

# Diffractive Production of Vector Mesons at HERA

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## OUTLINE

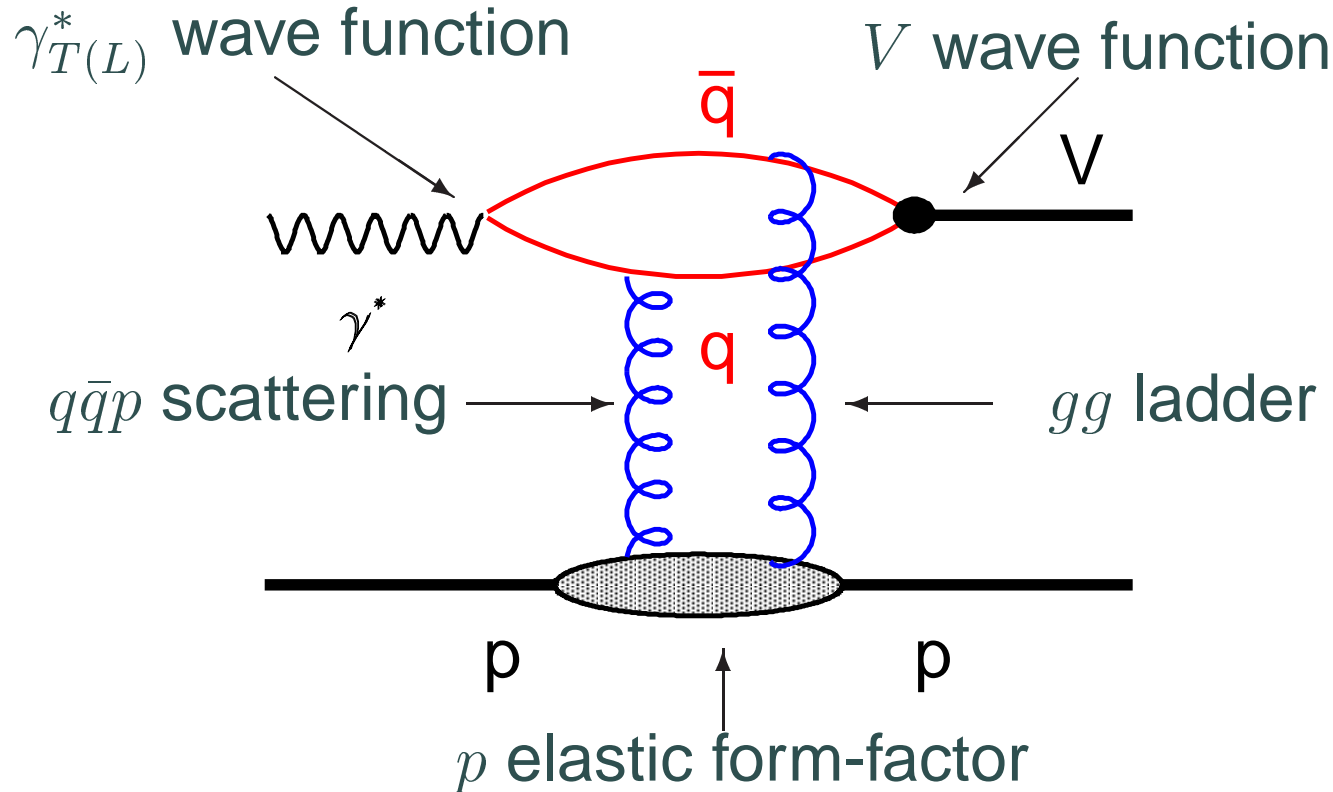
- **Introduction**
- **Motivation**
- **Experimental results:**  $\rho^0$ ,  $\phi$ ,  $J/\psi$ ,  $\psi(2S)$

different kinematic regimes: photoproduction, high  $Q^2$ , high  $t$

- **Summary and Outlook**

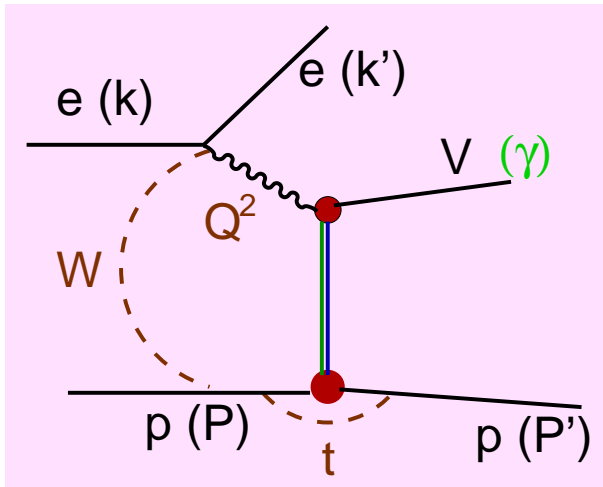
## INTRODUCTION

Aim is understand VM in QCD



- ⇒ understand dynamics of high energy scattering
- test pQCD in transition regime soft-hard
  - measure non-perturbative quantities (generalised) pdfs

## KINEMATICS



- $Q^2$  - virtuality of exchanged  $\gamma^*$

$$Q^2 = -q^2 = -(k - k')^2$$

- $W$  -  $\gamma^* p$  centre of mass energy

$$W = (q + p)^2$$

- $t$  - 4-momentum transfer squared at the  $p$  vertex

$$t = (P - P')^2$$

- $x$  - Bjorken variable

$$x = \frac{Q^2}{P \cdot q} = \frac{Q^2}{Q^2 + W^2}$$

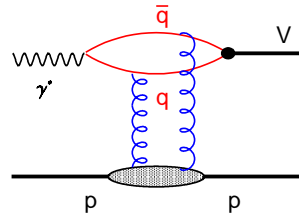
## MOTIVATION

**HERA regime:** collisions of 27.5 GeV  $e$  with 820 (920) GeV  $p$   
 $0 < Q^2 < 50 \text{ GeV}^2$  and  $30 < W < 300 \text{ GeV}$

- Large  $W$ , large  $Q^2$ , but  $Q^2 \ll W^2$   
⇒ strong interactions in the presence of a large scale (unitarity limit?)
- Small  $x$   
⇒ large density of partons, expect coherent effects
- Transition from soft to hard interactions  
⇒ transition from a hadron to a state of partons (confinement?)

# EXCLUSIVE VECTOR MESON IN $ep$

Candidate for a Hard process

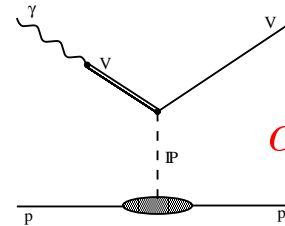


small  $q\bar{q} \Rightarrow$  either  $\gamma_L^*$  or  
 $V = c\bar{c}$  or  $b\bar{b}$

$$\sigma_L \sim \frac{\alpha_S^2}{Q^6} |xG(x, Q^2)|^2$$

- $\Rightarrow$  fast increase of  $\sigma_L$  with  $W^2$
- $\Rightarrow$   $Q^2$  dependence slower than  $1/Q^6$
- $\Rightarrow$  universal  $t$  dependence  $\sim e^{b_{2g}t}$   
 $b_{2g} \simeq 4 \text{ GeV}^{-2}$  and  $\alpha'_{IP} \simeq 0$
- $\Rightarrow \rho : \omega : \Phi : J/\psi = 9 : 1 : 2 : 8$

In the soft Pomeron picture based on Regge phen. and VDM



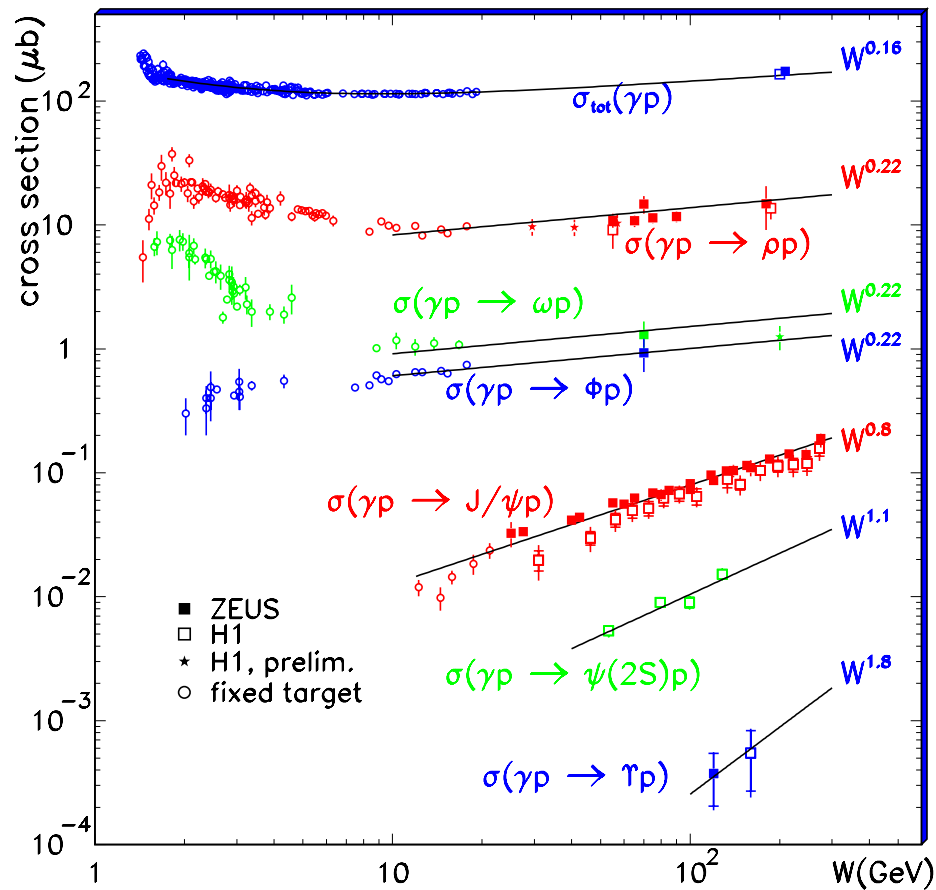
$$\alpha_{IP}(t) = \alpha_{IP}(0) + \alpha'_{IP} \cdot t$$

$$\sigma_{\text{tot}} \sim s^{\alpha_{IP}(0)-1}$$

$$\frac{d\sigma_{\text{el}}}{dt} \sim \frac{\sigma_{\text{tot}}^2}{16\pi} e^{2(b_0^{\text{el}} + \alpha'_{IP} \ln s)t}$$

- $\Rightarrow \alpha_{IP}(0) = 1 + \epsilon \simeq 1.10$   
and  $\alpha'_{IP} = 0.25 \text{ GeV}^{-2}$
- $\Rightarrow \frac{\sigma_{\text{el,D}}}{\sigma_{\text{tot}}} \sim s^\epsilon$
- $\Rightarrow$  shrinkage of the  $t$  slope

