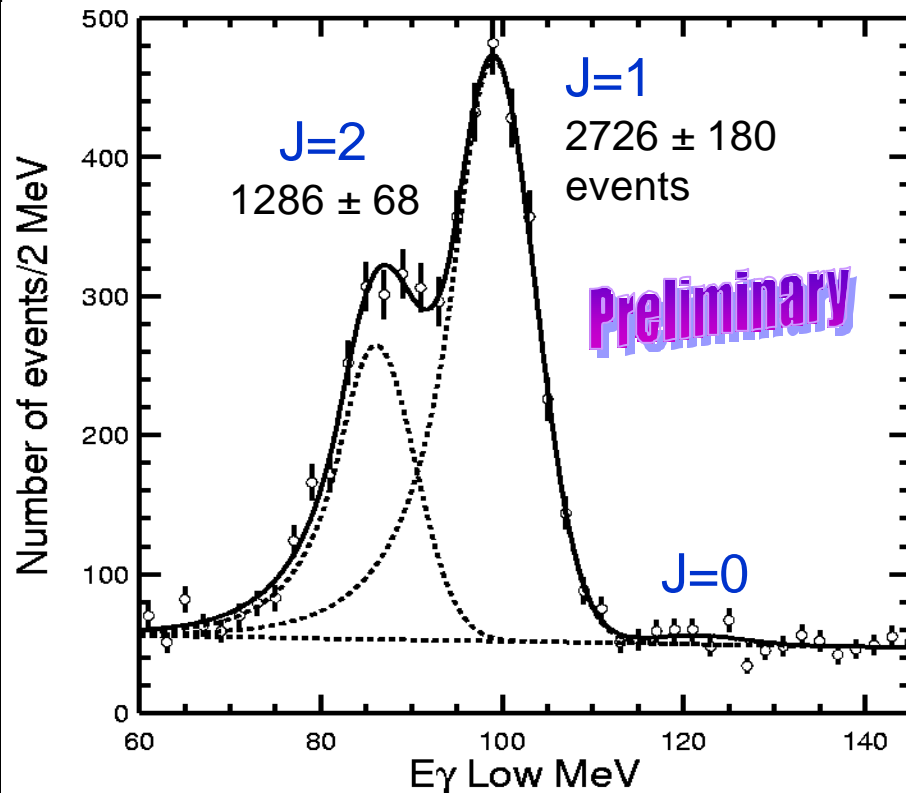
# Exclusive 2 $\gamma$ -cascades

- Signal variables





$$\Upsilon(3S) \mapsto \gamma \chi_b(2P_J) \mapsto \gamma \gamma l^+ l^-$$



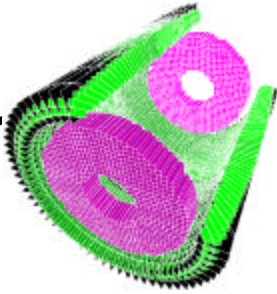
Efficiency  
( $\mu\mu+ee$ )/2  
~32%

Energy resolution from the fit:  
4.6 $\pm$ 0.2 MeV @ 100 MeV

- Energy calibrated to  $\pm 0.34\%$  with the photon-recoil mass and known  $\Upsilon(nS)$  masses

- Photon-line energies  $\mapsto M(\chi_b(2P_J))$ 
  - $E_{J=2} = 86.09 \pm 0.30 \pm 0.29$  MeV
  - $E_{J=1} = 99.08 \pm 0.17 \pm 0.34$  MeV

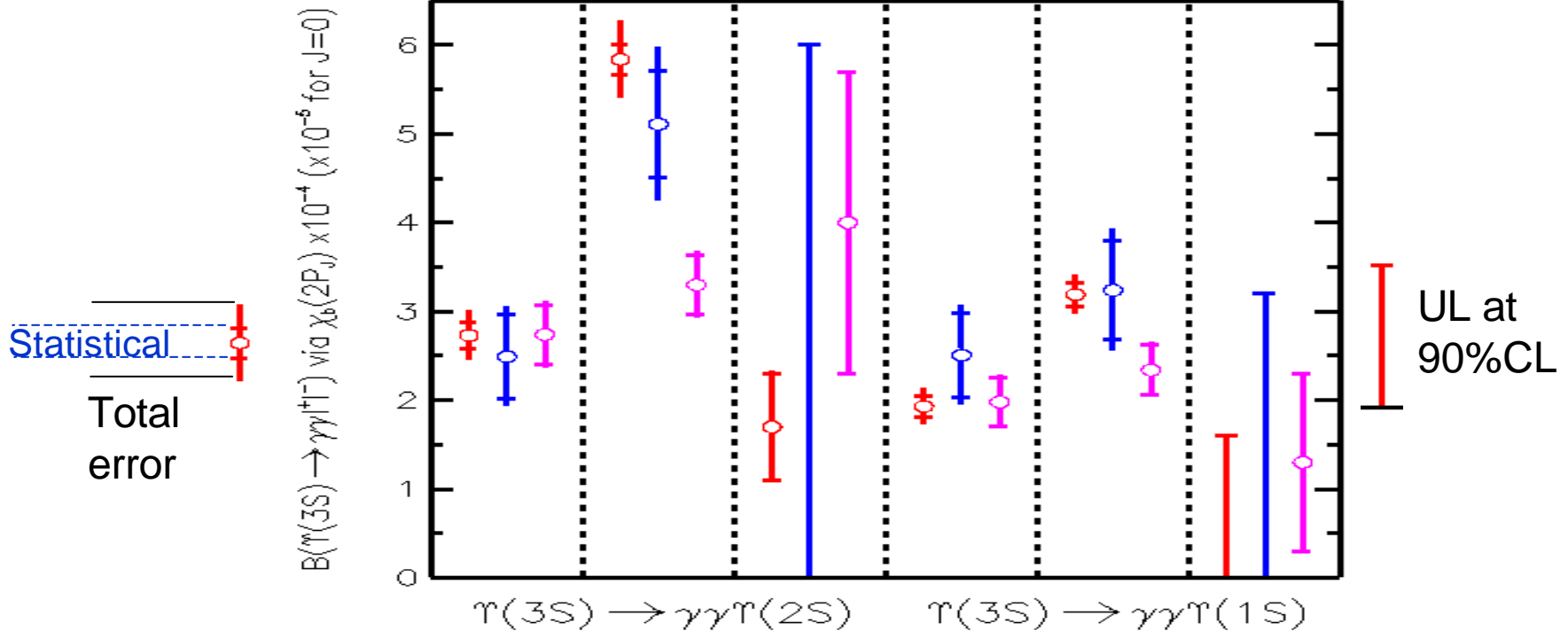
- More precise than previous measurements
- Consistent with CUSB and CLEO-II results



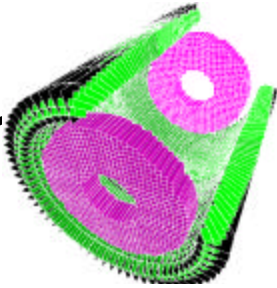
$$\Upsilon(3S) \mapsto \gamma \chi_b(2P_J) \mapsto \gamma \gamma l^+ l^-$$

CLEO-III, CLEO-II, CUSB

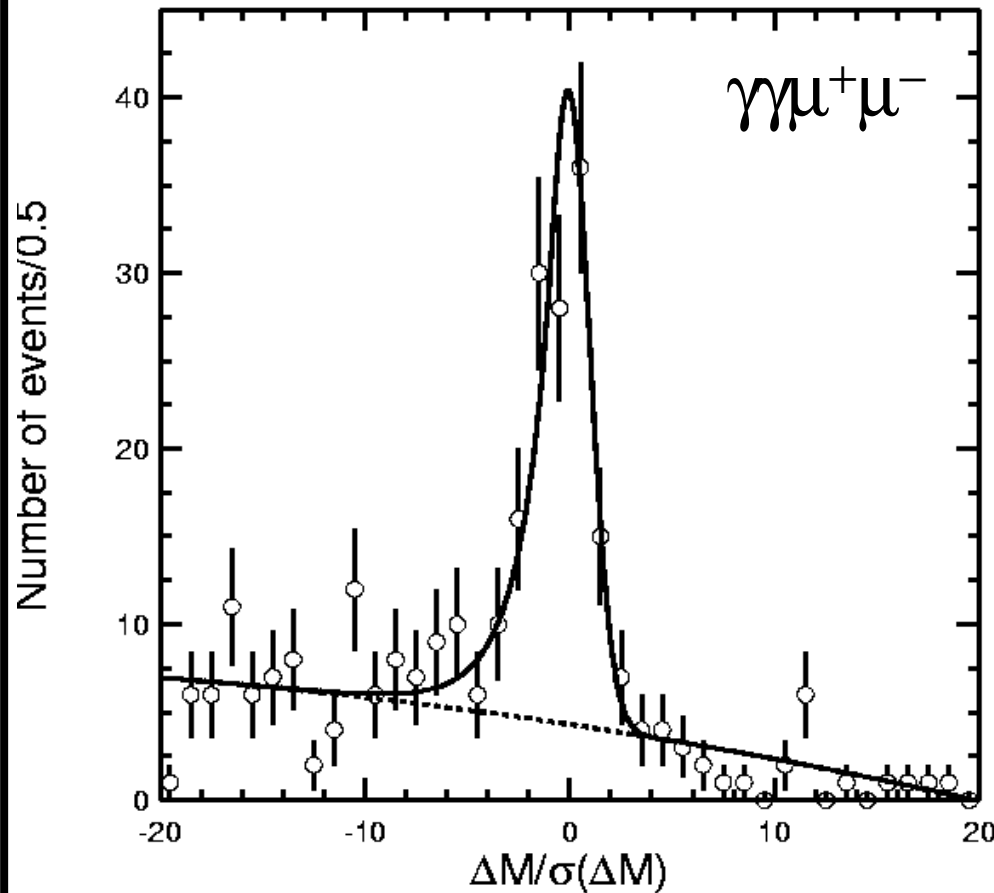
J=2 J=1 J=0 J=2 J=1 J=0



- Throughout this talk:  $B(\gamma\gamma l^+ l^-) = ( B(\gamma\gamma\mu^+\mu^-) + B(\gamma\gamma e^+e^-) )/2$
- Good agreement with CLEO-II
- CUSB measurements low for J=1
- Much improved errors in CLEO-III



$$\Upsilon(3S) \mapsto \gamma \chi_b(1P_J) \mapsto \gamma \ell^+ \ell^-$$



- Since cannot resolve J=2,1,0 states, fit the recoil mass distribution instead

Events ( $\mu\mu+ee$ ):  $167 \pm 19$  events

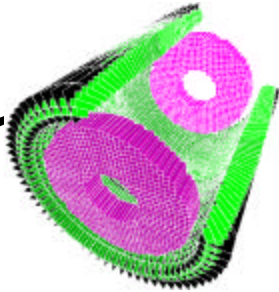
Efficiency:  $\sim 37\%$

$$B(\Upsilon(3S) \mapsto \gamma \chi_b(1P_{2,1}) \mapsto \gamma \Upsilon(1S) \mapsto \gamma \ell^+ \ell^-) = (5.2 \pm 0.5 \pm 0.5) 10^{-5}$$

$$B(\Upsilon(3S) \mapsto \gamma \chi_b(1P_{2,1}) \mapsto \gamma \Upsilon(1S)) = (2.1 \pm 0.2 \pm 0.2) 10^{-3}$$

CUSB

$$(1.2 \pm 0.4 \pm 0.1) 10^{-3}$$

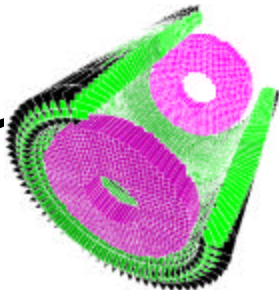


## Comparison of the measured E1 transition rates with the potential models

	$\langle 2P r 3S \rangle$		$\langle 1P r 2S \rangle$		$\langle 1P r 3S \rangle$		$\langle 1S r 2P \rangle$	
	$GeV^{-1}$		$GeV^{-1}$		$GeV^{-1}$		$\langle 2S r 2P \rangle$	
DATA	2.7±0.2		1.9±0.2		0.050±0.006		0.096±0.005	
	World Average				This measurement			
Model	NR	rel	NR	rel	NR	rel	NR	rel
Kwong,Rosner [13]	2.7		1.6		0.023		0.13	
Fulcher [14]	2.6		1.6		0.023		0.13	
Büchmuller et al.[15]	2.7		1.6		0.010		0.12	
Moxhay,Rosner [16]	2.7	2.7	1.6	1.6	0.024	0.044	0.13	0.15
Gupta et al.[17]	2.6		1.6		0.040		0.11	
Gupta et al.[18]	2.6		1.6		0.010		0.12	
Fulcher [19]	2.6		1.6		0.018		0.11	
Danghighian et al.[20]	2.8	2.5	1.7	1.3	0.024	0.037	0.13	0.10
McClary,Byers [21]	2.6	2.5	1.7	1.6			0.15	0.13
Eichten et al.[22]	2.6		1.7		0.110		0.15	
Grotch et al.[23]	2.7	2.5	1.7	1.5	0.011	0.061	0.13	0.19

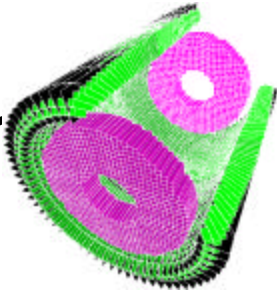
- Potential models:**

- easily reproduce the large E1 matrix elements
- have trouble predicting small elements  
(see  $\Upsilon(3S) \mapsto \gamma \chi_b(1P_J) \quad \Delta n=2$ )



## Searches for $\pi^0, \eta$ transitions

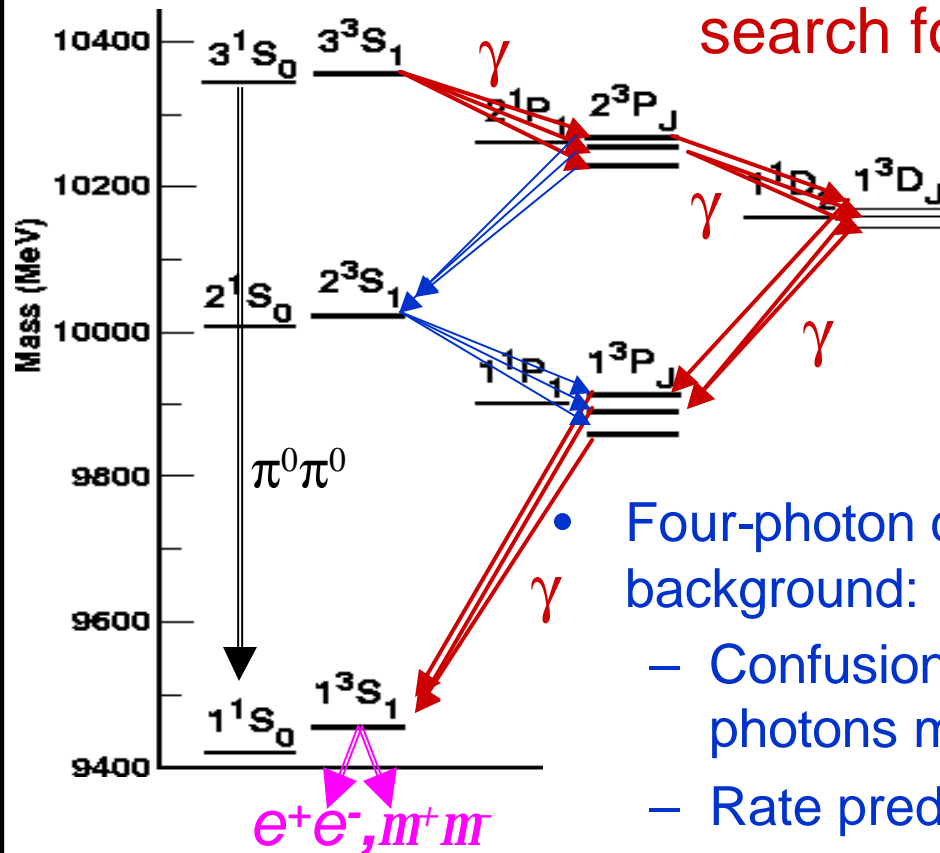
- Also could contribute to  $\gamma\gamma\ell^+\ell^-$  events
- Suppress photon transitions
- Look at  $M\gamma\gamma$
- No signal found
- At 90% C.L.
  - $B(\Upsilon(3S) \mapsto \pi^0 \Upsilon(1S)) < 0.17 \cdot 10^{-3}$
  - $B(\Upsilon(3S) \mapsto \eta \Upsilon(1S)) < 0.9 \cdot 10^{-3}$
  - $B(\Upsilon(3S) \mapsto \pi^0 \Upsilon(2S)) < 1.2 \cdot 10^{-3}$



ICHEP ABS948, CLEO CONF 02-06

# Exclusive 4 $\gamma$ -cascades

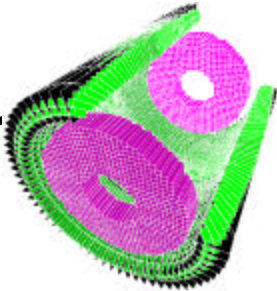
- Can use four-photon E1 cascade to search for  $\Upsilon(1^3D_J)$ !



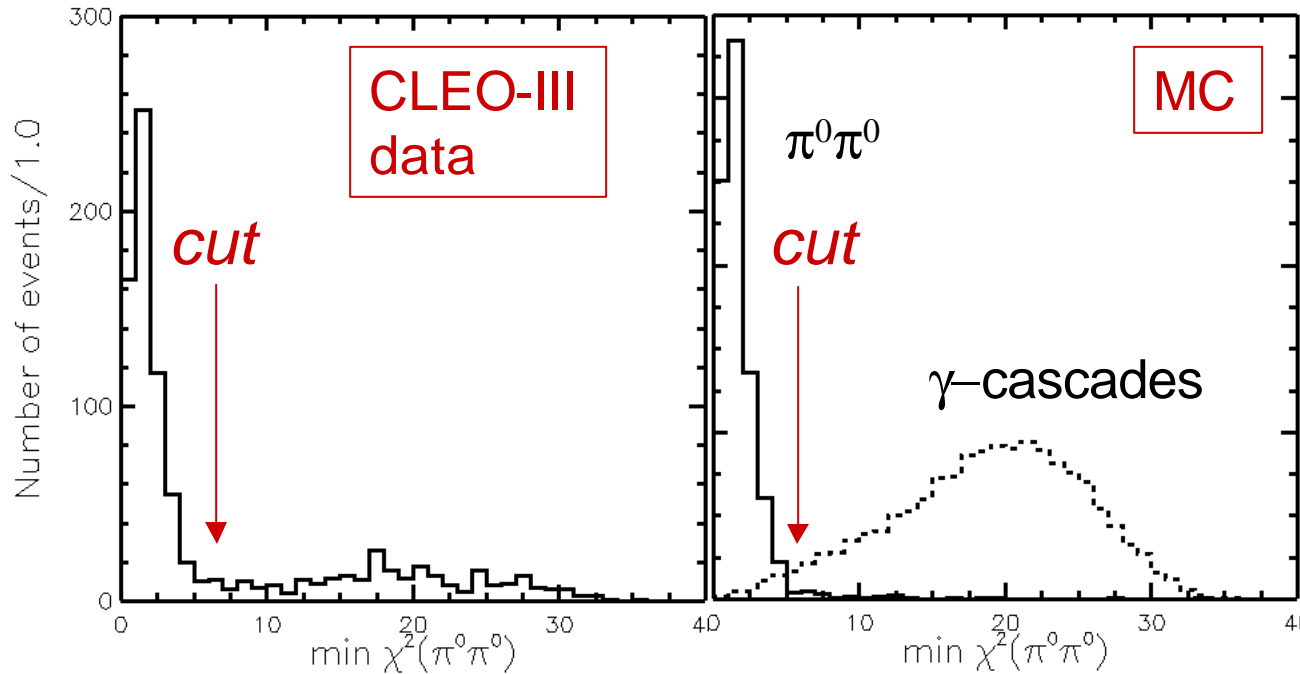
Signal rate predicted by  
Godfrey&Rosner:  $3.76 \cdot 10^{-5}$   
PR D64, 097501 (2001)

- Four-photon cascade via the  $\Upsilon(2S)$  – the main background:
  - Confusion in ordering of the observed photons makes these two cascades similar
  - Rate predicted by Godfrey&Rosner:  $3.84 \cdot 10^{-5}$

- Also  $\Upsilon(3S) \mapsto \pi^0\pi^0 \Upsilon(1S)$  is a potential background



# $\Upsilon(3S) \mapsto \pi^0 \pi^0 \Upsilon(1S)$



- Selection/  
rejection of  
 $\Upsilon(3S) \mapsto$   
 $\pi^0 \pi^0 \Upsilon(1S)$