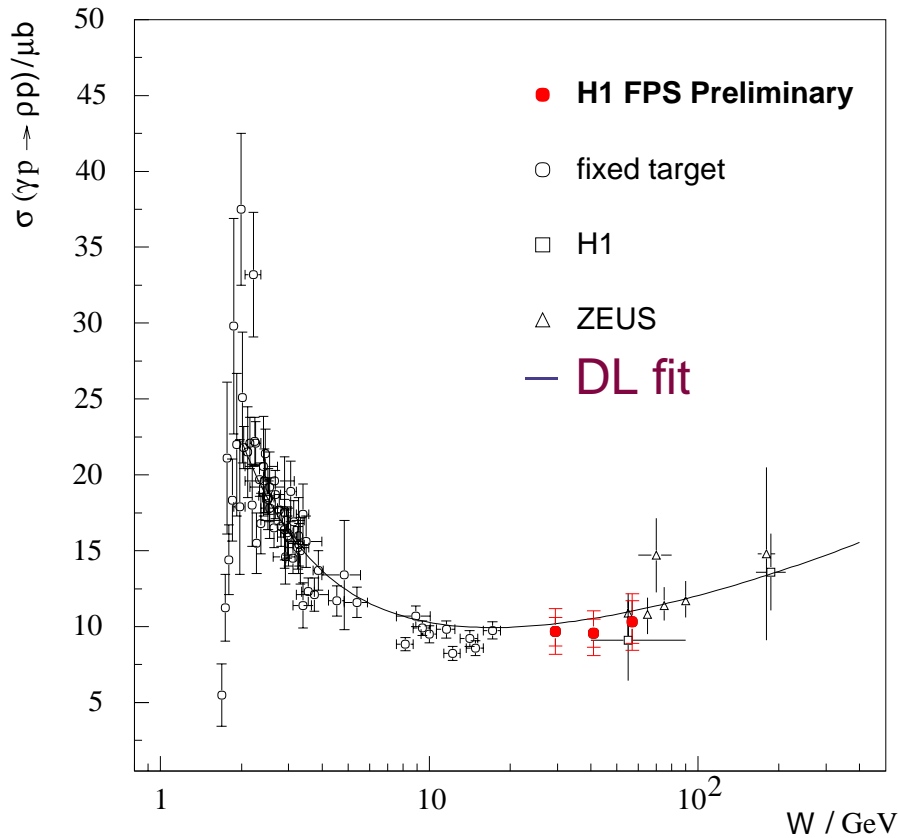
# EXCLUSIVE VECTOR MESON IN $\gamma p$



$\Rightarrow$  at  $Q^2 = 0$ , change of regime with mass

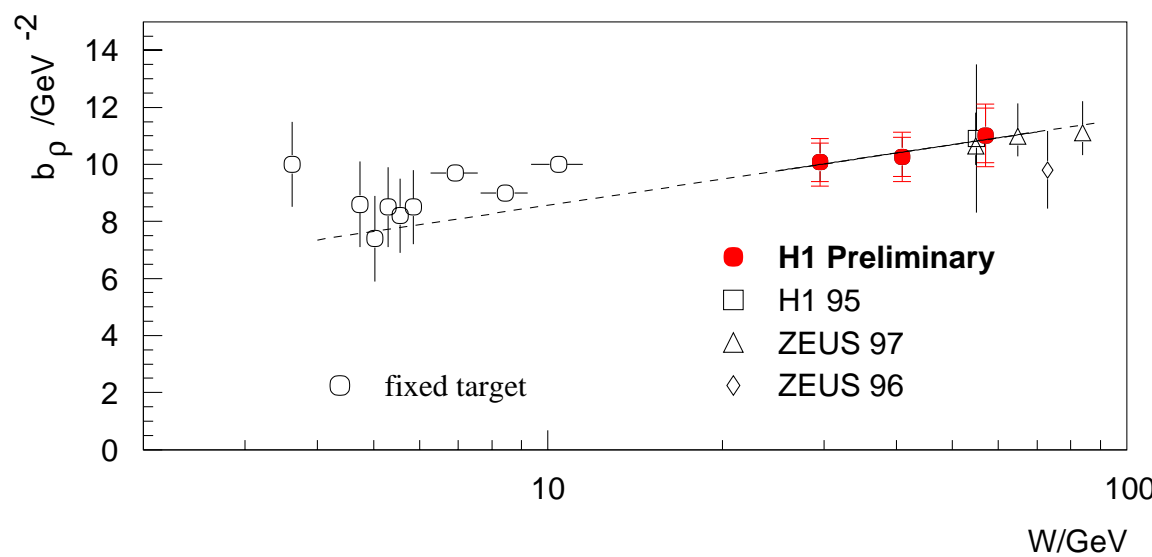
# EXCLUSIVE $\rho^0$ MESON IN $\gamma p$

$p$  measured in forward proton spectrometer



$$\sigma_{\gamma p \rightarrow \rho p} \propto W^{0.22}$$

Fit  $\frac{d\sigma}{dt} \propto e^{b \cdot t}$

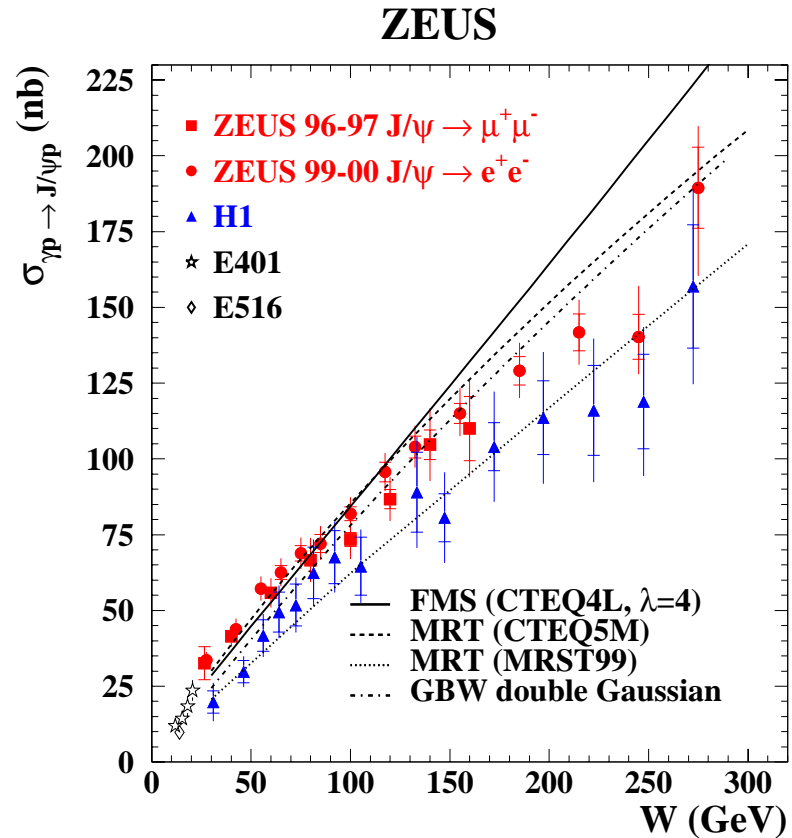
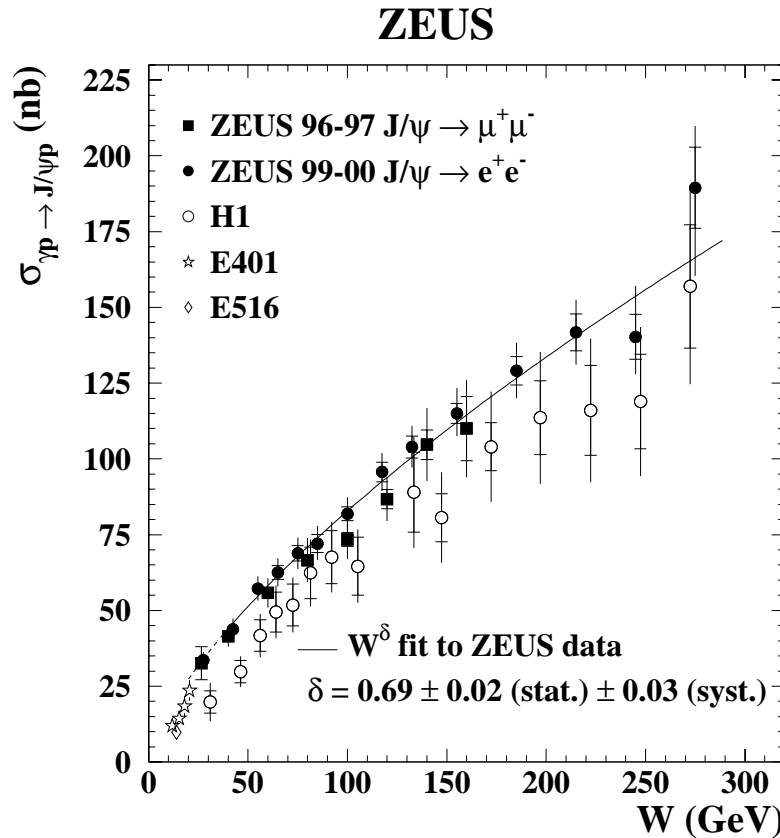