Events ( $\mu\mu+ee$ ):  
 $737 \pm 28$  events

Efficiency:  
 $\sim 14\%$

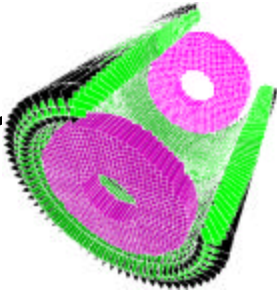
$$B(\Upsilon(3S) \mapsto \pi^0 \pi^0 \Upsilon(1S) \mapsto \gamma \gamma l^+ l^-) = (5.7 \pm 0.2 \pm 0.4) 10^{-4}$$

$$B(\Upsilon(3S) \mapsto \pi^0 \pi^0 \Upsilon(1S)) = (2.33 \pm 0.09 \pm 0.16) 10^{-2}$$

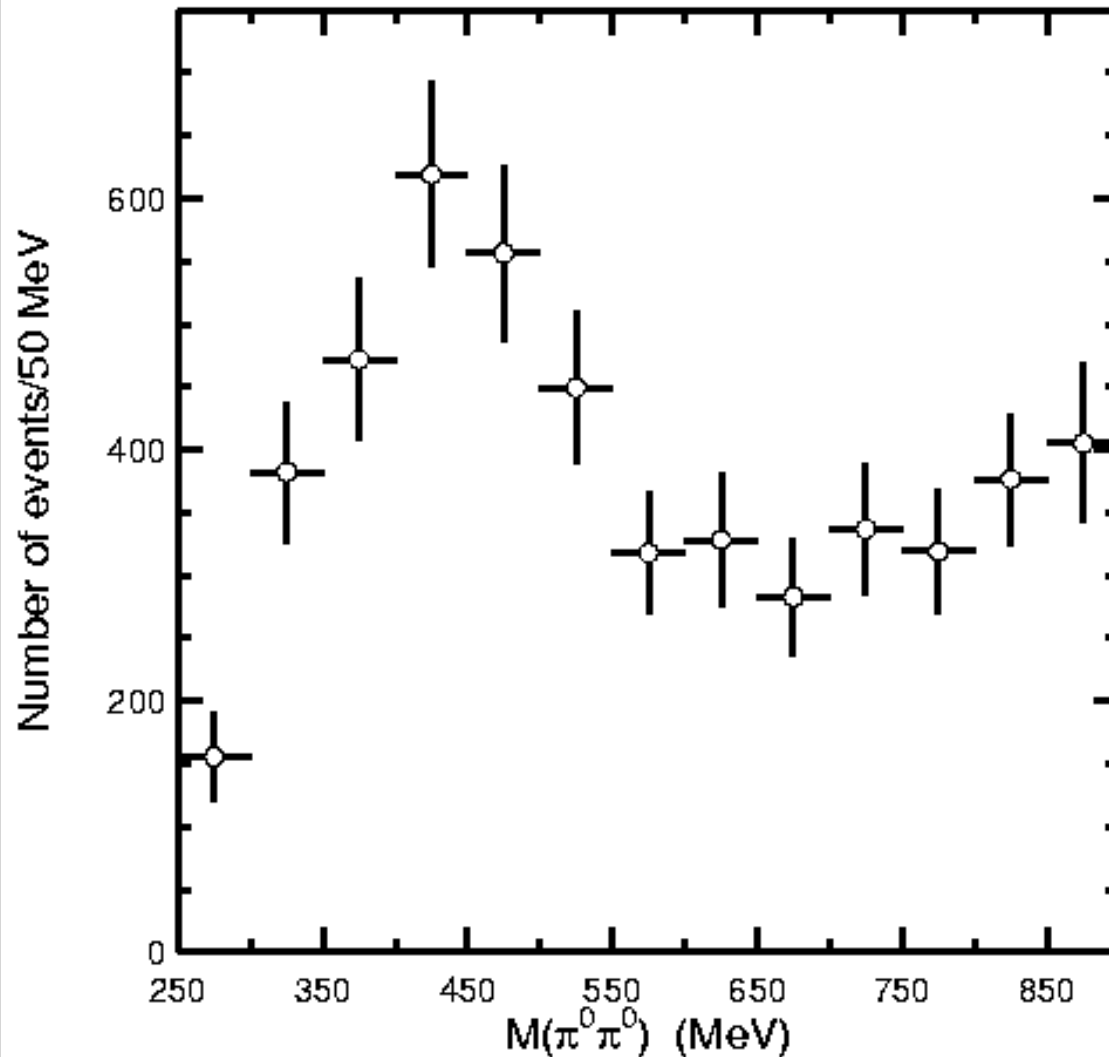
$$\text{CLEO-II} \quad (2.03 \pm 0.28 \pm 0.19) 10^{-2}$$

$$\text{CUSP} \quad (2.3 \pm 0.4 \pm 0.3) 10^{-2}$$

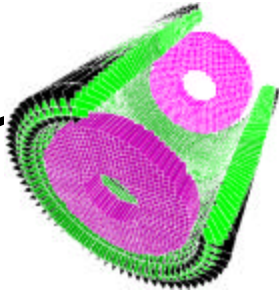




$$\Upsilon(3S) \mapsto \pi^0 \pi^0 \Upsilon(1S)$$



- Double-peak structure in  $M(\pi^0 \pi^0)$  confirmed here with more data



## Selection of $\Upsilon(1D)$ events

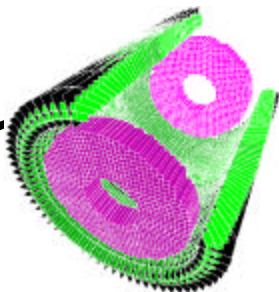
$$c_{1D}^2 = \min_{M_{1D}, J_{2P}, J_{1P}} \sum_{i=1}^4 \left( \frac{E_{gi} - E_{gi}^{\text{expected}}(M_{1D}, J_{2P}, J_{1P})}{s(E_{gi})} \right)^2$$

- Implements constraints to the well known masses:  $M_{3S}$ ,  $M_{2P_J}$ ,  $M_{1P_J}$ ,  $M_{1S}$
- In addition to  $\chi^2_{1D}$  value also obtain “most likely” mass of  $\Upsilon(1D)$  for each event
- To suppress cascades through the  $\Upsilon(2S)$

calculate:

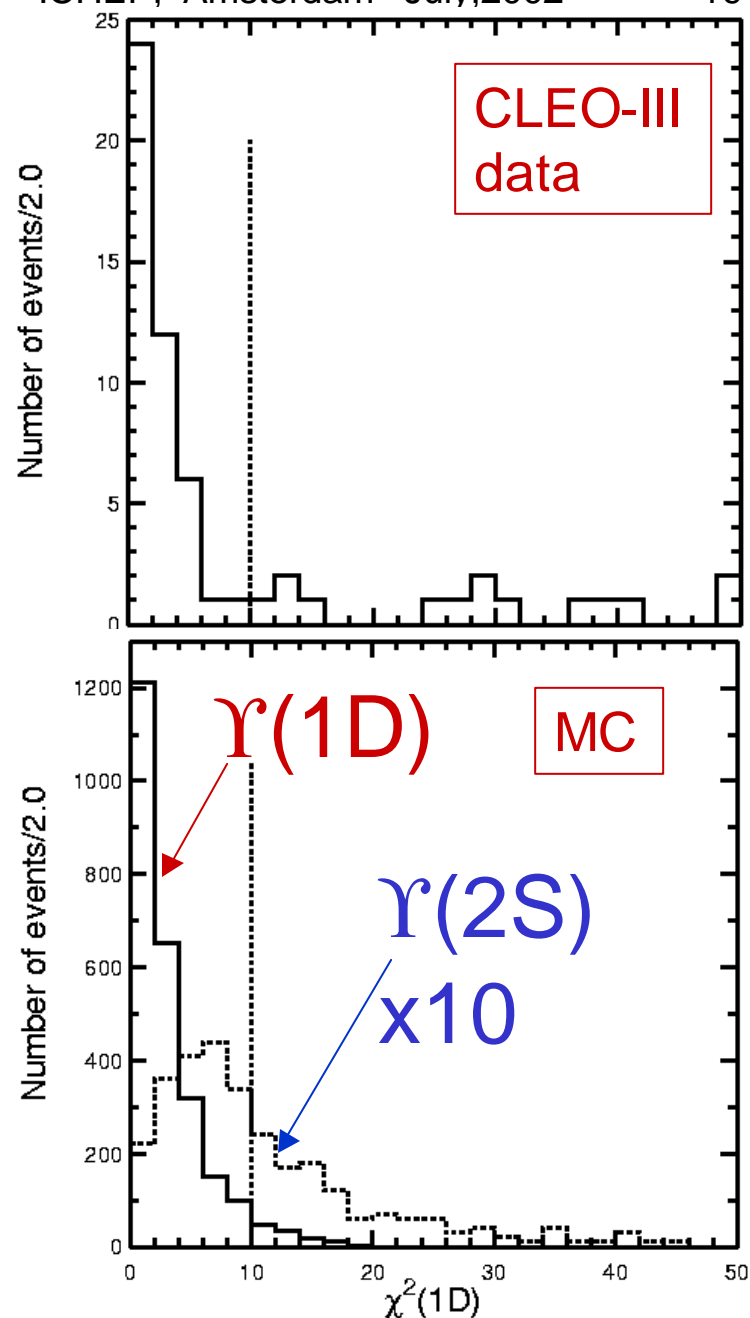
$$c_{2S}^2 = \min_{J_{2P}, J_{1P}} \sum_{i=1}^4 \left( \frac{E_{gi} - E_{gi}^{\text{expected}}(M_{2S}, J_{2P}, J_{1P})}{s(E_{gi})} \right)^2$$

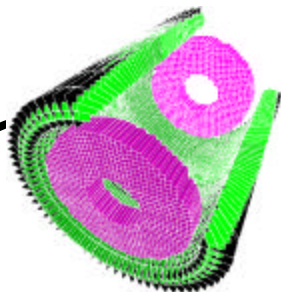
$$c_{2S}^2 > 12$$



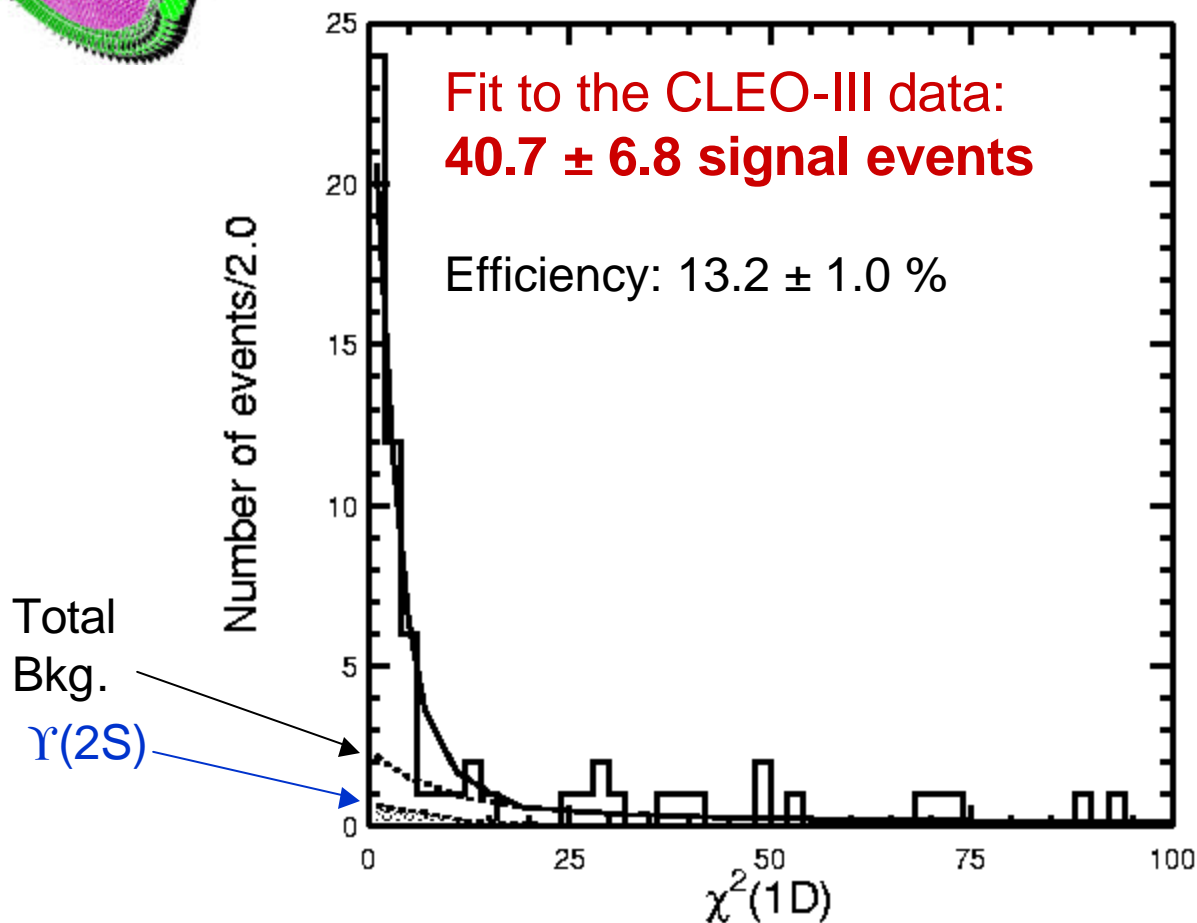
## Inclusive $\Upsilon(1D)$ signal

- No background source can produce as narrow a peak as observed in the data
- For  $\chi^2_{1D} < 10$ :
  - 44 events in the data
  - 1.6-3.0 events due to  $\Upsilon(2S)$
  - 0.8 events due to  $\Upsilon(3S) \mapsto \pi^0\pi^0\Upsilon(1S)$
  - 1.8-3.7 of other backgrounds (e.g. radiative Bhabhas and  $\mu$ -pairs) estimated from the tail of the distribution
  - **Total background 10-14%**





# Inclusive $\Upsilon(1D)$ signal



The signal significance estimated from the fits with the background contributions alone:

$\mu\mu+ee$ : **9.7s**

$\mu\mu$ :  $8.2\sigma$

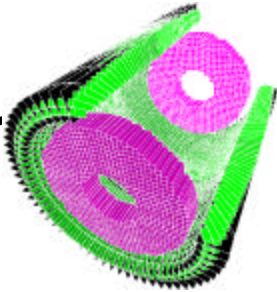
$ee$ :  $5.2\sigma$

$$B(\Upsilon(3S) \rightarrow gg\Upsilon(1D) \rightarrow gg\gamma\Upsilon(1S) \rightarrow gggg\ell^+\ell^-) = (3.3 \pm 0.6 \pm 0.5) 10^{-5}$$

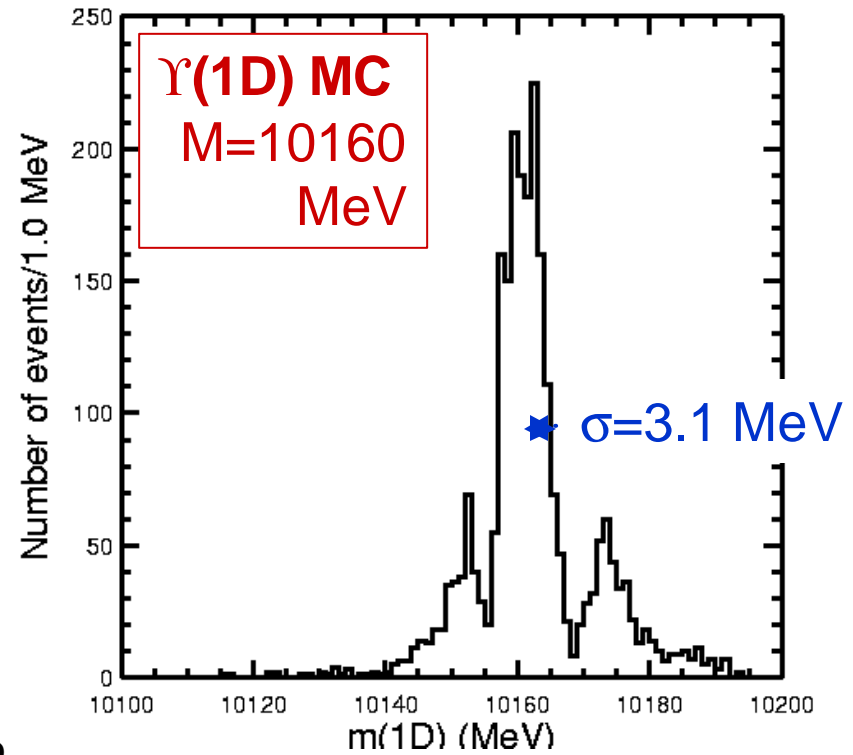
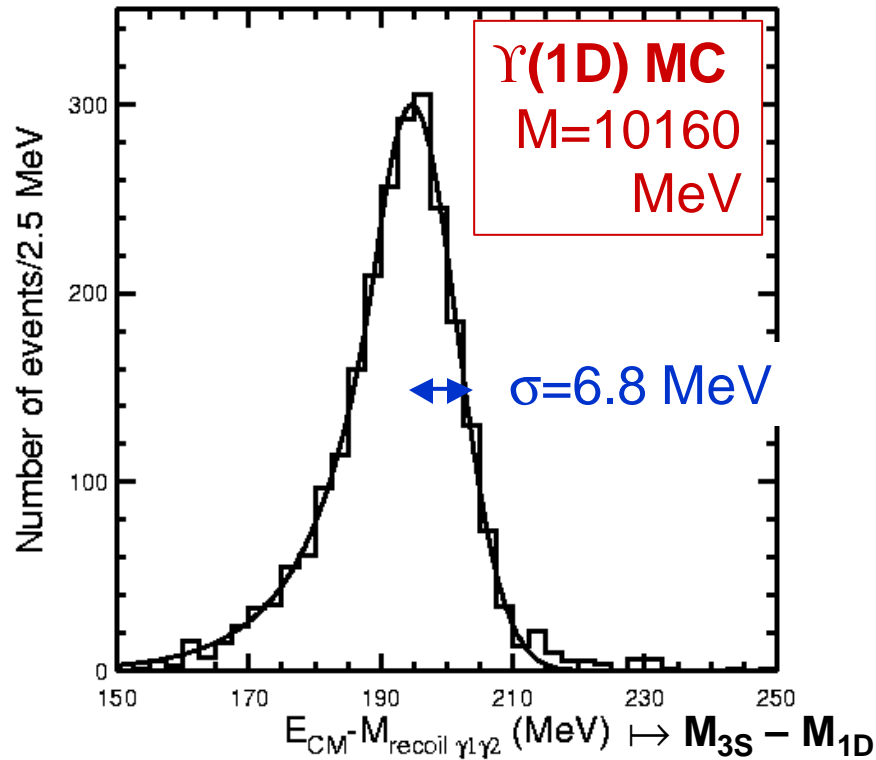
Godfrey&Rosner

3.8

$10^{-5}$

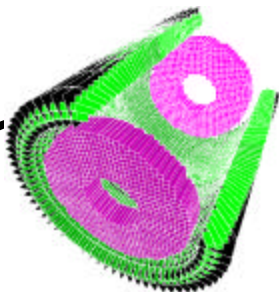


## $\Upsilon(1D)$ mass estimators



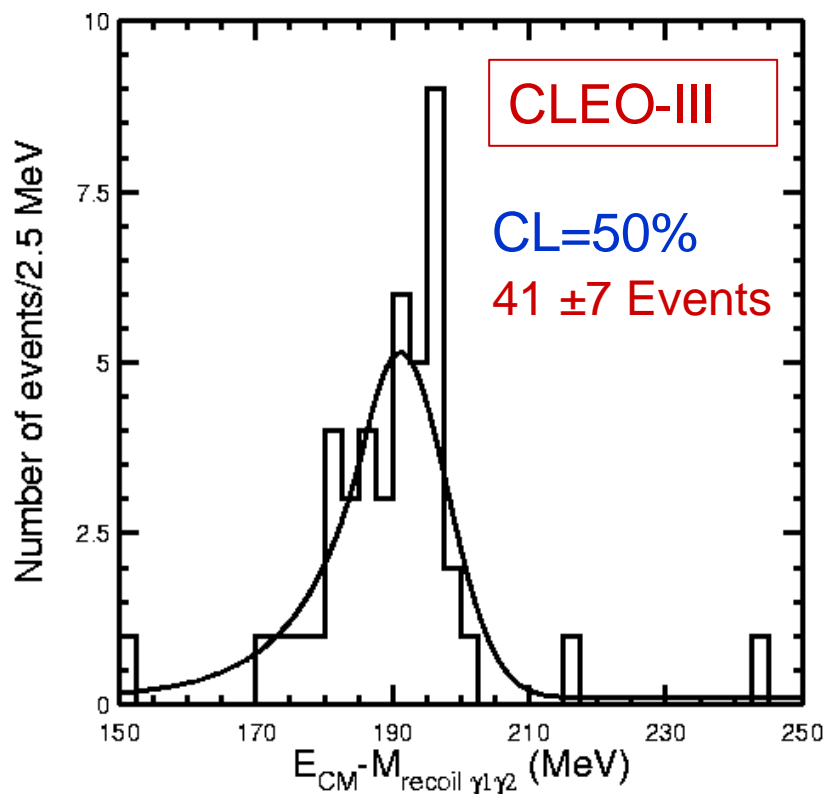
- Recoil mass against the two lowest energy photons:
  - Worse resolution
  - Simple shape