$$b_\rho \propto \ln(W)$$

⇒ consistent with Regge phenomenology

# EXCLUSIVE $J/\psi$ IN $\gamma p$

$$\sigma(\gamma p \rightarrow J/\psi p)$$

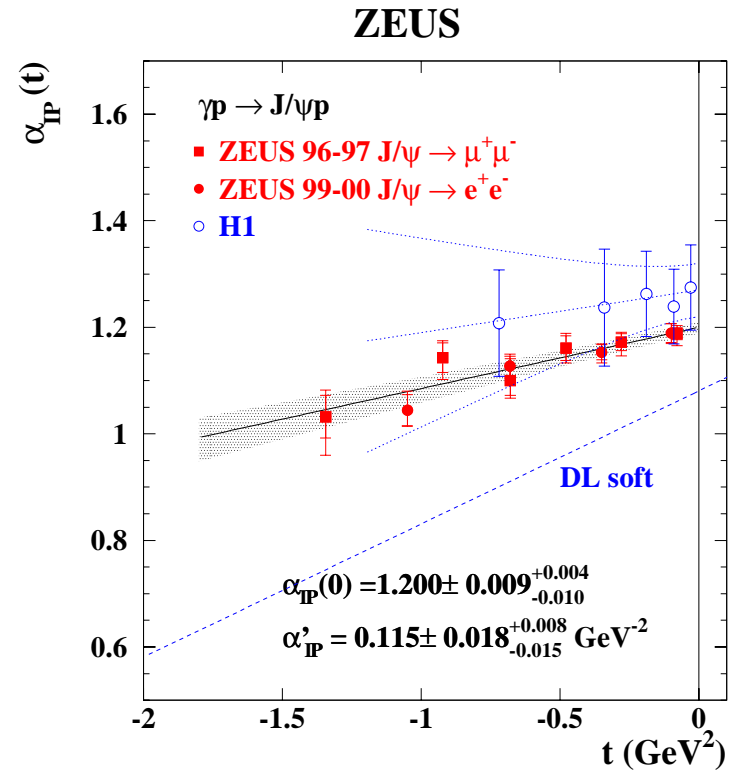
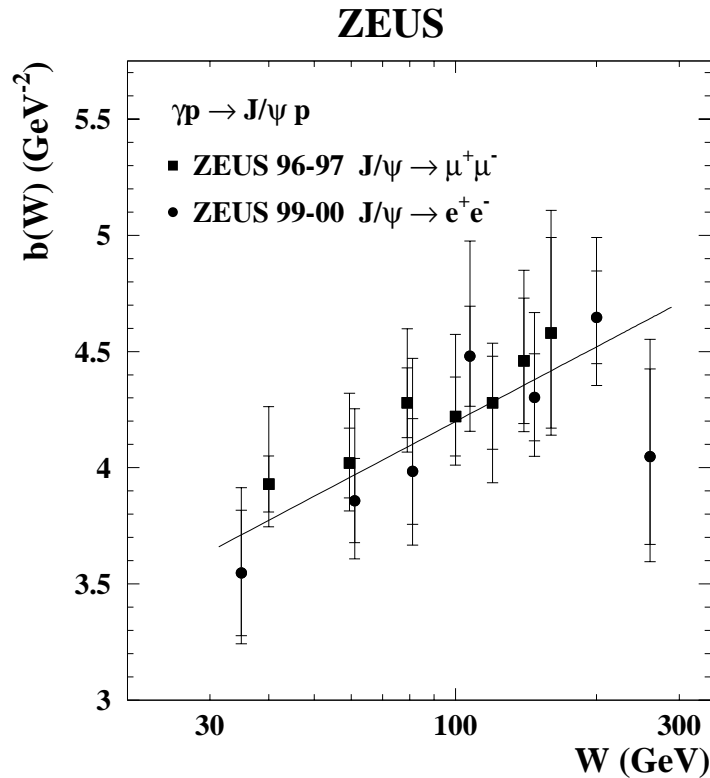


- Fast rise with  $W$ , compatible with pQCD calculations
- Deconvolution of  $xg(x, Q^2)$  from data still not possible



# EXCLUSIVE $J/\psi$ IN $\gamma p$

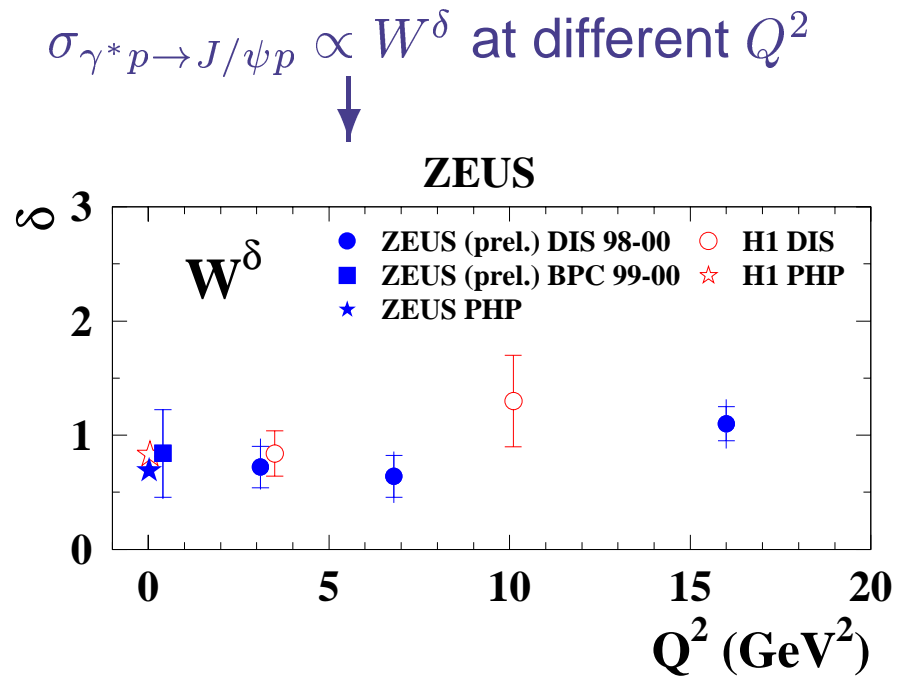
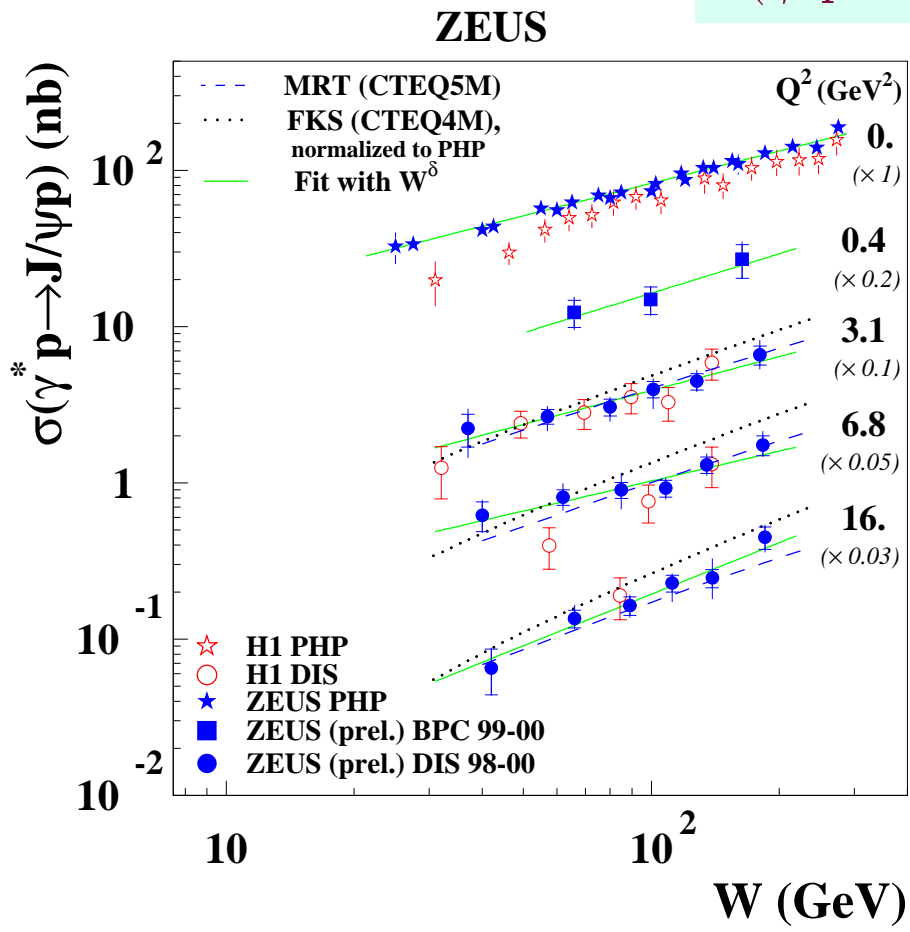
## $b$ slopes and $IP$ trajectory



- $b \simeq 4 \text{ GeV}^2 \Rightarrow$  point-like interaction
- Slight rise of  $b$  with  $W$

# EXCLUSIVE $J/\psi$ IN $\gamma^*p$

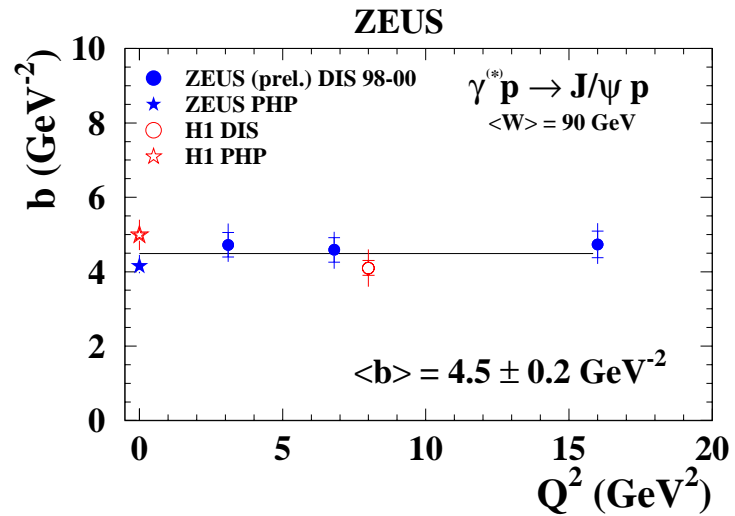
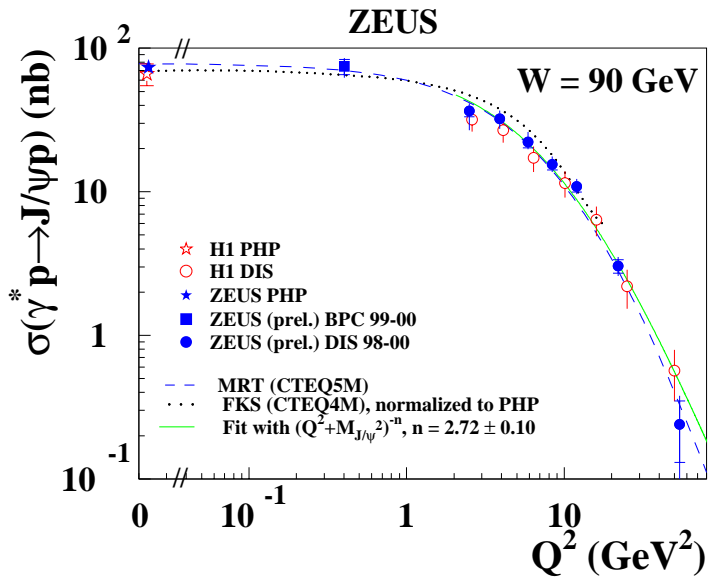
$$\sigma(\gamma^*p \rightarrow J/\psi p)$$



●  $\delta \sim$  flat with  $Q^2$

# EXCLUSIVE $J/\psi$ IN $\gamma^*p$

$\sigma(\gamma^*p \rightarrow J/\psi p)$  and  $b$  vs  $Q^2$

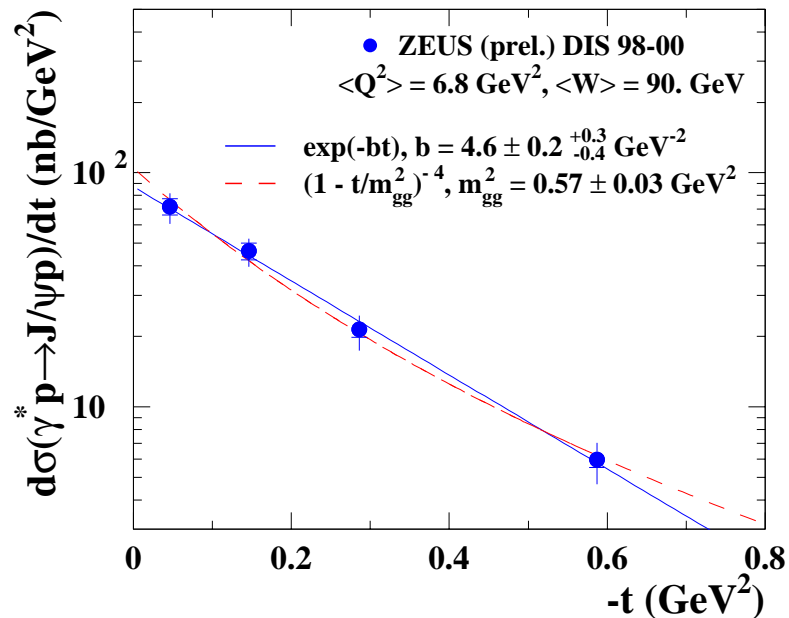


$\Rightarrow J/\psi$  production has all the characteristics of a hard process

# EXCLUSIVE $J/\psi$ IN $\gamma^*p$

“Proton elastic form-factor”