- Most likely mass (constrained to 2P, 1P masses):
  - Better resolution
  - Satellite peaks due to wrong  $J_{2P}, J_{1P}$  minimizing the  $\chi^2_{1D}$

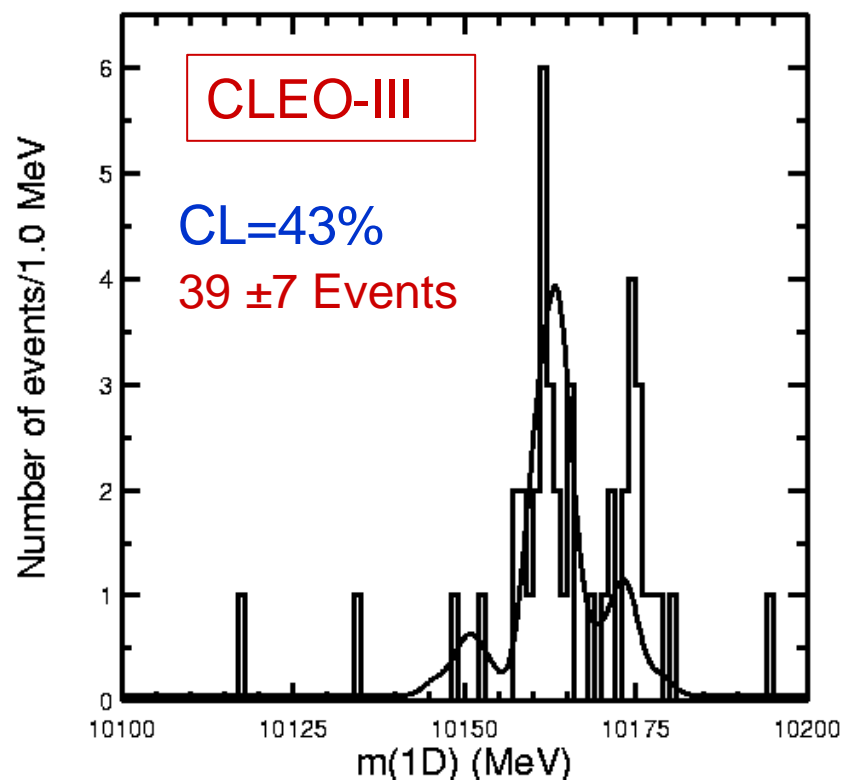


## $\Upsilon(1D)$ mass analysis

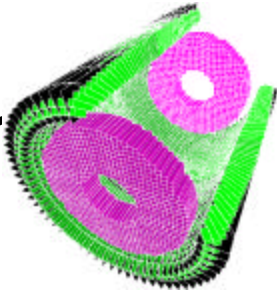
- Single-peak fits:



$M = 10163.4 \pm 1.3$  MeV

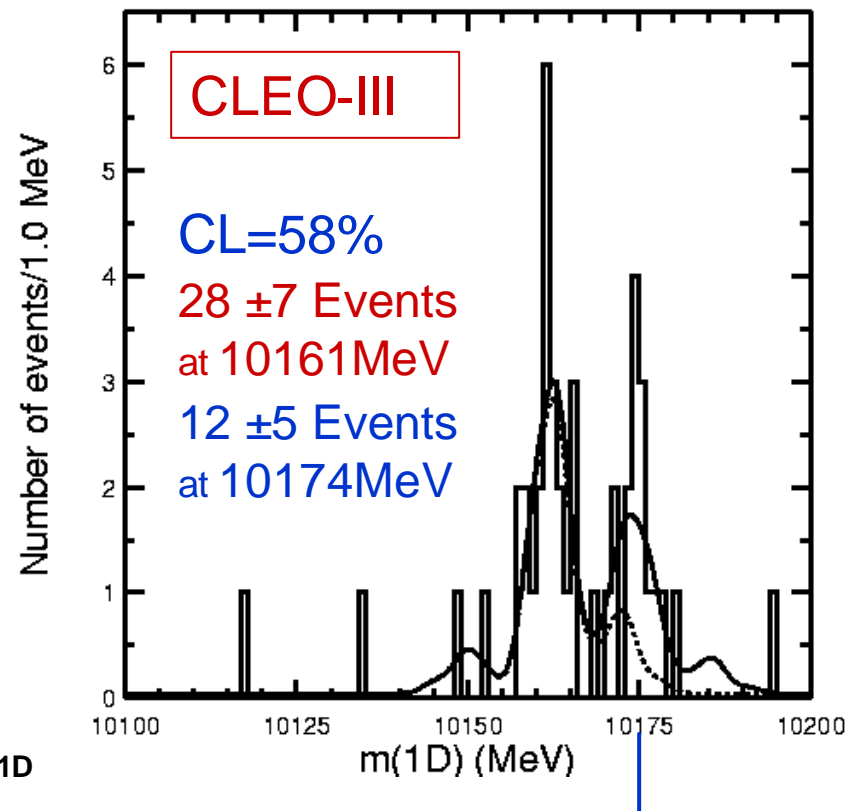
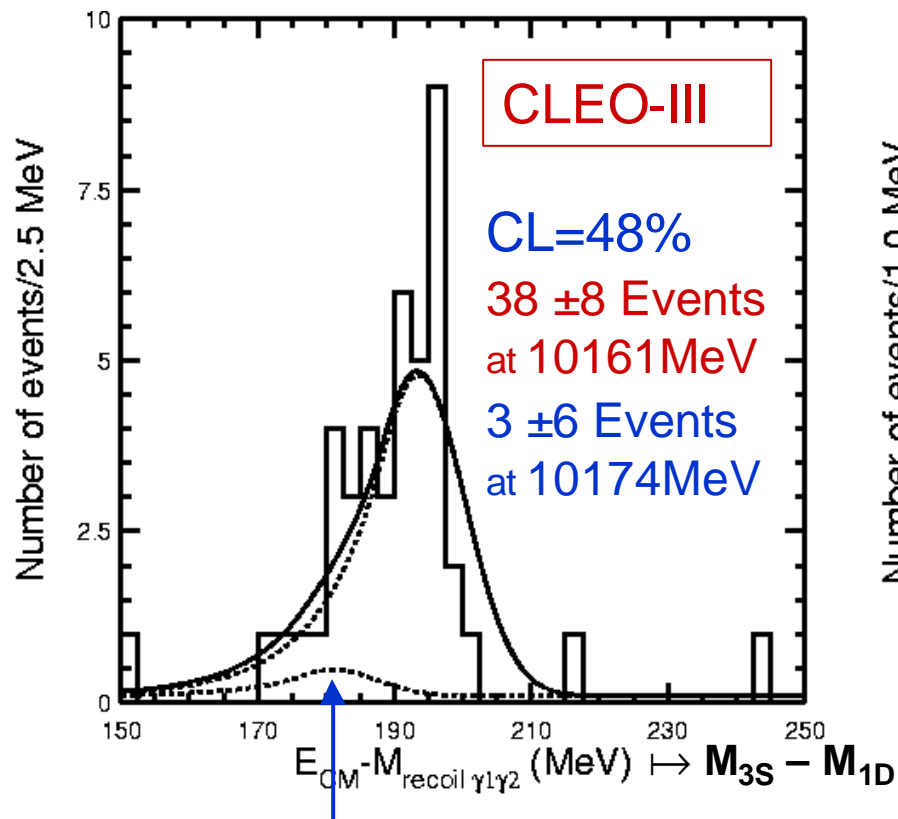


$M = 10162.0 \pm 0.5$  MeV



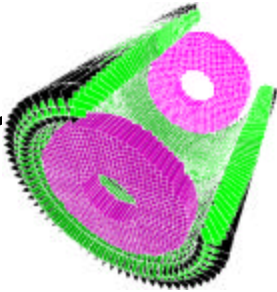
## $\Upsilon(1D)$ mass analysis

- Double-peak fits:



Masses fixed from the fit  
shown on the right

$M = 10161.2 \pm 0.7 \text{ MeV}$



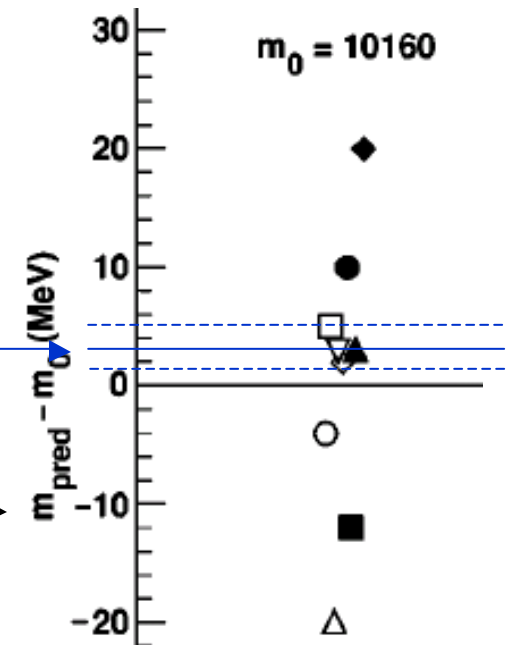
## $\Upsilon(1D)$ mass analysis

- No compelling evidence for more than one state
- Significance of the peak at 10162: **6.8s**
- Mass averaged over different fits:  **$10162.2 \pm 1.6$  MeV**
- Inconsistent with the  $\Upsilon(1D_3)$
- Could be the  $\Upsilon(1D_2)$  or  $\Upsilon(1D_1)$
- The theory predicts the rate ratio:  $\Upsilon(1D_2)/\Upsilon(1D_1)=6$
- **Thus, the  $\Upsilon(1D_2)$  is the most likely interpretation**

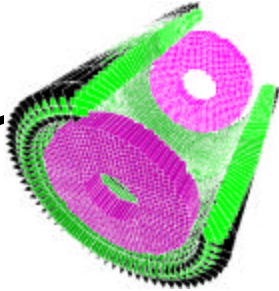
All calculations of the fine splitting predict the  $\Upsilon(1D_2)$  mass from  $-0.5$  to  $-1.0$  MeV below the center-of-gravity of the triplet

→ Our mass measurement is consistent with the c.o.g.  $\sim 10163 \pm 2$  MeV

Spread in the predictions of the center-of-gravity of the triplet 1D states by various potential models (from Godfrey&Rosner)

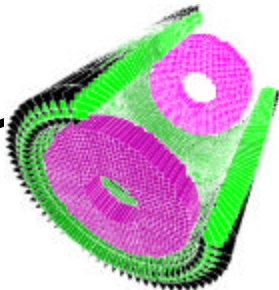






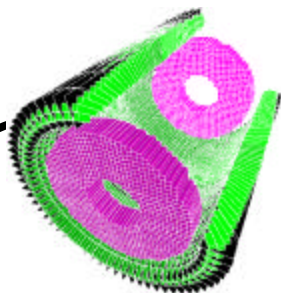
## Summary

- **No evidence for hindered M1 transitions**  
 $\Upsilon(3S) \mapsto gh_b(1S)$  found in contradiction with many theoretical estimates of the transition width
- **Much improved results for:**
  - $B(\Upsilon(3S) \mapsto gC_b(1P_{2,1}) \mapsto gg\Upsilon(1S))$
  - $B(\Upsilon(3S) \mapsto gC_b(2P_{2,1,0}) \mapsto gg\Upsilon(2,1S))$
  - $B(\Upsilon(3S) \mapsto p^0p^0\Upsilon(1S))$
  - **Upper limits on:**
    - $B(\Upsilon(3S) \mapsto p^0\Upsilon(1S))$
    - $B(\Upsilon(3S) \mapsto h\Upsilon(1S))$
    - $B(\Upsilon(3S) \mapsto p^0\Upsilon(2S))$



## Summary

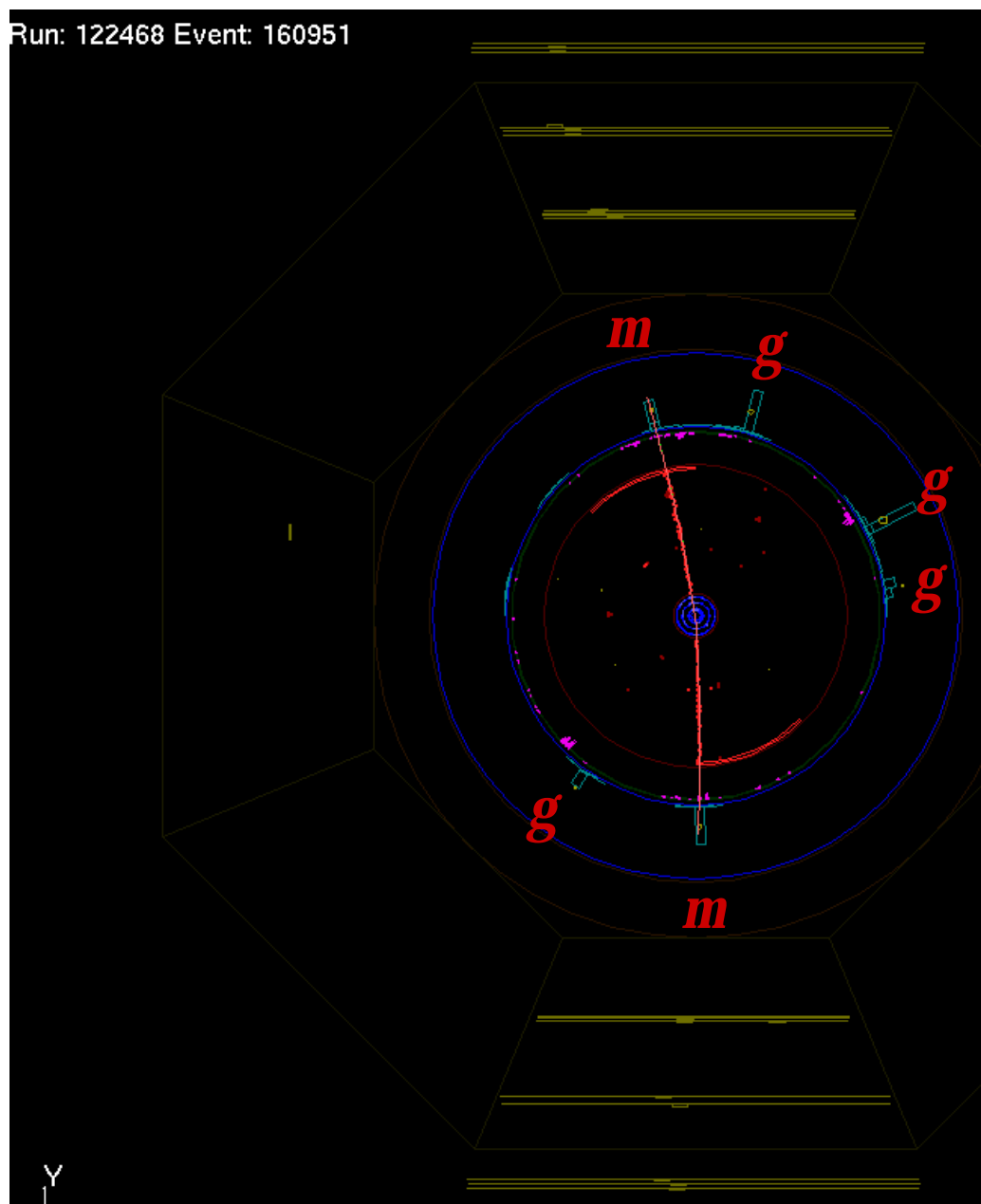
- **First observation of  $\Upsilon(1D)$ :**
  - Signal is **9.7s significant**
  - Inclusive (i.e. sum over all J) product branching ratio for production in  $gggg\ell^+\ell^-$   **$(3.3\pm 0.6\pm 0.5) 10^{-5}$**
  - In agreement with the prediction by **Godfrey&Rosner ( $3.8 10^{-5}$ )**
  - Evidence for a state at  **$10162.2 \pm 1.6$  MeV**
  - Likely interpretation:  **$\Upsilon(1D_2)$**
  - The mass is consistent with the predictions of some of the potential models
  - First new narrow bb state observed in 19 years
  - The only long-lived L=2 meson we know

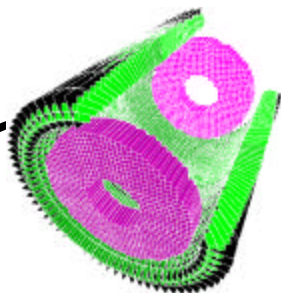


(Extra Slide)

Run: 122468 Event: 160951

$\gamma\gamma\mu\mu$   
1D candidate

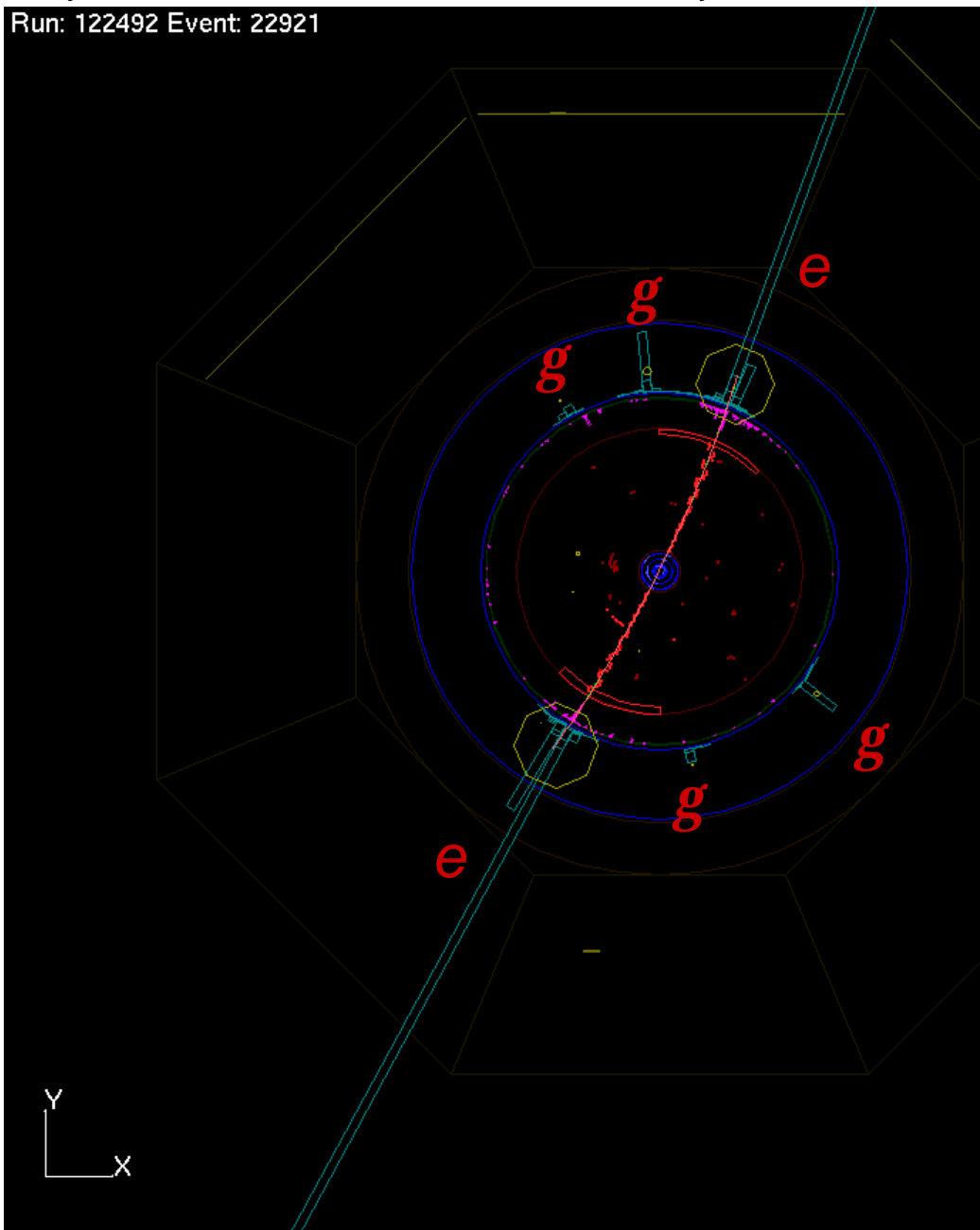


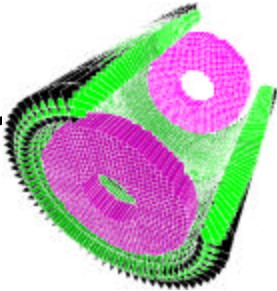


(Extra Slide)