**ZEUS**

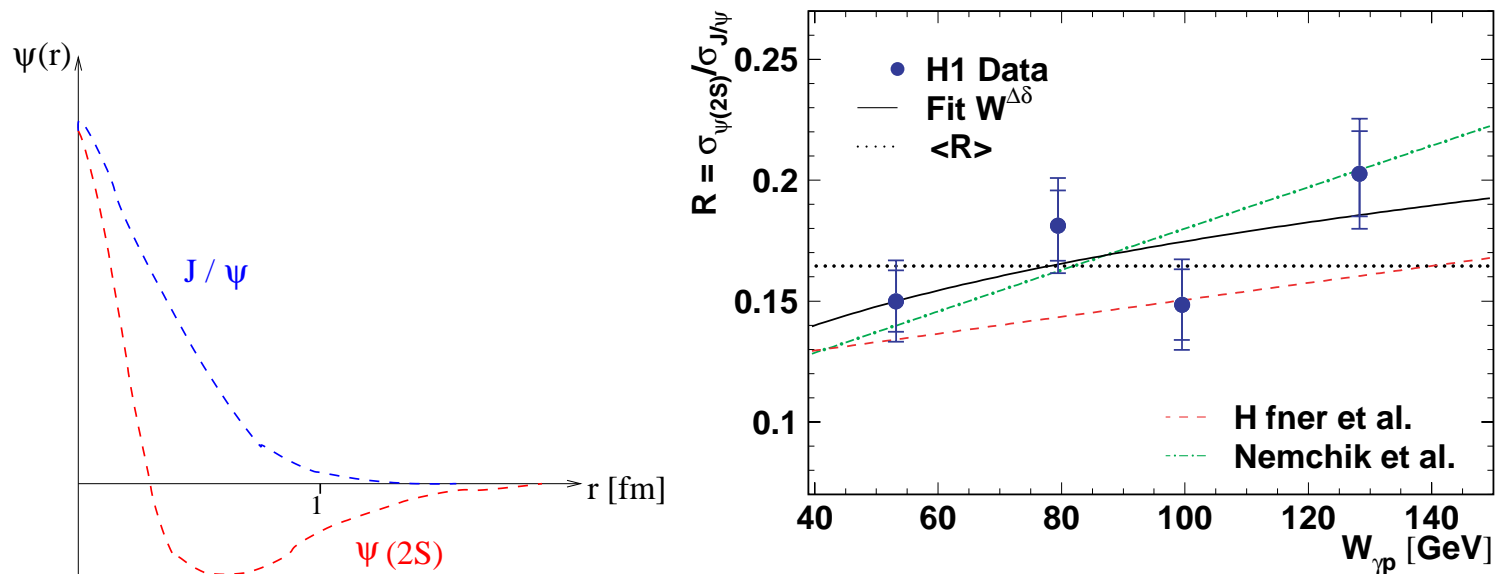


- In QCD, at high  $Q^2$ ,  $b$  related to transverse gluon distribution
- Form  $(1 - t/m_{2g}^2)^{-4}$  suggested as elastic form-factor of  $2g$  exchange by Frankfurt and Strikman

# DIFFRACTIVE $\psi(2S)$ IN $\gamma p$

Ratio  $R = \frac{\sigma(\psi(2S)p)}{\sigma(J/\psi p)}$  sensitive to the radial wave function of charmonium

QCD prediction  $R \simeq 0.17$

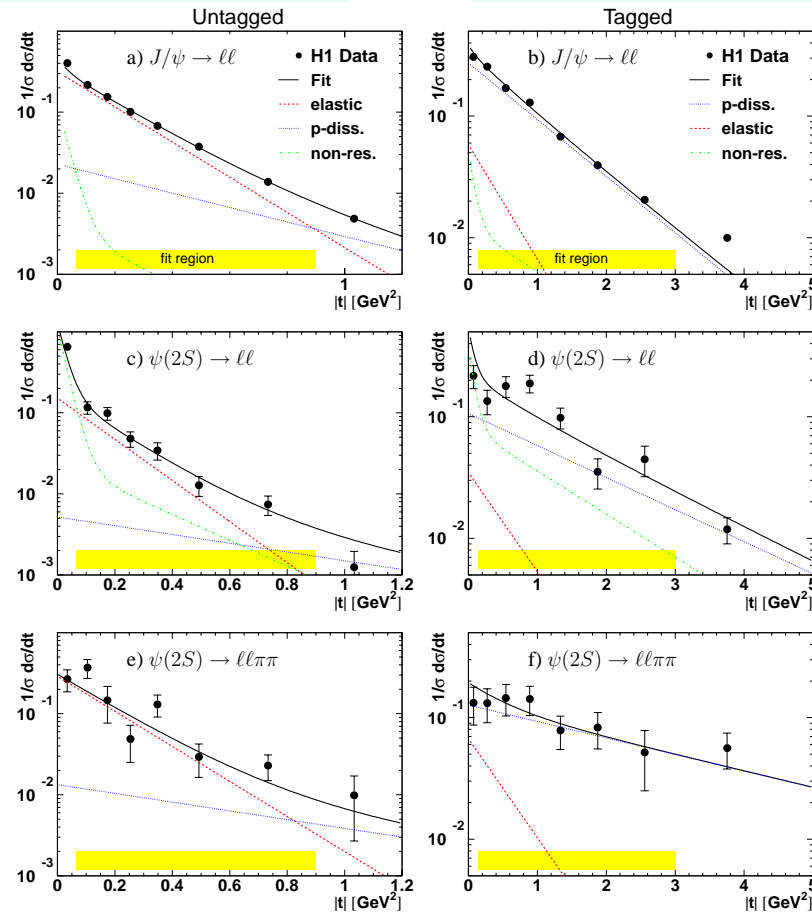


- $\psi(2S)$  production suppressed
- Energy dependence similar to  $J/\psi$

# DIFFRACTIVE $\psi(2S)$ IN $\gamma p$

$$d\sigma/dt(\gamma p \rightarrow \psi p)$$

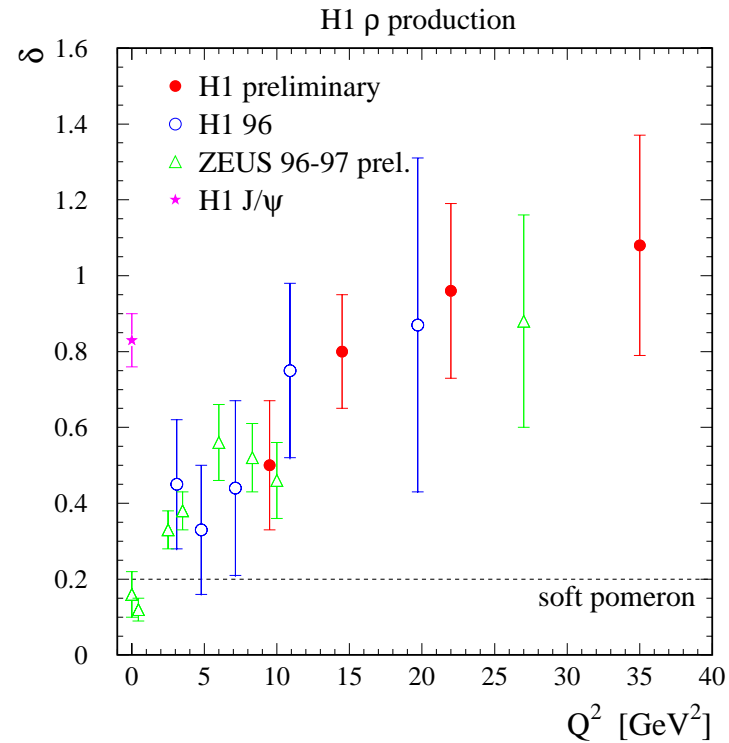
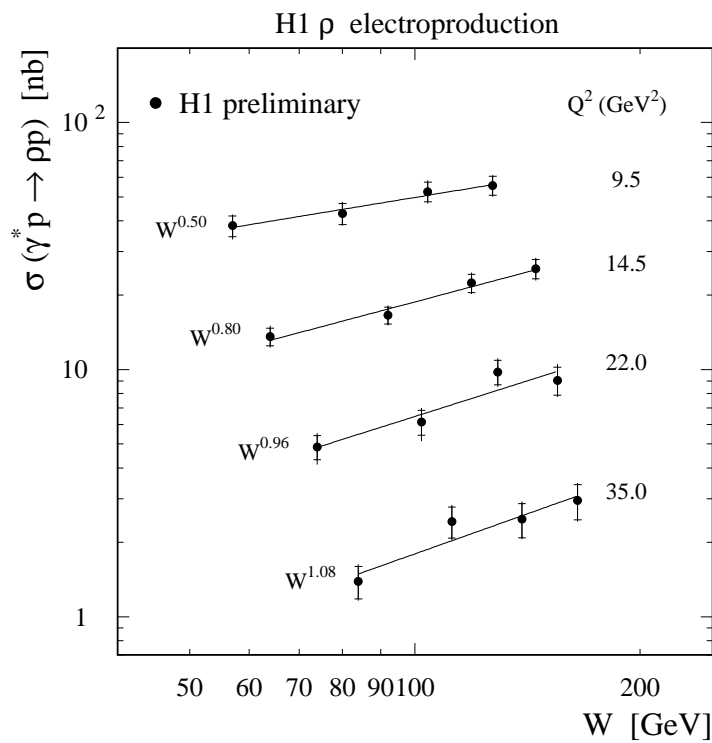
$$d\sigma/dt(\gamma p \rightarrow \psi Y)$$



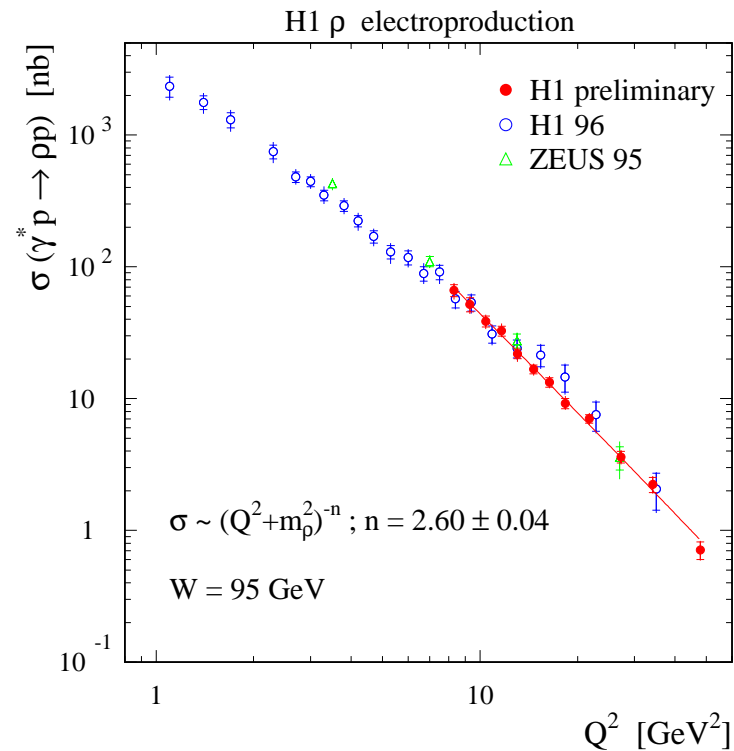
- t-dependence of  $J/\psi$  and  $\psi(2S)$  is similar

# EXCLUSIVE $\rho^0$ MESON IN $\gamma^*p$

## $W$ dependence of $\sigma(\gamma^*p \rightarrow \rho^0 p)$



$\sigma_{\gamma^*p \rightarrow \rho p} \propto W^\delta$  at different  $Q^2$ ,  $\Rightarrow \delta$  increases with  $Q^2$

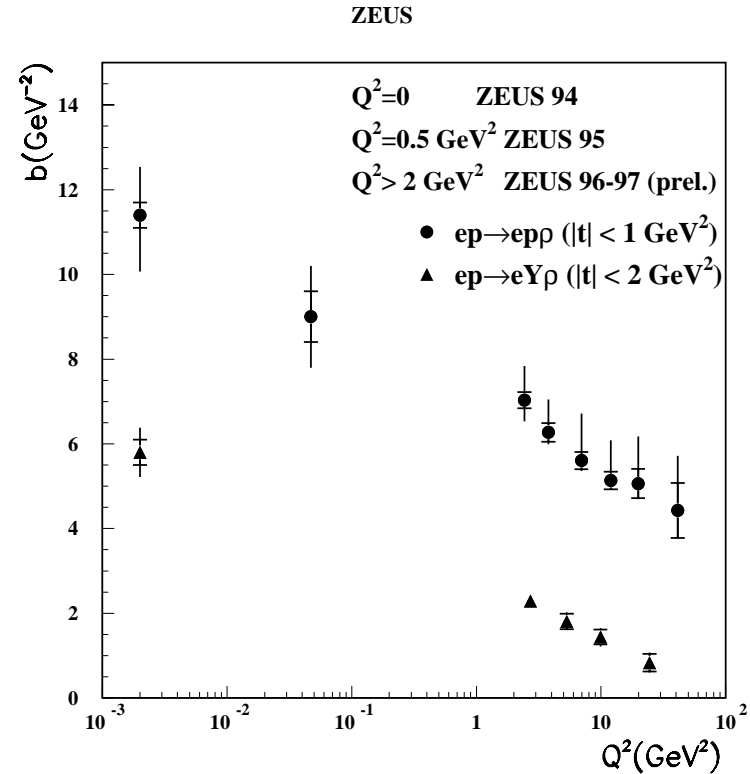
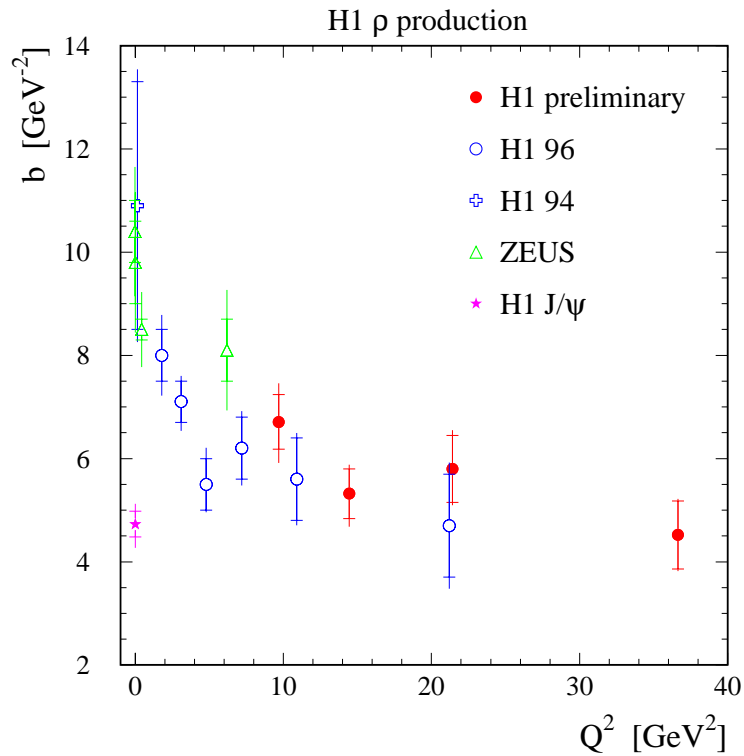
EXCLUSIVE  $\rho^0$  MESON IN  $\gamma^* p$  $Q^2$  dependence of  $\sigma(\gamma^* p \rightarrow \rho^0 p)$ 

$$\Rightarrow \text{pQCD: } \sigma_L \sim \frac{\alpha_S^2}{Q^6} |xG(x, Q^2)|^2$$



$\rho^0$  MESON IN  $\gamma^*p \rightarrow \rho^0p$  AND  $\gamma^*p \rightarrow \rho^0Y$

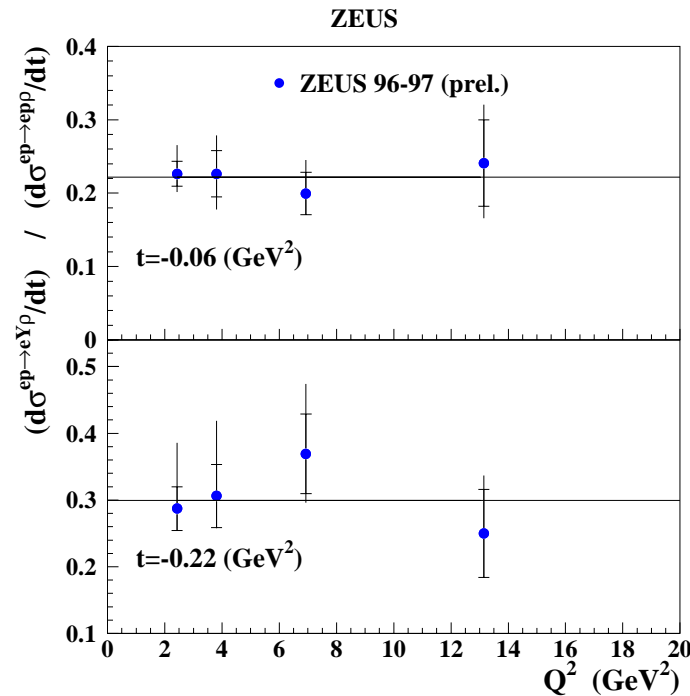
$Q^2$  dependence of  $b$



$\Rightarrow b$  decreases with  $Q^2$  for both exclusive and p-diss. reaction

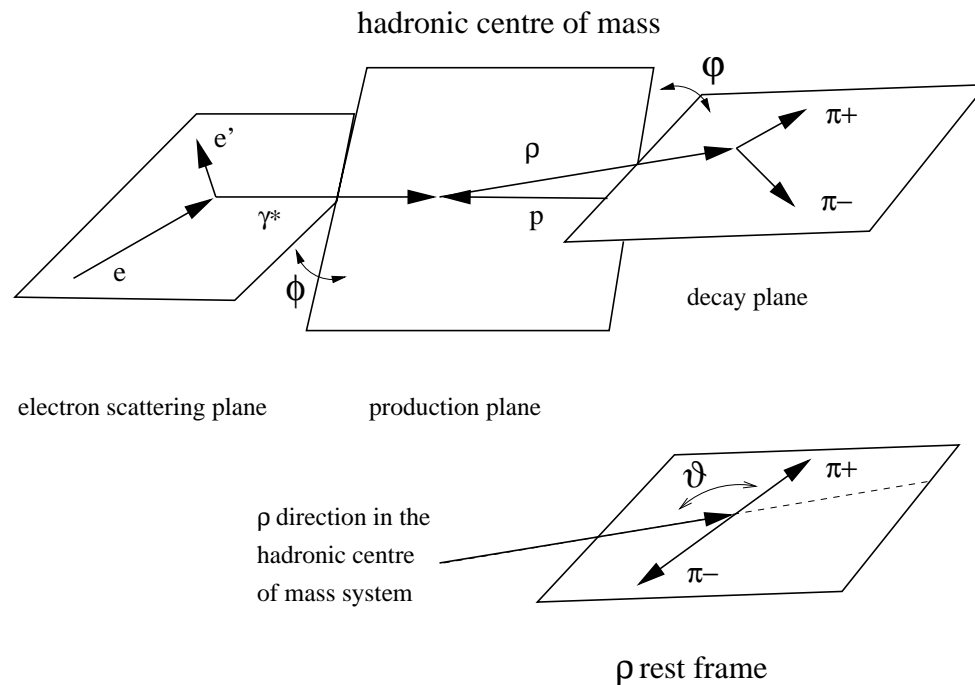
$\rho^0$  MESON IN  $\gamma^*p \rightarrow \rho^0p$  AND  $\gamma^*p \rightarrow \rho^0Y$

$Q^2$  dependence of  $d\sigma/dt$



⇒ Vertex factorization holds: probability of proton dissociation is independent of the projectile

# DECAY ANGULAR DISTRIBUTIONS



## Helicity angles

$\theta_h, \phi_h$  - angles of decay particle in the meson rest frame

$\Phi$  - angle between scattering and production plane

- Angular distributions are related to the spin of  $\gamma^*$  and meson

Angular distr.  $\rightarrow$  spin density matrix elements  $r_{ij}^{kl} \rightarrow$  helicity amplitudes  $T_{\lambda_V M \lambda_\gamma}$

## DECAY ANGULAR DISTRIBUTIONS

### Spin Matrix Elements

s-channel helicity conservation (SCHC):

- the VM retains the  $\gamma^*$  helicity.  $R = \sigma_L/\sigma_T$  is related to the spin density matrix elements  $r_{00}^{04}$  (good approximation).