$\gamma\gamma ee$  1D  
candidate

Run: 122492 Event: 22921

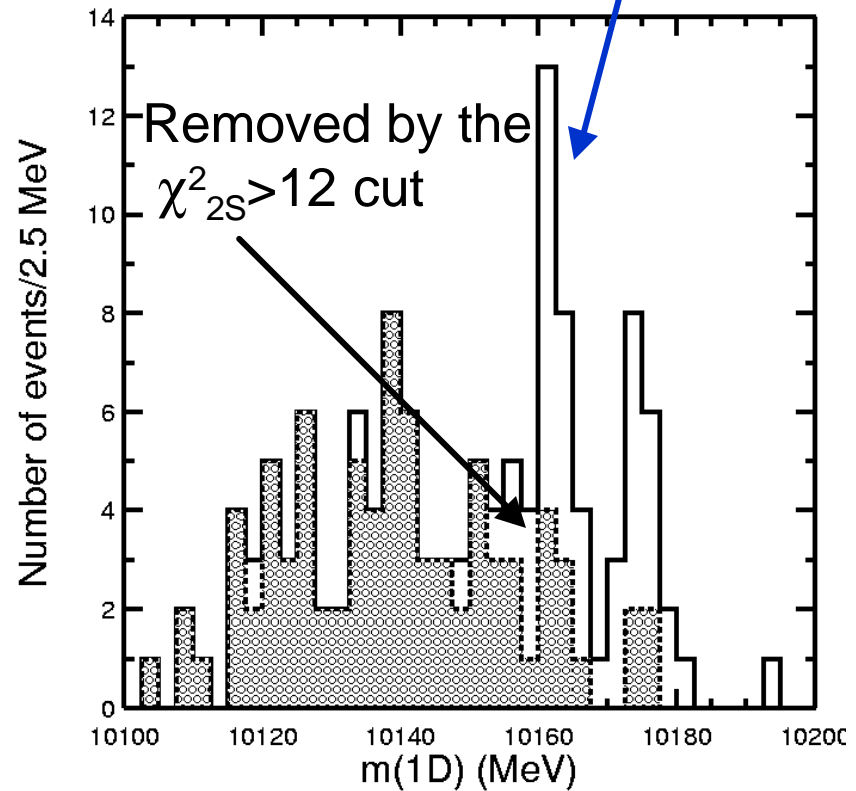
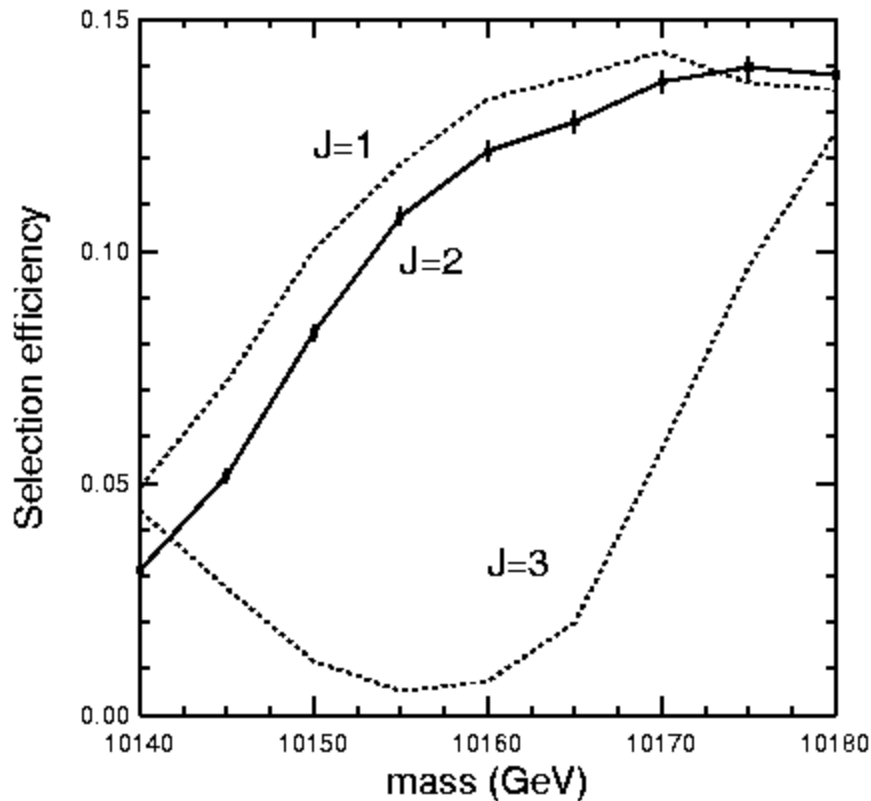




(Extra Slide)

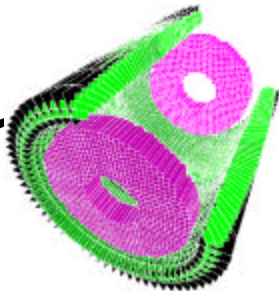
# Ruling out $J_{1D}=3$

No  $\chi^2_{2S} > 12$  cut



- No efficiency for the  $\Upsilon(1D_3)$  at 10162 MeV because of the  $\chi^2_{2S} > 12$  cut

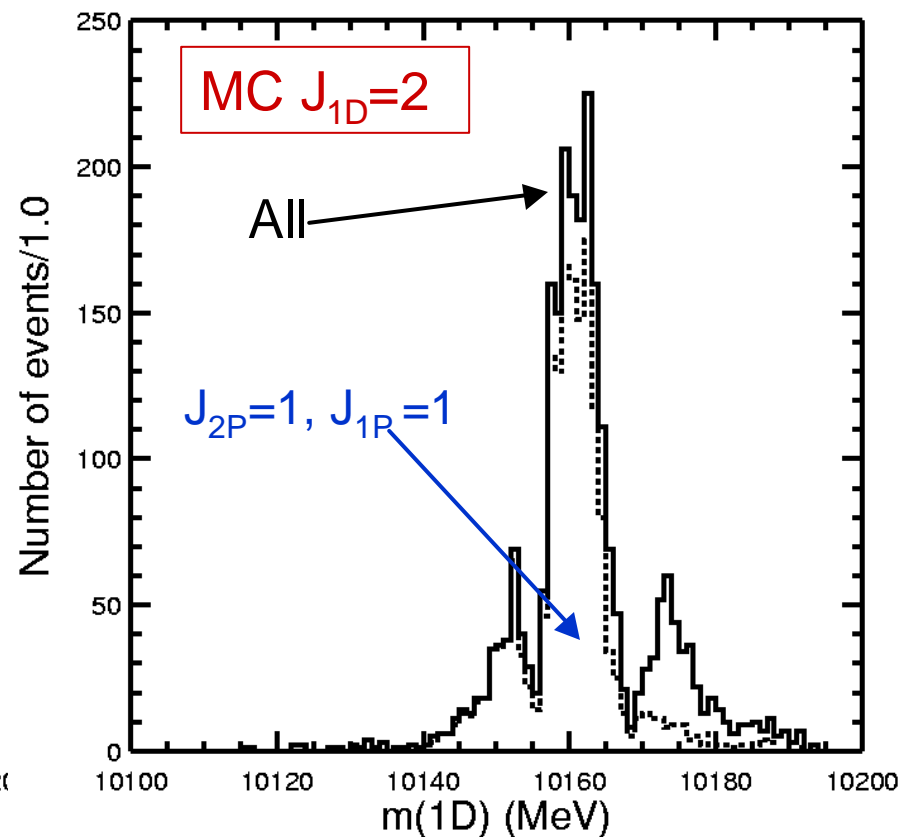
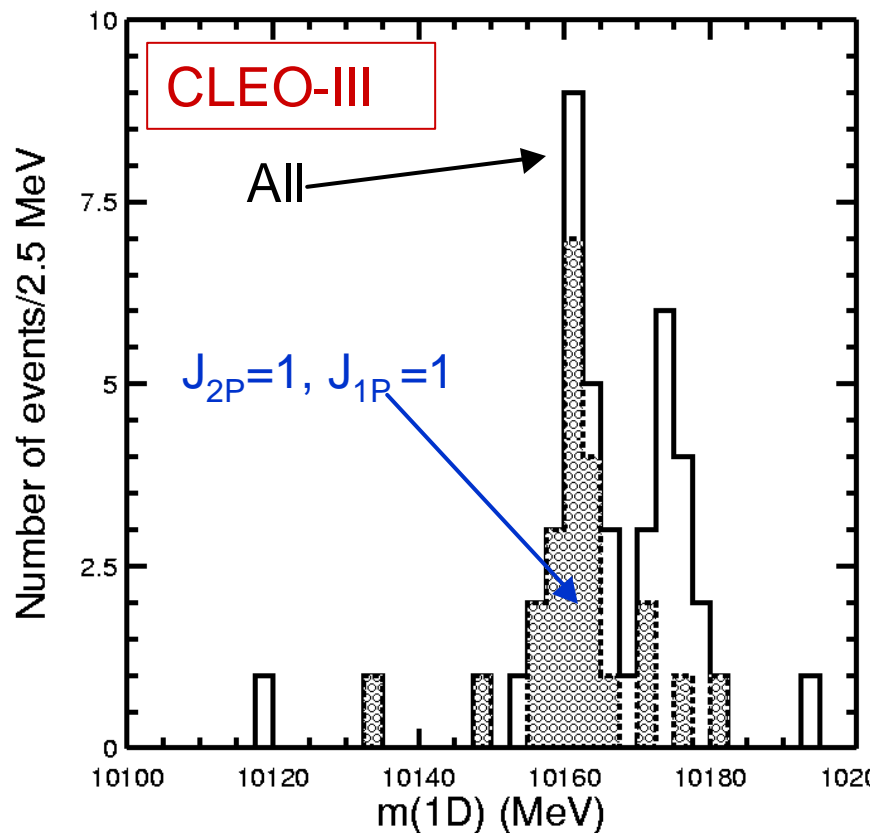
- However the  $\chi^2_{2S} > 12$  cut does not change the 10162 peak amplitude much

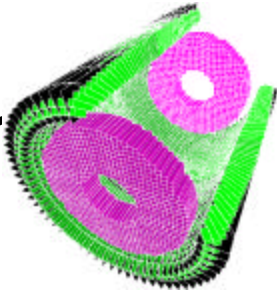


(Extra Slide)

## Probing $J_{1D}$ via $J_{2P}$ , $J_{1P}$

- The peak at 10162 MeV has a large fraction of  $J_{2P}=1$ ,  $J_{1P}=1$  events, as expected for  $J_{1D}=1$  or 2.

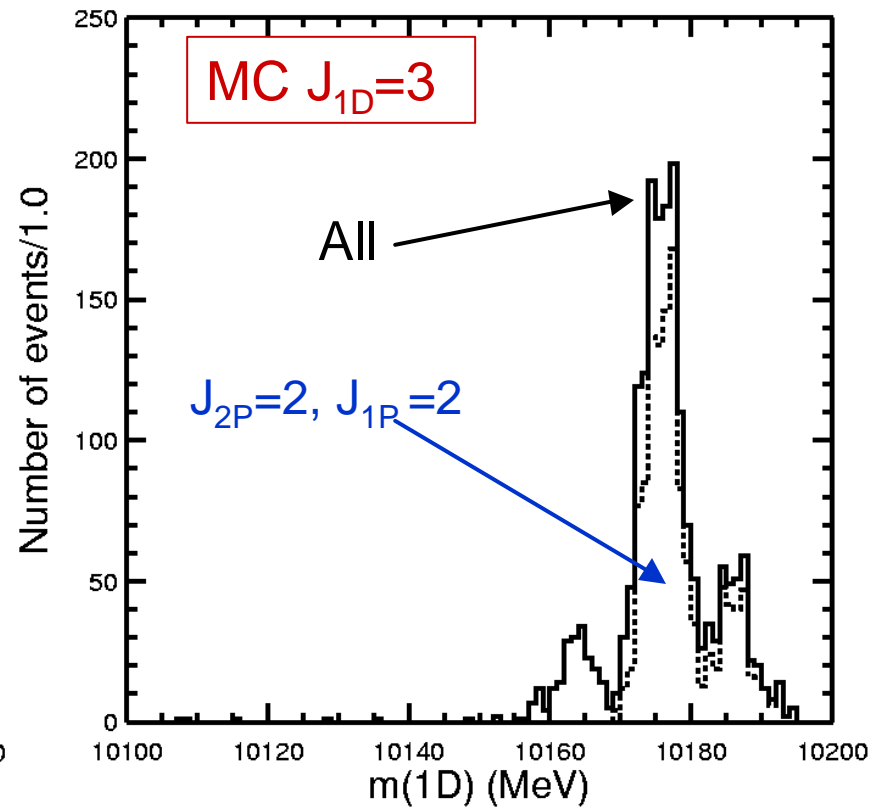
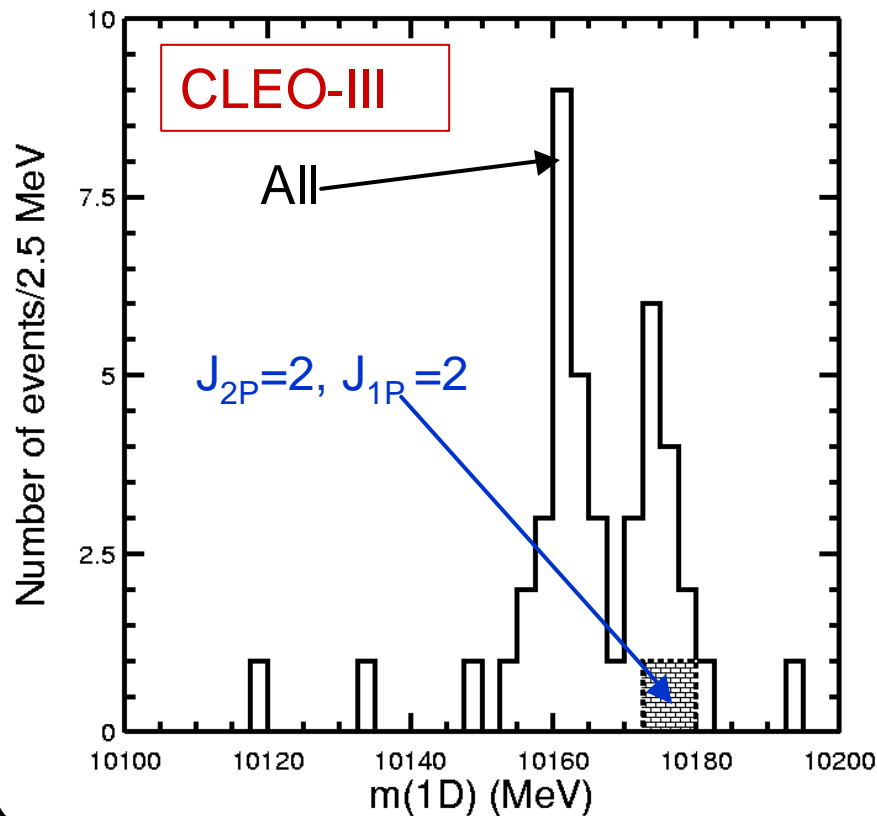


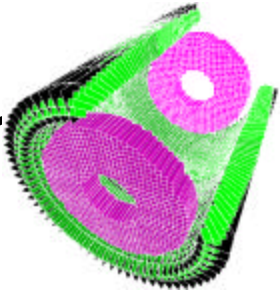


(Extra Slide)

## Probing $J_{1D}$ via $J_{2P}$ , $J_{1P}$

- The peak at 10174 MeV has a too small a fraction of  $J_{2P}=2$ ,  $J_{1P}=2$  events, to be  $J_{1D}=3$  (in fact a few such events expected due to the satellite from  $J_{1D}=2$  at 10162 MeV)



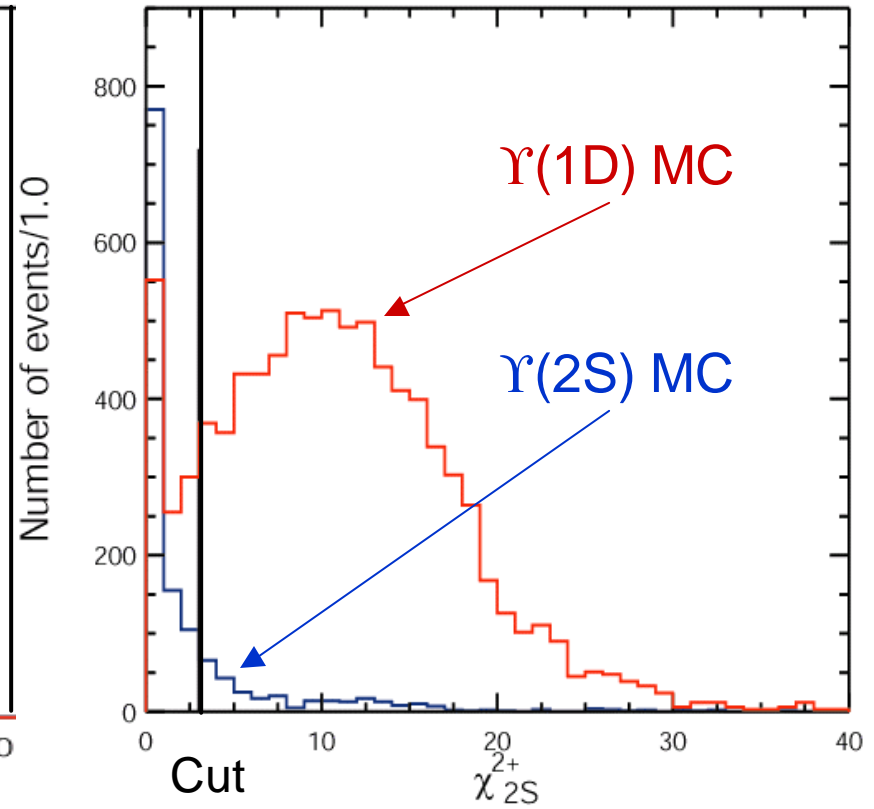
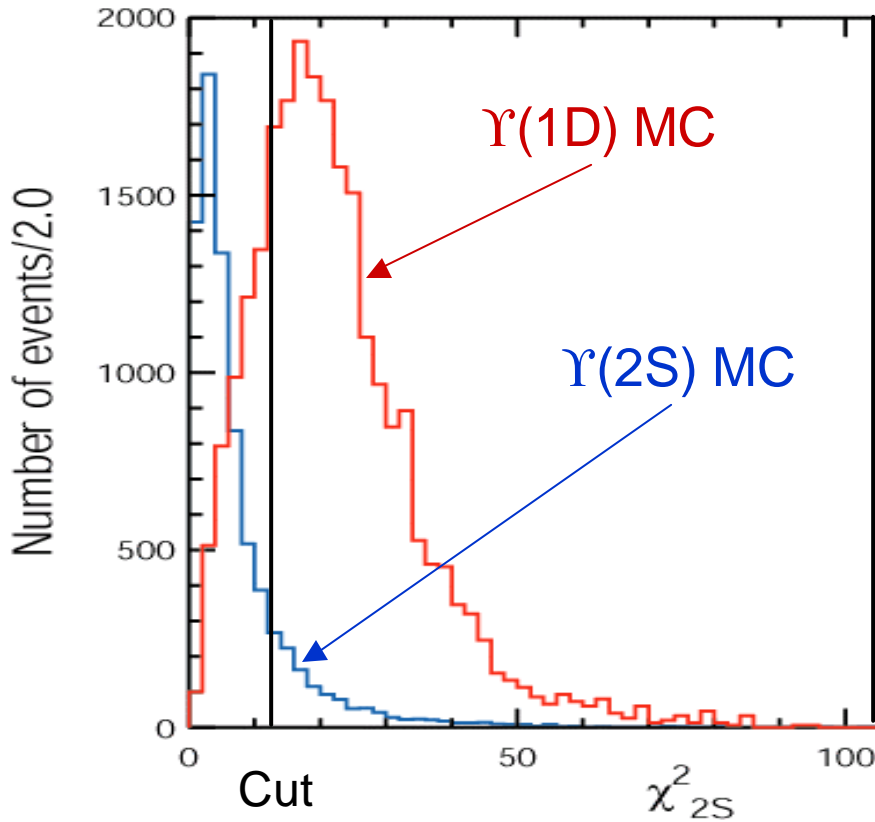


(Extra Slide)

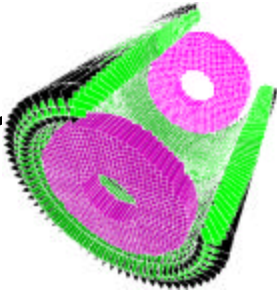
# Suppression of the $\Upsilon(2S)$ $4\gamma$ -cascades

$$c_{2S}^2 = \min_{J_{2P}, J_{1P}} \sum_{i=1}^4 \left( \frac{E_{gi} - E_{gi}^{\text{expected}}(M_{2S}, J_{2P}, J_{1P})}{s(E_{gi})} \right)^2 \quad c_{2S}^{2+} = \min_{J_{2P}, J_{1P}} \sum_{i=1}^4 \left( \frac{\max\{0, E_{gi} - E_{gi}^{\text{expected}}(M_{2S}, J_{2P}, J_{1P})\}}{s(E_{gi})} \right)^2$$

$c_{2S}^2 > 12$   $c_{2S}^{2+} > 3$







(Extra Slide)

# Signal for $\Upsilon(3S) \rightarrow \gamma \chi_b(2P_0) \rightarrow \gamma \gamma \Upsilon(2S)$

- Projection of 2D-fit to  $\delta = (\text{Recoil Mass}(2\gamma) - M_{\Upsilon(2S)})/\sigma$  vs.  $E_{\gamma_{\text{Low}}}$  onto  $E_{\gamma_{\text{Low}}}$  with the  $\exp(-\delta^2/2)$  as weight