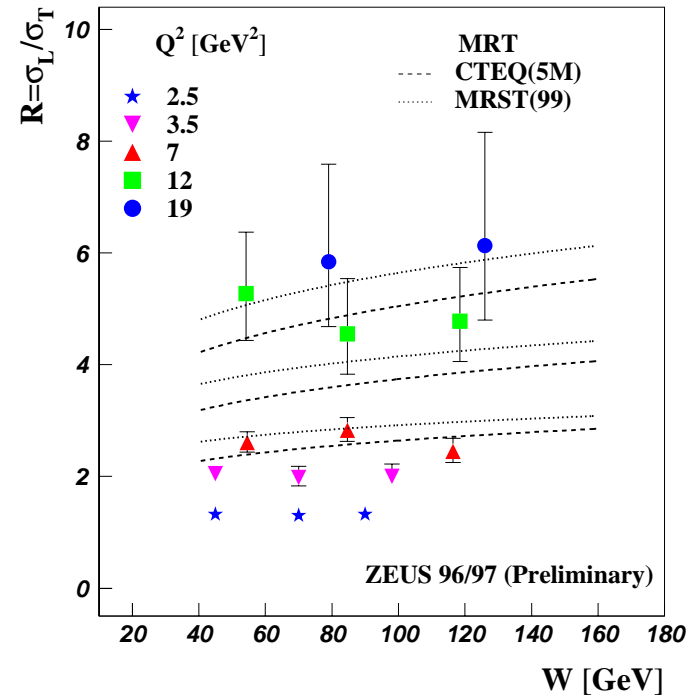
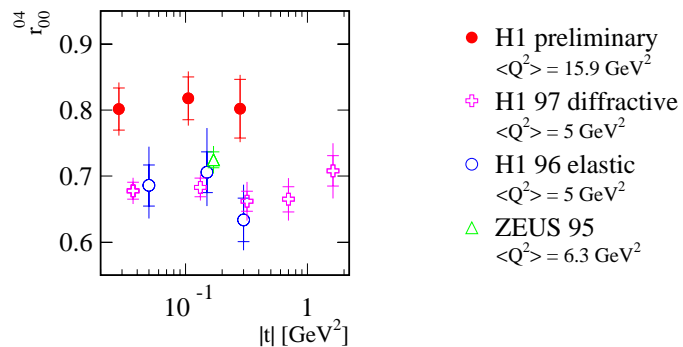
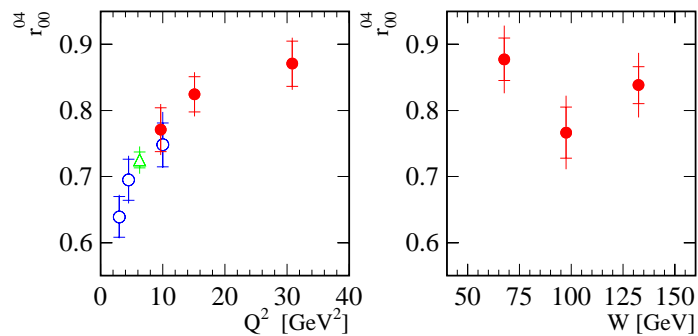
pQCD:

- during the interaction, the orbital angular momentum of the  $q\bar{q}$  can be modified through the transfer of transverse momentum carried by gluons;
- the helicity of the outgoing vector meson can be different from that of the incoming photon, helicity flip between photon and meson is possible.

# EXCLUSIVE $\rho^0$ MESON IN $\gamma^*p$

Spin Matrix Elements  $\sigma(\cos(\theta)) \rightarrow r_{00}^{04}$

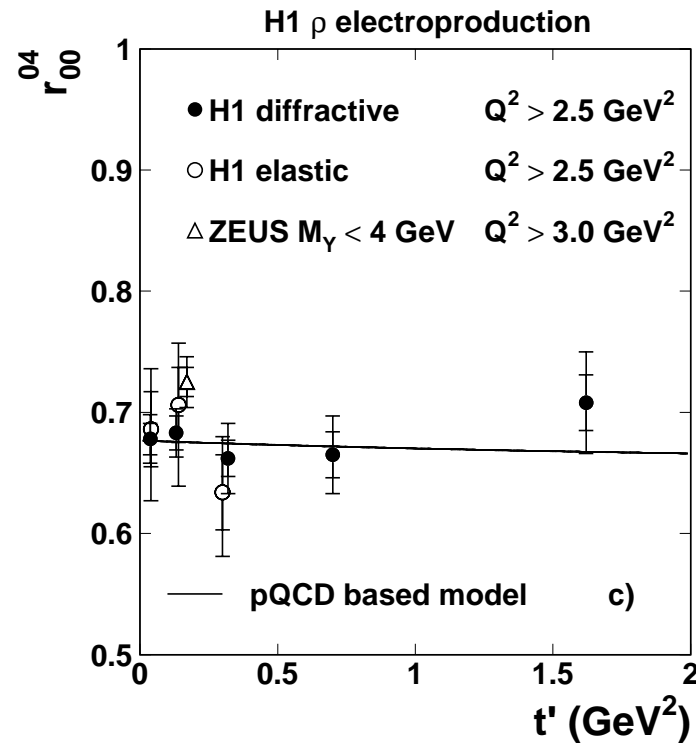
$$R = \sigma_L / \sigma_T \simeq \frac{r_{00}^{04}}{\epsilon(1-r_{00}^{04})}$$



•  $\sigma_L$  and  $\sigma_T$  have the same  $W$  and  $t$  dependence

$\rho^0$  MESON IN  $\gamma^*p \rightarrow \rho^0 Y$

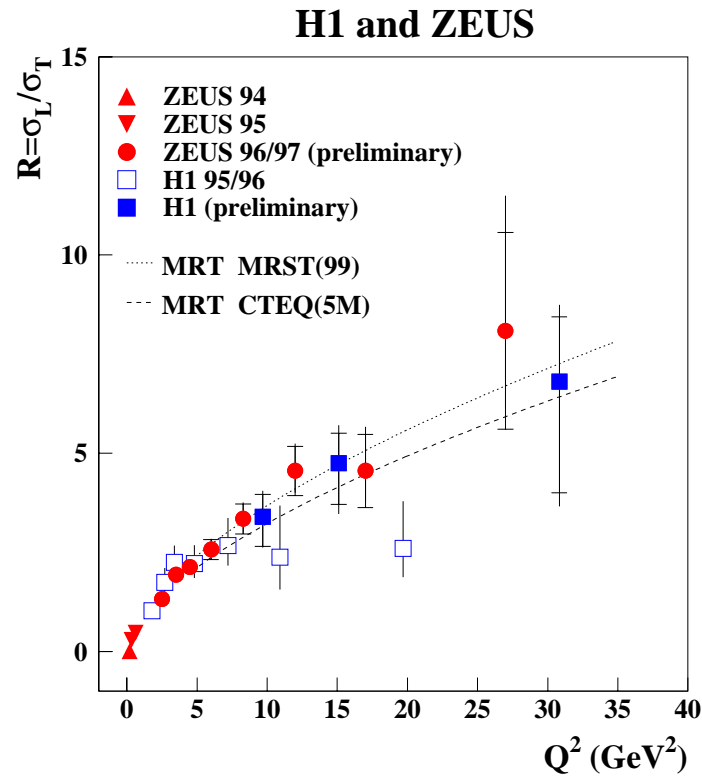
Spin Matrix Elements,  $\sigma(\cos(\theta)) \rightarrow r_{00}^{04}$



- $\sigma_L$  and  $\sigma_T$  have the same  $t$  dependence

# EXCLUSIVE $\rho^0$ MESON IN $\gamma^*p$

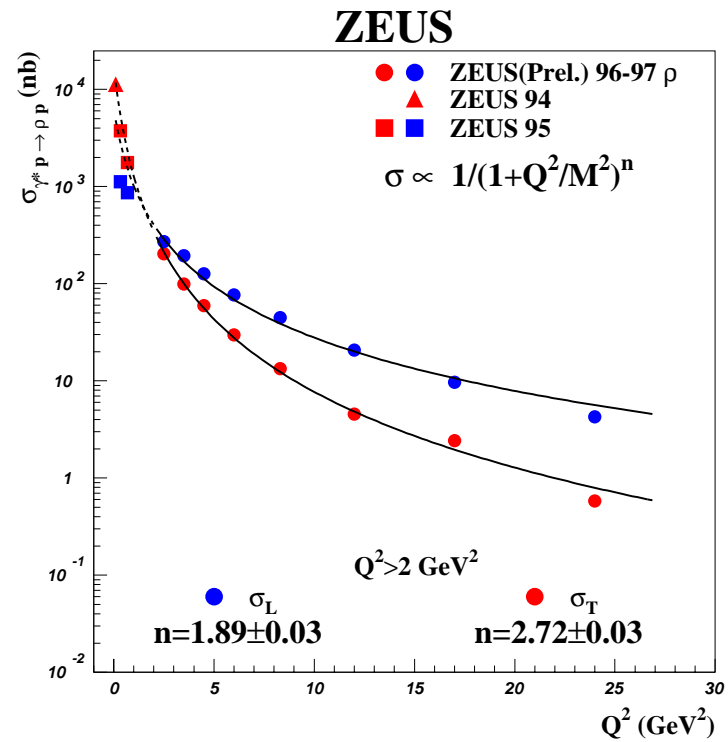
$$R = \sigma_L / \sigma_T \simeq \frac{r_{00}^{04}}{\epsilon(1-r_{00}^{04})}$$



⇒ Continuous rise of  $R$  with  $Q^2$

# EXCLUSIVE $\rho^0$ MESON IN $\gamma^*p$

$Q^2$  dependence of  $\sigma(\gamma^*p \rightarrow \rho^0 p)$



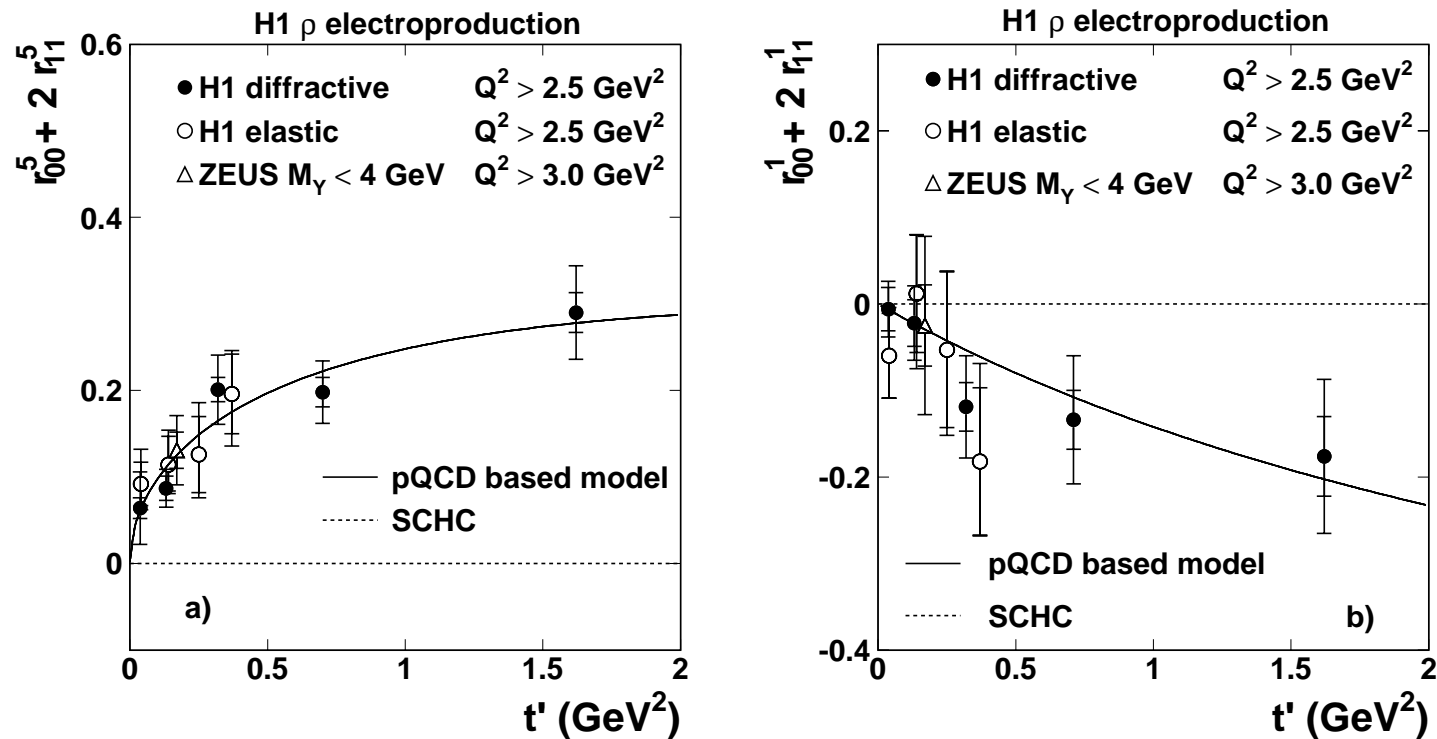
$\Rightarrow$  As expected different  $n$  for  $\sigma_L$  and  $\sigma_T$



$\rho^0$  MESON IN  $\gamma^*p \rightarrow \rho^0 Y$

$\sigma(\cos(\Phi)) \rightarrow r_{00}^5 + 2r_{11}^5$ , sensitive to single flip

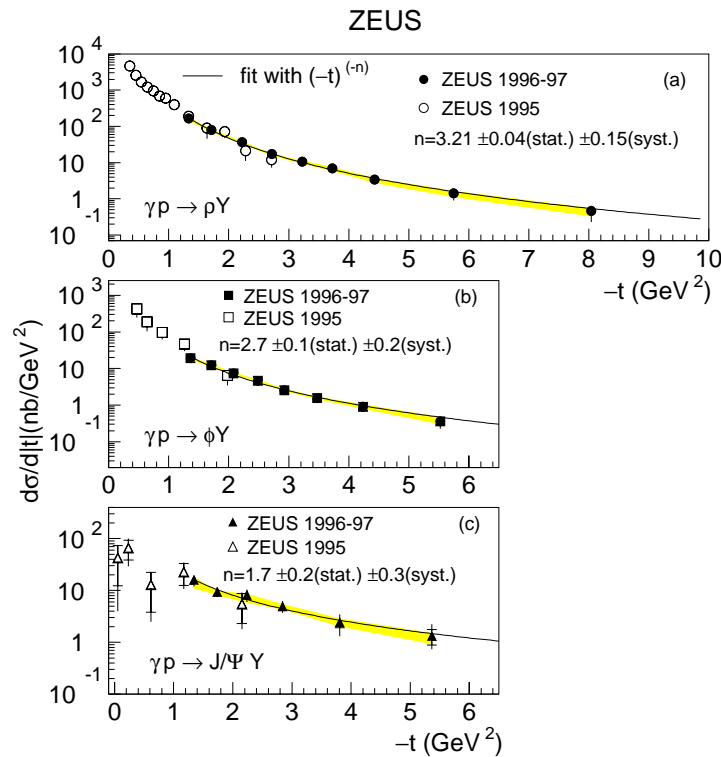
$\sigma(\cos(2\Phi)) \rightarrow r_{00}^1 + 2r_{11}^1$ , sensitive to double flip



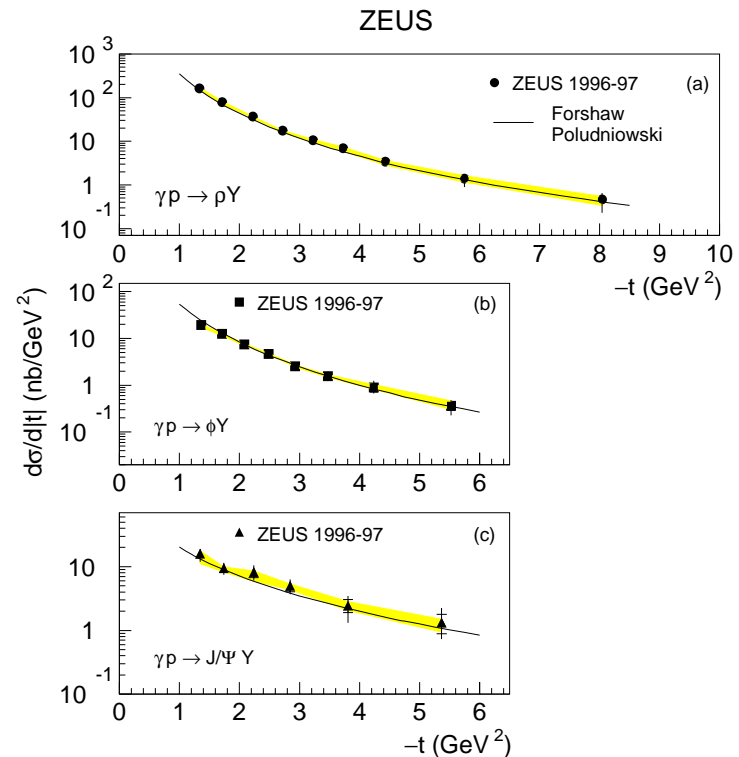
- Ratio of single flip to non flip amplitudes increases with  $t'$
- Probability of double flip is still small in this  $t$ -kinematic domain

# HIGH- $t$ $\rho^0$ , $\phi$ AND $J/\psi$ IN $\gamma p$

## $t$ dependence of $\sigma(\gamma p \rightarrow VY)$



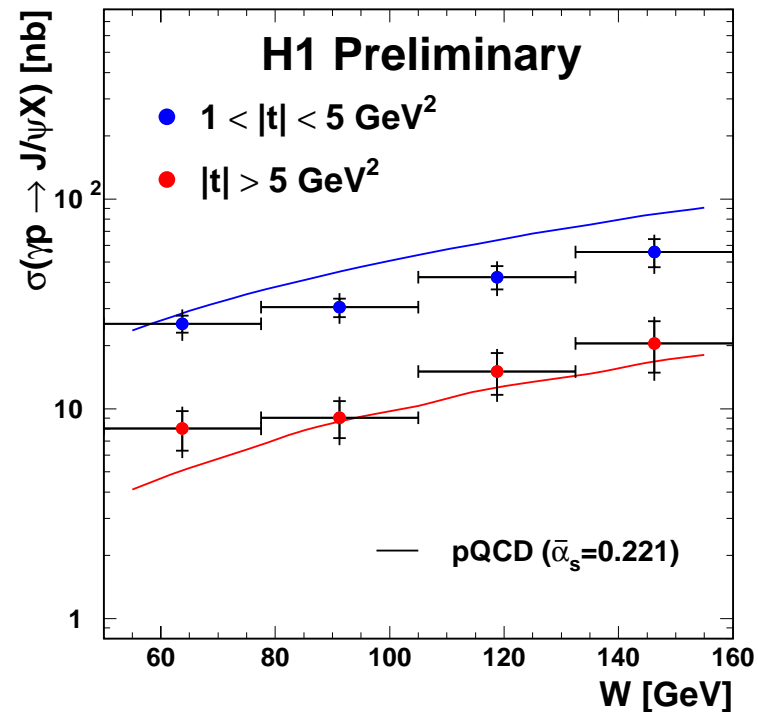
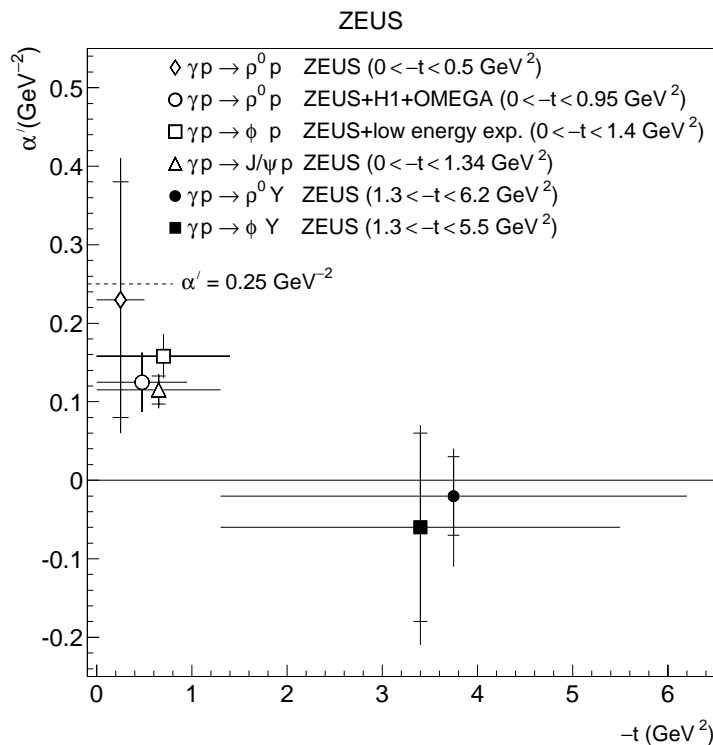
⇒ power-like behavior



⇒ BFKL-type behavior

# HIGH- $t$ $\rho^0$ , $\phi$ AND $J/\psi$ IN $\gamma p$

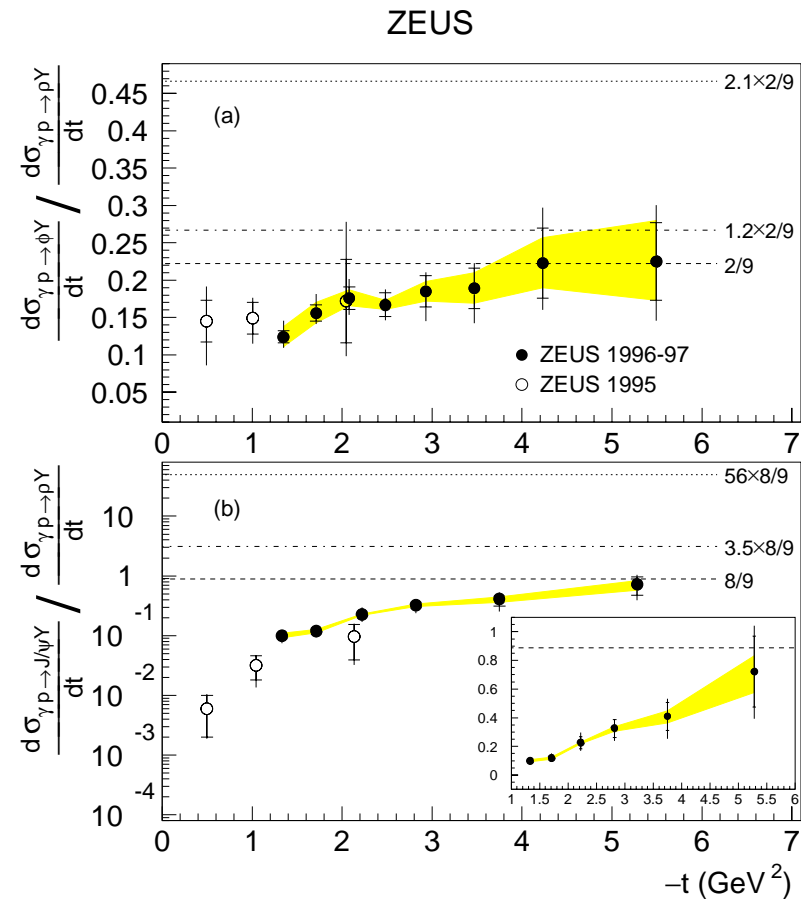
## $t$ dependence of $\alpha'_P$ and $W$



$\alpha'_P$  small  $\equiv W$  dependence not change with  $t$ , described by pQCD  
 $\Rightarrow t$  provides an hard scale;  $\alpha_P(t)$  is not linear in  $t$

# HIGH- $t$ $\rho^0$ , $\phi$ AND $J/\psi$ IN $\gamma p$

$\sigma_{\phi, J/\psi} / \sigma_{\rho}$  ratios at high  $t$  and  $SU(4)$



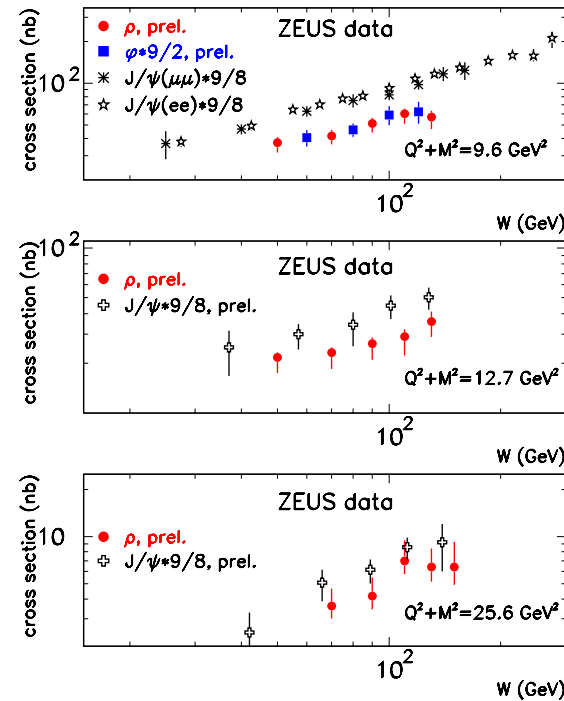
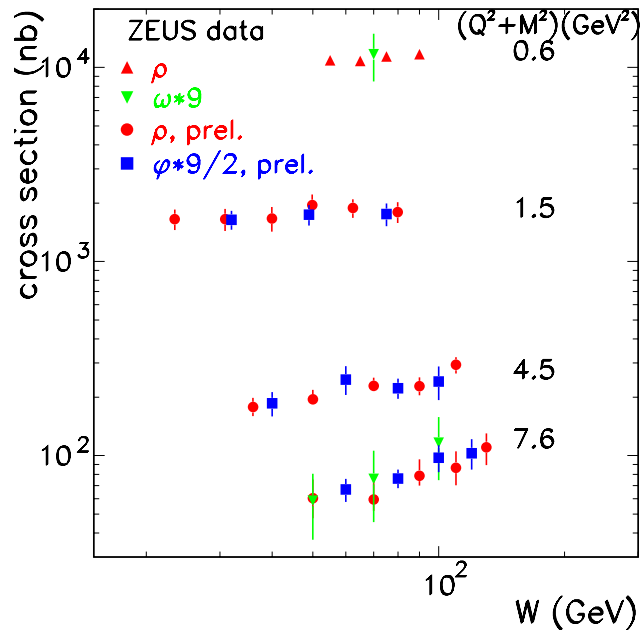
- Indication of VM production flavour independent at high  $t$

# OVERVIEW OF VECTOR MESON PRODUCTION

$SU(4)$ ?

- Scaling of  $\sigma(\gamma^*p \rightarrow Vp)$  in  $Q^2 + M_V^2$  proposed by H1

$\Rightarrow$  does not work for  $J/\psi$

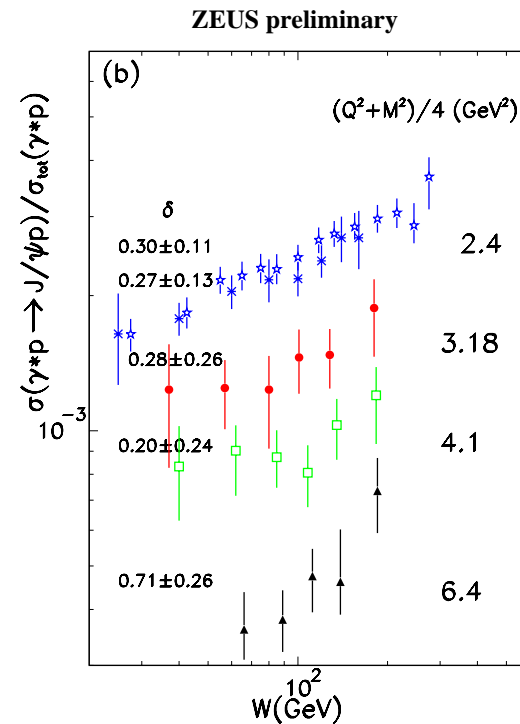
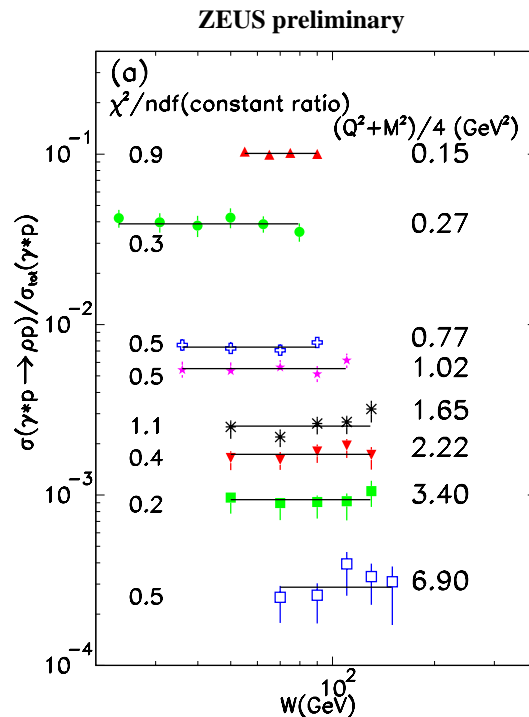


- No simple universality for  $VM$  if  $Q^2 + M^2$  is used as a scale
- Naive  $SU(4)$  may be altered by  $VM$  wave function effects

# OVERVIEW OF VECTOR MESON PRODUCTION

$$r_V = \frac{\sigma(\gamma^*p \rightarrow Vp)}{\sigma_{\text{tot}}(\gamma^*p)} \text{ vs } W$$

Expectation:  $\sigma_{\text{tot}}(\gamma^*p) \propto W^\delta$ ,  $\sigma(\gamma^*p \rightarrow Vp) \propto W^{2\delta} \Rightarrow r_V \propto W^\delta$

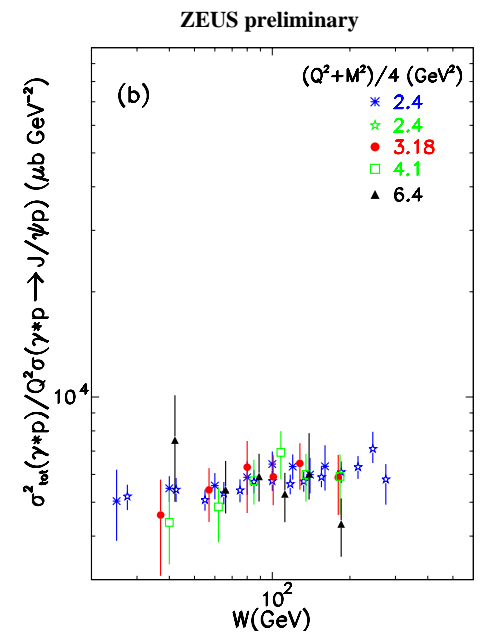
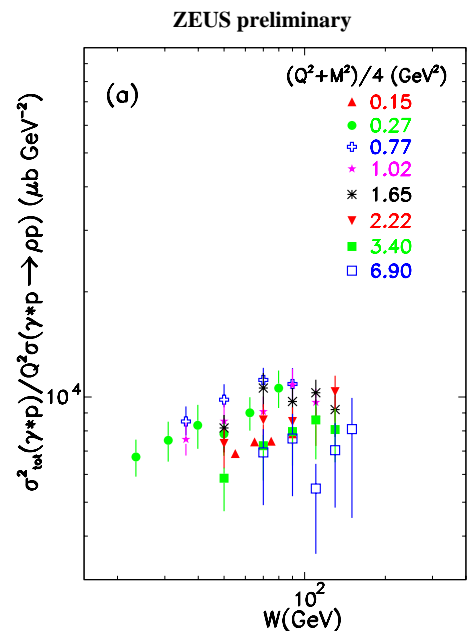


- $W$  independence of  $r_V$  for  $\rho^0$  can not be explained by pQCD or Regge; it is similar to inclusive diffraction

# OVERVIEW OF VECTOR MESON PRODUCTION

$$\frac{\sigma_{\text{tot}}^2(\gamma^*p)}{\sigma(\gamma^*p \rightarrow Vp)} \text{ vs } W$$

Typical for  $2g$  exchange:  $\frac{\sigma_{\text{tot}}^2(\gamma^*p)}{\sigma(\gamma^*p \rightarrow Vp)} \propto \frac{1/(\alpha_S^2(Q^2)Q^4)}{1/(Q^6 \cdot b)} \propto Q^2 \cdot b$



- Confirms the perturbative/non-perturbative  $2g$  exchange as opposed to VDM mechanism

## Summary and Outlook

- Experimentally much progress has been achieved,
  - high precision in large kinematic region
- Theoretically the overall picture looks correct,
  - large uncertainties
  - full NLO calculations are missing
- For the near future:
  - increased statistics at high  $Q^2$  will help (HERA II),
  - $\Upsilon$  remains to be investigated (HERA II)